

1 Thermophysical properties and containerless
2 solidification of Al-22.5wt%Cu in reduced gravity
3 using the ISS-EML

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10 **Abstract**

11
12 A hypoeutectic Al-22.5wt%Cu sample was processed using the Electro-Magnetic Levitator on
13 board of the International Space Station within the frame of the European Space Agency
14 project NEQUISOL and THERMOLAB. In total, 21 cycles of melting and cooling were carried
15 out, with the last cycle yielding a primary undercooling of 20 K and a eutectic undercooling
16 of 35 K. The density, surface tension, and viscosity of the sample were measured using the
17 oscillating drop method and a high-speed camera. The evolution of the temperature was
18 recorded with a pyrometer. The oscillations of the drop appear modulated, indicative of a
19 nonlinear behavior caused by rotation, precession, nutation of the sample, and high initial
20 deformation. After being returned to Earth, preliminary investigation of the surface of the
21 sample showed unexpected features. The growth direction of the visible dendrites seemed
22 to differ from the characteristic $\langle 100 \rangle$ expected of primary α -Al. Furthermore, two distinct
23 eutectic morphologies were observed: a typical lamellar eutectic, as well as an undulated
24 structure, akin to what is formed in rapidly solidified eutectic Al-33wt%Cu droplets.

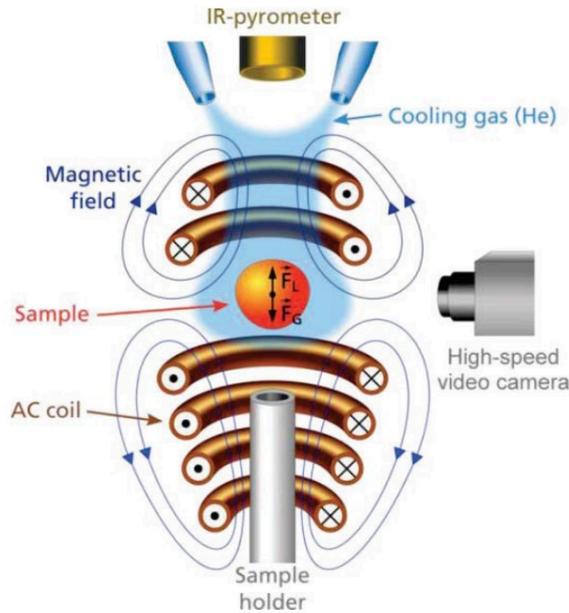
25 Introduction

26 The optimization of high temperature metallurgical processes such as atomization, welding
27 or casting requires building numerical models using the thermophysical properties of
28 metallic liquids as inputs. The development of an accurate database for the density, gas-
29 liquid surface tension, and viscosity of metals and alloys is then essential to understand and
30 analyze the physics of these processes [1]. The THERMOLAB project aims at measuring these
31 thermophysical properties, in weightlessness conditions, for samples processed in the
32 Electromagnetic Levitator onboard of the International Space Station (ISS-EML) [2]. In this
33 work, the microstructure and thermophysical properties of hypoeutectic Al-22.5wt%Cu are
34 investigated through containerless melting and solidification cycles carried out using the ISS-
35 EML equipment. First, the method to process the sample will be described, followed by the
36 presentation of the results for the measurements of the density, surface tension, and
37 viscosity. A discussion of the undercooled solidified microstructure will finally be presented.

38

39 Experimental Method

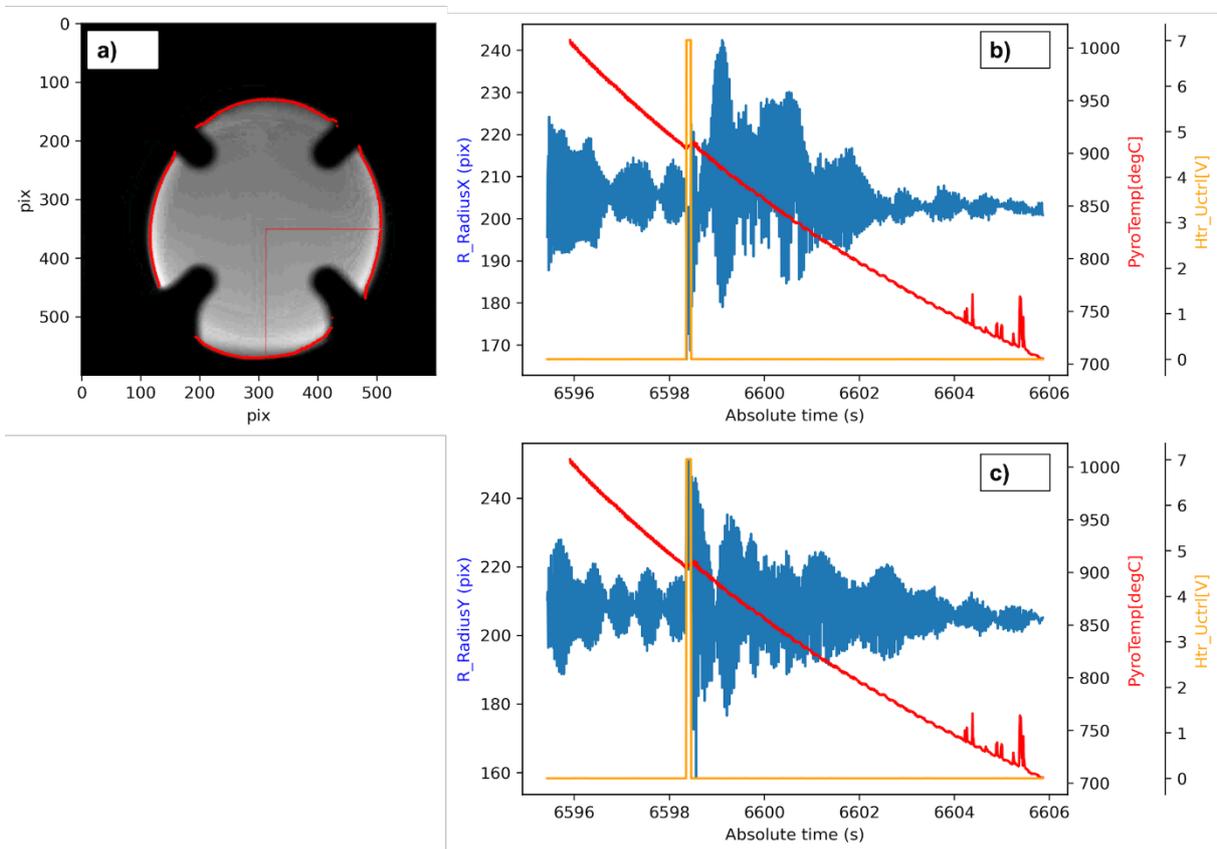
40 Electromagnetic levitation (EML) is a powerful solidification and thermophysical property
41 measurement technique for analysis of electrically conducting samples such as metals and
42 semiconductors. By avoiding contact with any container walls and operating under high
43 purity environment, heterogeneous nucleation is strongly reduced and a large range of
44 undercoolings can be achieved. In EML, the sample is placed within a conical levitation coil
45 typically consisting of five to seven water-cooled copper windings with one or two
46 counterwindings at its top as shown in Fig.1.



47
48 *Figure 1. Schematic of EML, taken from [3]*

49 Eddy currents are induced in the sample by the electromagnetic field generated by the
50 levitation coils. The sample is heated and melted by ohmic losses, whereas the interaction of
51 these eddy currents with the electromagnetic field leads to a displacement force on the
52 sample that is opposite to the gravitational force. The temperature of the sample is
53 monitored continuously with a contactless pyrometer. To cool the sample below its liquidus
54 temperature and induce solidification, a jet of high purity helium is used. Detailed
55 information on the EML technique can be found in [4]. A radial camera with a frame rate of
56 100 Hz is looking from the side of the sample with a window of 10x10 cm or 608x608 pixels.
57 This camera is combined with a pyrometer for temperature measurements. The Al-
58 22.5%wtCu sample was made from high purity Al-6N and Cu-6N and had a mass of $m =$
59 439.870 mg and a diameter around 6.4 mm. In total, 21 cycles of melting and solidification
60 were performed during the ISS experiments. The sample is assumed to be oxygen saturated.
61 An oxide layer was always observed at the start of melting, disappearing when reaching high
62 enough temperatures and reforming at the end of the cooling cycles. The data was trimmed
63 of the end part when the oxide layer was visibly forming again as it would impact the surface
64 tension measurement. For each cycle, the available dataset consists of the combination of
65 the video recorded by the radial camera and the temperature profile measured by the
66 pyrometer. The density, ρ , the surface tension, σ , and the viscosity, η , of the sample are
67 estimated using the oscillating drop technique [5]. The oscillations of the sample are
68 triggered by a short pulse stretching the drop along the heating field and resulting in

69 damped surface oscillations until reaching the equilibrium spherical shape. A cycle typically
 70 takes around 2 minutes and the total mass loss for the sample was 0.760mg after all the
 71 cycles. The droplet oscillations were monitored by the radial video camera and the images
 72 were first analyzed using a dedicated software “TeVi” (SEA Datentechnik GmbH), which
 73 detects the edges of the sample from the image contrast as well as the horizontal and
 74 vertical radii, R_x and R_y , as shown in Fig.2.a. However, the strong initial deformation causes
 75 the sample to rotate around a varying axis and translate which leads to inconsistent
 76 oscillations. Fig.2.b and Fig.2.c show that the oscillations appear really scattered between
 77 the horizontal and vertical radius. The evolution of both radius are shown in number of
 78 pixels as a function of the number of seconds since the pyrometer was started.
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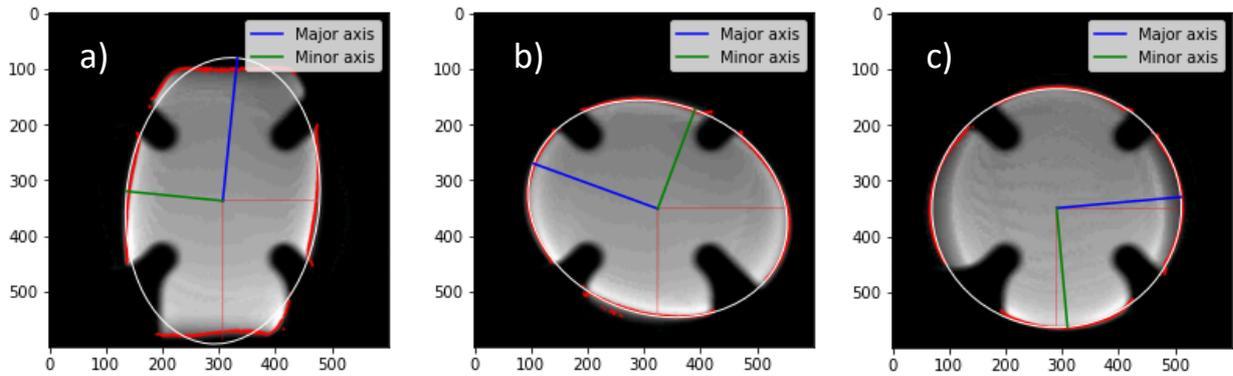


80
 81 *Figure 2. Cycle 14 a): Frame at t=3.08s after pulse. b) Oscillations of R_x . c) Oscillations of R_y*

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 83
 84

85 It was then chosen to fit an elliptical profile to the detected edges of the sample and
86 measure a and b , the minor and major axes of the sample, for each time frame. An example
87 of this process can be seen in Fig.3, with the elliptical profile shown in white.

88



89

90 *Figure 3. Cycle 14 a): Frame at t=0.00s after pulse. b) t= 1.45s after pulse. c) t=4.78s after pulse*

91

92 It is then possible to apply the oscillating drop method using four different radii,
93 R_x, R_y, a and b , to calculate the thermophysical properties. The effective radius R_{eff} is then
94 introduced as $R_{eff} = \sqrt{ab}$ to minimize the scatter caused by the overlap between the
95 rotation and translation of the sample and its oscillations [6]. After Fast Fourier Analysis of
96 the oscillations, a high pass filter at 30 Hz was applied to eliminate the background noise and
97 reveal the relative damped oscillations as shown in Fig.4. The calculations of the
98 uncertainties for the three properties have been conducted according to the Evaluation of
99 Measurement Data-Guide to the Expression of Uncertainty in Measurement principal
100 handbook while assuming the sources of errors to be independent [7].

101

102 The surface of the sample was observed using a Tescan Vega scanning electron
103 microscope. Images of the microstructure were acquired with an acceleration voltage of 20
104 kV and a working distance of 10 mm using the backscatter electron detector.

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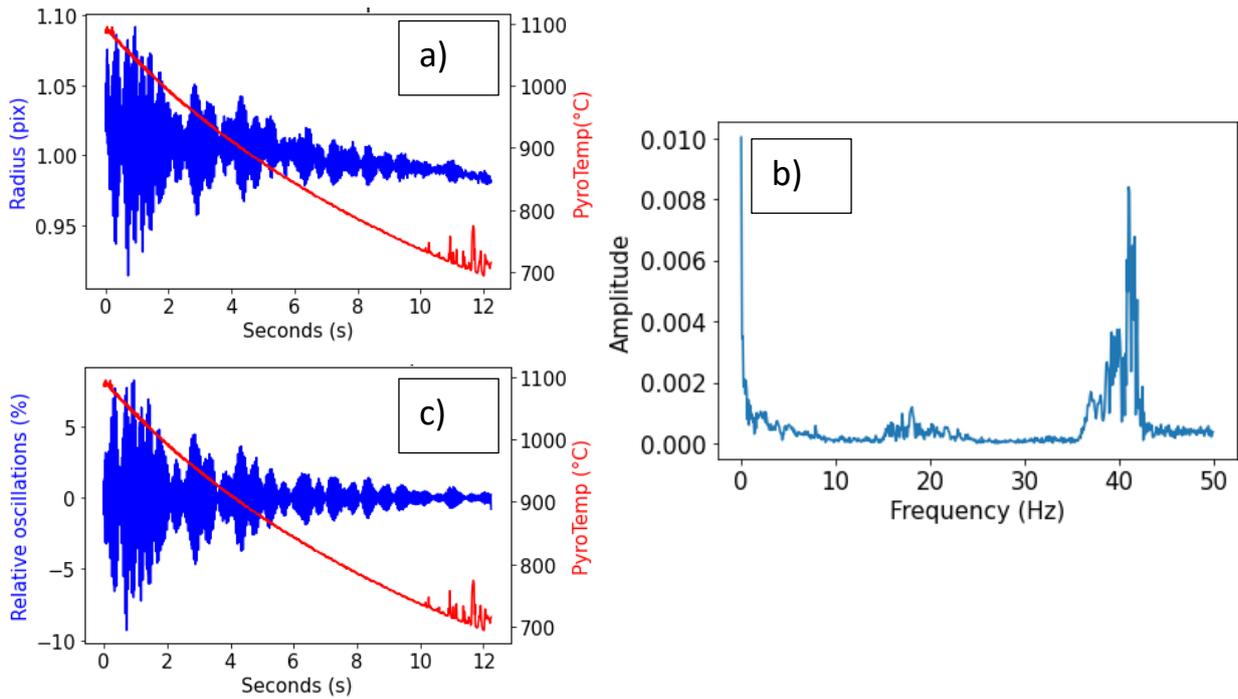
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112 Figure 4. Cycle 11 a) Before filtering b) Fast Fourier Transform of the oscillations c) After High-pass 30Hz filtering

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114 Results and Discussion

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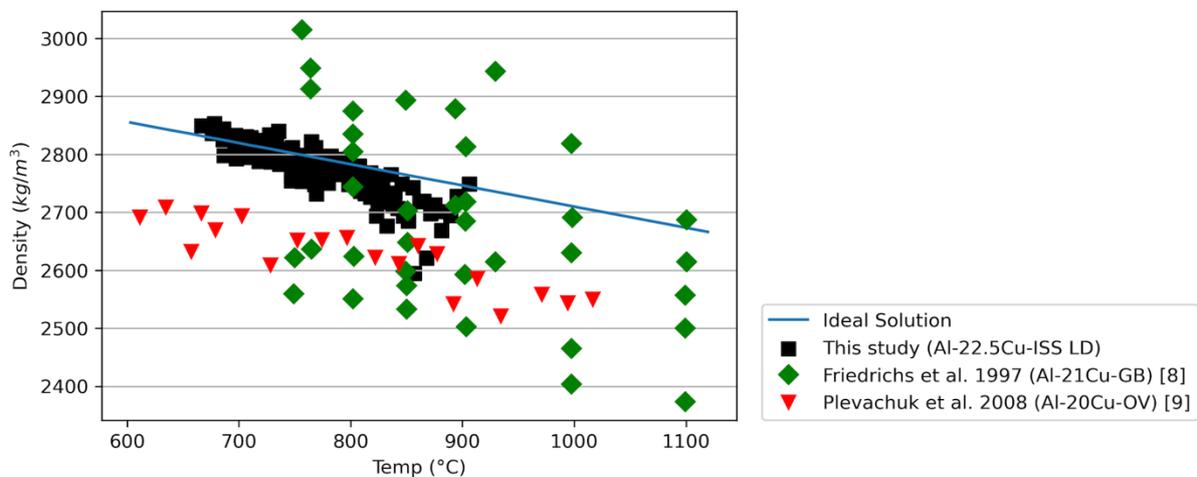
117 Density

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119 The density of the sample can be calculated as $\rho(T) = \frac{m}{V(T)}$ with m and V the mass and
 120 volume of the sample respectively. The volume is traditionally calculated by fitting Legendre
 121 polynomials to the detected edges of the sample to determine the radius profile $R(\theta)$ of the
 122 drop, with θ the angle from the horizontal axis. By assuming a vertical axis of symmetry for
 123 the rotation, the volume can then be calculated from $V = \frac{2}{3}\pi \int_0^\pi R(\theta)^3 \sin(\theta) d\theta$. However,
 124 the strong deformations and the coils hide part of the edges of the drop making the volume
 125 to be underestimated with this method. Instead, we approximated the volume of the sample
 126 by assuming its shape to be described by an ellipsoid symmetrical along its major axis, also
 127 called a spheroid, such as $V = \frac{4}{3}\pi a^2 b$. Other modes of oscillations can arise from the initial
 128 strong deformation leading to nonlinearities and invalidating the spheroidal shape
 129 assumption for the drop. The beginning of the oscillations is then trimmed so that the
 130 relative oscillations are under 5% to get more accurate measurements of the density in the
 131 domains where the oscillations are linear. The density measurements are averaged in

132 segments where the temperature drops by 10°C which leads to 10-15 measurements per
 133 cycle. In Fig.5, the results of this study for the density of Al-22.5wt%Cu are presented
 134 alongside with the Ideal solution model and experimental results obtained using the Gas
 135 Bubble Pressure method and the Oscillating Vessel method [8-9]. The results agree
 136 especially well with the ideal solution predictions. The increased scattering of the results
 137 obtained in this study at the highest temperatures is understood to be caused by the
 138 residual nonlinearities affecting the spheroidicity of the drop and the precision of the edge
 139 detection algorithms.

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142 *Figure 5. Density results*

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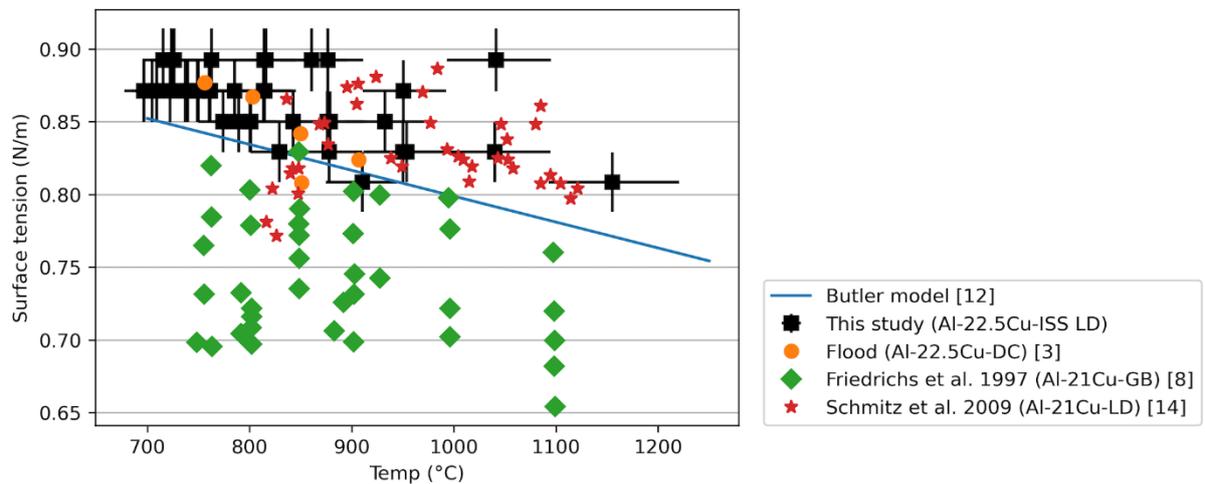
144 Surface tension

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146 During the oscillations, the shape of the drop can be described mathematically using
 147 spherical harmonics [10]. In microgravity the oscillations modes are degenerated and so
 148 assume the same frequency while the second mode of oscillation is the most stable and the
 149 dominating one. Under the linear assumption of small amplitude oscillations, Rayleigh
 150 derived the relationship between the surface tension of the fluid and the natural frequency
 151 of the oscillations f_n , and the mass of the sample such as $\sigma(T) = \frac{3\pi m f_n(T)^2}{8}$ [11]. The natural
 152 frequency can be determined as the maximum peak frequency of a Fast Fourier Transform
 153 (FFT) of the oscillation's amplitude. The uncertainty on the measurement of the natural
 154 frequency from the FFT is inversely proportional to the signal duration. Inversely, the
 155 uncertainty on the temperature is proportional to the signal duration as the temperature

156 drops during the oscillations. An optimal condition is found for a signal duration between
 157 two and three seconds. For the analysis, the oscillations of each cycle are then segmented
 158 into two-second-long portions. This allows for multiple surface tension measurements per
 159 cycle while the temperature decreases. In Fig.6, the results of this study for the surface
 160 tension are compared with the Butler model and experimental results obtained using the
 161 Gas Bubble Pressure method, the Levitated Drop method and the Discharge Crucible method
 162 [8, 12-14]. The measurements obtained in this study are consistent with the results
 163 predicted by the Butler model [12] and the experimental results obtained by P.Flood [13]
 164 and Shmitz et al. [14]. The limited resolution of 0.5 Hz of the FFT for two-second long
 165 segments of data recorded at an acquisition rate of 100 Hz is limiting the precision in the
 166 identification of the maximum peak frequency. Additionally, the rotation of the droplet and
 167 the nonlinear deformation may cause splitting of the resonance peak and frequency shift
 168 [15].

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171 *Figure 6. Results for surface tension*

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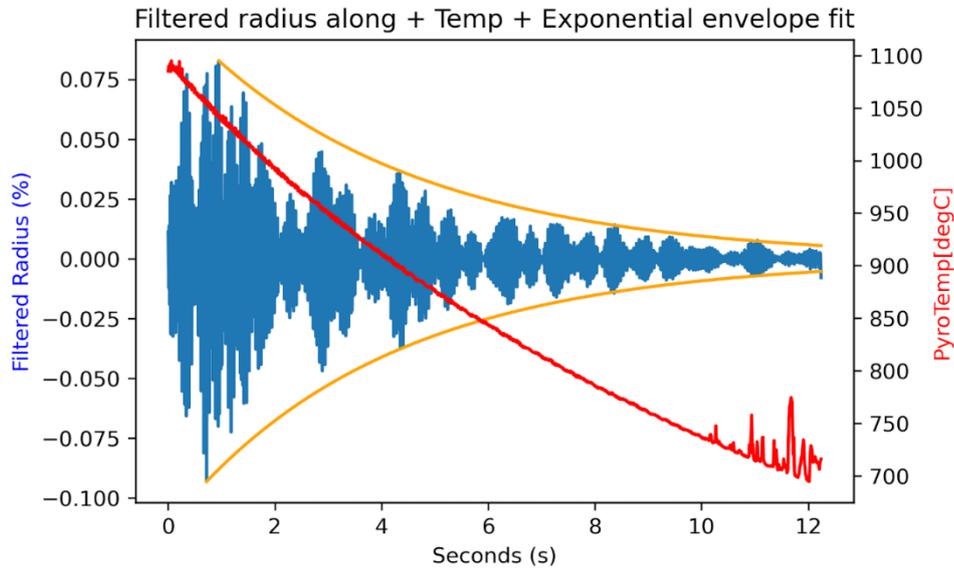
173 **Viscosity**

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175 The viscosity is related to the damping constant τ of the oscillations according to the Lamb
 176 equation such as $\eta(T) = \frac{3m}{20\pi R_o(T)\tau(T)}$ [16]. The effective radius oscillations seem modulated
 177 suggesting strong nonlinearities at play as shown in Fig.3 and Fig.4. As a result, significant
 178 information on the evolution of the peaks of the amplitude is missing which prevents from
 179 accurately fitting an exponential decay envelope to segments of the signal to determine the

180 decay constant and thus, the viscosity. To minimize the uncertainty, this fitting is then done
181 on selected peaks of the whole data set from the initial deformation as opposed to the
182 segmented analysis for the surface tension. Fig.7 show a result of this fitting executed on the
183 lower and upper envelope.

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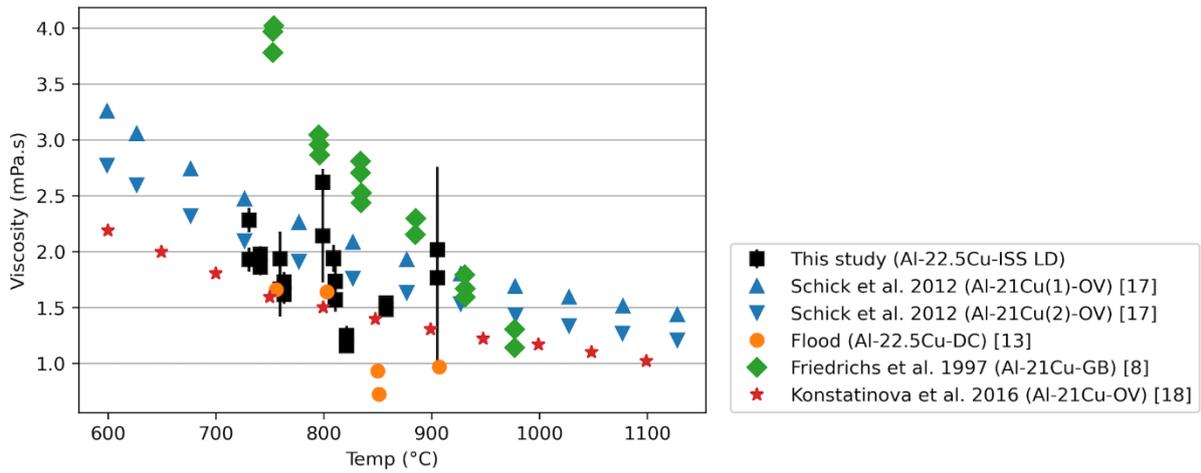


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186 *Figure 7. Cycle 11) Exponential fitting of the envelope of the relative oscillations of the effective radius.*

187

188 In Fig.8, the results of this study for the viscosity of Al-22.5wt%Cu are compared with
189 experimental results obtained using the Gas Bubble Pressure method, the Oscillating Vessel
190 method, and the Discharge Crucible method [8,13,17-18]. The measurements are agreeing
191 well with the experimental results by P.Flood [13], Schick et al. [17] and Konstatinova [18].
192 The uncertainties on the measurements are understood to be due to the high initial
193 deformation giving rise to nonlinear oscillations. This phenomenon coupled with the sample
194 rotation and translation modulate the perceived damped oscillations from the side camera
195 point of view and affect the precision of the Lamb equation in predicting the viscosity of the
196 sample. More research is needed to consider the nonlinearities arising from initial high
197 deformation and leading to the temporary presence of other modes of oscillations to correct
198 the lamb equation and modify the modelling of the damped oscillations [19].



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200 *Figure 8. Viscosity results*

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202 **Microstructure**

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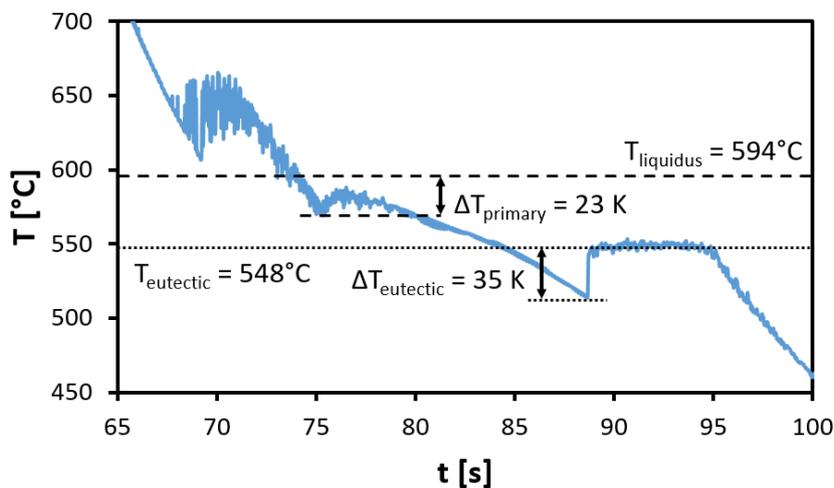
204 Figure 9 shows the temperature-time profile of the last solidification cycle performed on the

205 Al-22.5wt%Cu sample. Two recalescence events are observed. The first one corresponds to

206 the primary nucleation of α -Al, with an undercooling of $\Delta T_{\text{primary}} = 23$ K. The second event is

207 the nucleation of the α -Al-Al₂Cu eutectic with an undercooling of $\Delta T_{\text{eutectic}} = 35$ K.

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210 *Figure 9. Cycle 21) Temperature-time profile of the last solidification cycle.*

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212 Due to the uniqueness and rarity of this sample, it is important to perform a complete

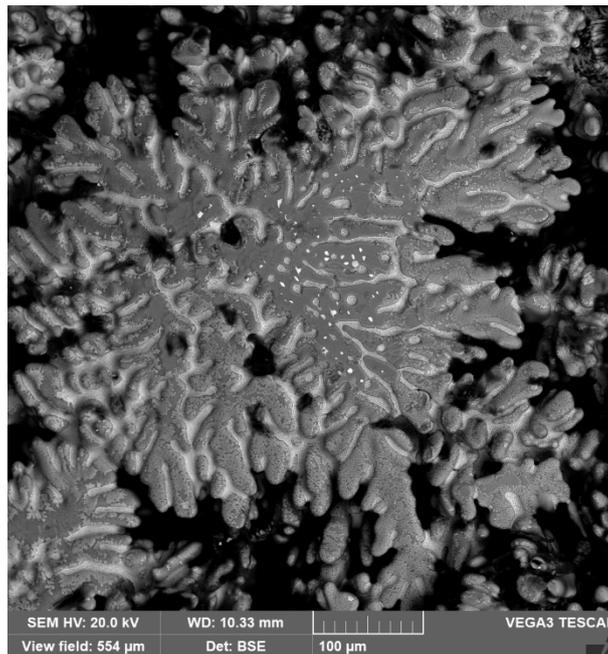
213 characterization, which starts with non-destructive techniques. At the time of writing, the

214 sample was at Los Alamos National Laboratory to undergo neutron diffraction, which will be

215 followed by synchrotron X-ray tomography at the Canadian Light Source. Microstructure

216 analysis is thus far limited to the surface of the sample. Figure 10 shows an SEM micrograph
217 of the region that solidified last. A large amount of porosity due to solidification shrinkage
218 can be seen. The dendrites present are well developed. However, their morphology is
219 reminiscent of seaweed. They do not exhibit the 4-fold symmetry typical of the $\langle 100 \rangle$ of α -Al.
220 None of the primary dendrites visible at the surface of the sample seem to grow along $\langle 100 \rangle$.
221 Deviation from $\langle 100 \rangle$ growth direction in Al-Cu alloys has been previously reported. α -Al
222 dendrites have been shown to grow along $\langle 111 \rangle$ directions in impulse atomized Al-4.5wt%Cu
223 [20-21]. This has been attributed to the anisotropy of attachment kinetics. However, the
224 cooling rates and undercoolings involved in atomization are much higher than what was
225 experienced by the Al-22.5wt%Cu sample in the MSL-EML. This deviation is perplexing as
226 both Al and Cu have interfacial energy anisotropies that favour $\langle 100 \rangle$ growth. This will be
227 investigated further with X-ray tomography and the following full metallographic study of
228 the sample.

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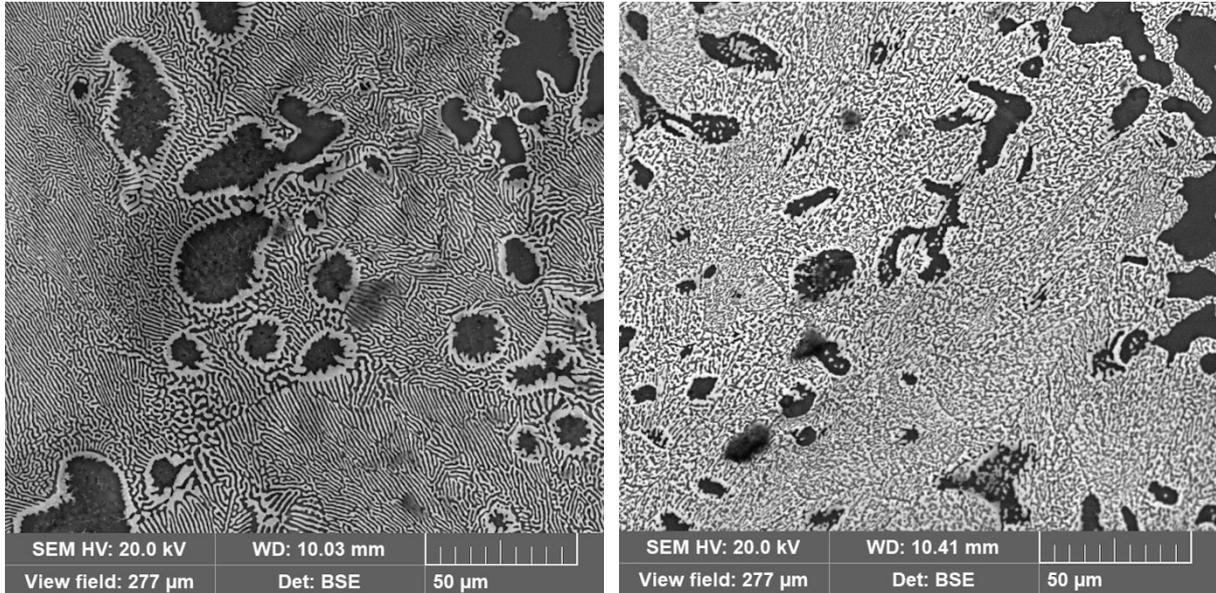
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231 *Figure 10. SEM micrograph of the end of solidification.*

232 Figure 11 shows a close-up of the microstructure on the opposite side of Figure 8. As
233 expected, it is composed of primary α -Al surrounded by an α -Al- Al_2Cu eutectic. However,
234 two distinct morphologies are observed. Figure 8 left displays a typical lamellar eutectic,
235 while the eutectic in Figure 8 right can be defined as undulated. Such a morphology has been
236 observed in rapidly solidified eutectic Al-33wt%Cu droplets [22]. It is assumed that the

237 undulated eutectic grows rapidly during recalescence, while the lamellar eutectic grows
238 post-recalescence during the eutectic plateau. Since the eutectic undercooling is fairly small
239 (35 K), the amount of undulated eutectic should be much less than that of the lamellar. This
240 will be further explored during the metallographic analysis of the sample.

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242

243 *Figure 11. Two different eutectic morphologies observed at the surface of the Al-22.5wt%Cu sample: regular lamellar (left)*
244 *and undulated (right).*

245

246 Conclusions

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248 Overall, the measurements obtained for the density, surface tension and viscosity of the Al-
249 22.5wt%Cu sample with the Oscillating Drop Method applied on data obtained from the ISS-
250 EML are consistent with results in the literature. However, the low weight of the sample
251 made it difficult to constrain the initial deformation to a linear regime. The only available
252 images were taken from the side camera and were showing a nonlinear behavior for the
253 damped oscillations coupled with rotation and translation of the sample. This affected the
254 precision of the method for the calculation of the three thermophysical properties. The
255 shape of the sample was likely non spheroidal at the beginning due to the presence of higher
256 modes of oscillations than the assumed second mode. This nonlinear regime was avoided for
257 density calculation as it would lead to uncertainties. The rotation and nonlinearities also lead
258 to peak splitting and frequency shifting in the FFT contributing to the uncertainty of the
259 surface tension measurement. Finally, the viscosity measurement of the sample was the

260 most affected by the nonlinear behavior. The oscillations were appearing modulated leading
261 to an important loss of information on their amplitude and making a segmented exponential
262 fitting of the envelope non accurate. The decay constant was then obtained through fitting
263 of the envelope on selected peaks for the whole data available for each cycle to minimize
264 the uncertainties and the Lamb equation was then applied to calculate the viscosity. Further
265 research is needed to develop corrections on the traditional Rayleigh and Lamb equation
266 [9,14] and compensate for the information loss on the perceived damped oscillations due to
267 the rotation, translation and nonlinearities.

268 Preliminary investigation of the surface of the sample showed unexpected features. The
269 growth direction of the visible dendrites seemed to deviate from the characteristic $\langle 100 \rangle$
270 expected of primary α -Al. Furthermore, two distinct eutectic morphologies were observed: a
271 typical lamellar eutectic, as well as an undulated structure, akin to what is formed in rapidly
272 solidified eutectic Al-33wt%Cu droplets. A full microstructural study will be carried out after
273 the sample has undergone a full set of non-destructive analysis.

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