

**Clusters, Clouds, and Correlations:  
Relating Young Clusters to Giant Molecular Clouds in M33 and M31**

by

Joshua Caleb William Peltonen

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Department of Physics  
University of Alberta

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# Abstract

Using young clusters and giant molecular clouds (GMCs), we constrain timescales and distances involved in the star formation process in the Triangulum (M33) and Andromeda (M31) galaxies. The young star clusters come from the PHATTER catalogue, with 1214 clusters in M33 and the PHAT catalogue, with 1249 clusters in M31. These cluster catalogues have the ages of their members measured using colour-magnitude diagram fitting. In M33, we compare these clusters to 444 GMCs identified from an ALMA  $^{12}\text{CO}(2-1)$  survey with 35 pc resolution. In M31, we compare the clusters to 251 GMCs identified from a CARMA  $^{12}\text{CO}(1-0)$  survey with 20 pc resolution. By comparing the separations between clusters and their nearest GMC, we find that younger clusters are preferentially located near GMCs, while older clusters have separations consistent with a random distribution. Performing two-point correlation on our clusters and GMCs reveals that younger clusters are likely to be found at small separations from other young clusters. In M33, where the cluster catalogue contains younger clusters, we find through cross-correlation analysis that GMCs and clusters  $\leq 10$  Myr are statistically associated. This association between GMCs and clusters decreases until the clusters are  $\approx 18$  Myr, where the clusters and GMCs then appear randomly distributed for older clusters. This transition to a more random distribution can be explained by a simple drift model where the clusters have a velocity dispersion of 5-10  $\text{km s}^{-1}$  with respect to the surrounding gas. By finding the number of spatially overlapping young clusters and GMCs in M33, we estimate that clusters spend 4-6 Myr emitting feedback into their parent GMC. We then estimate that GMCs in M33 have a total lifetime of 11-15 Myr.

# Preface

Chapter 2 of this thesis has been published as Peltonen J., et al., 2023, “Clusters, clouds, and correlations: relating young clusters to giant molecular clouds in M33 and M31,” *Monthly Notices of the Royal Astronomical Society*, vol. 522, issue 4, 6137–6149. I was responsible for the data analysis as well as the manuscript composition. E. Rosolowsky was the supervisory author and was involved with concept formation and manuscript composition. All other authors contributed to manuscript edits.

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# Abbreviations

**ACA** Atacama Compact Array.

**ALMA** Atacama Large Millimeter/submillimeter Array.

**CARMA** Combined Array for Research in Millimeter-wave Astronomy.

**CMD** colour-magnitude diagram.

**CO** Carbon monoxide.

**GMC** Giant molecular cloud.

**HST** Hubble Space Telescope.

**IQR** Interquartile range.

**MAST** Mikulski Archive for Space Telescope.

**PHAT** Panchromatic Hubble Andromeda Treasury.

**PHATTER** Panchromatic Hubble Andromeda Treasury: Triangulum Extended Region.

**SCIMES** Spectral Clustering for Molecular Emission Segmentation.

**SFR** Star formation rate.

**UV** Ultraviolet.

**YSCC** Young star cluster candidate.

# Chapter 1

## Introduction

### 1.1 Motivation

The Universe is filled with galaxies, which are teeming with stars. These stars produce a majority of the optical light emitted from galaxies, which results in galaxies being observed primarily by their starlight. However, we know that galaxies began with gas and dust without stars. Therefore, understanding how this gas transformed into stars and the distribution of these stars, which defines galactic structure, is intimately linked with the star formation process. Understanding this star formation process that converts gas into stars is essential to understand galaxy evolution as it creates the stars we observe and shapes the gas that will birth the future generations of stars (Mo et al., 2010).

#### 1.1.1 Giant Molecular Clouds

The star formation process is facilitated by gravitational collapse. First, the hot, low-density gas that filled the early Universe and much of present galaxies must coalesce. As the gas condenses, it can also cool more efficiently, which reduces outward thermal pressure and allows collapse to continue. Many cooling processes require other particles. The simplest example is collisional excitation, where some of the kinetic energy of one particle is transferred to another through a collision putting this particle in an excited state. Then when this particle naturally decays to the

ground state, the emitted photon has a chance of escaping the region, which will lower the total kinetic energy of the particles. Another example is Bremsstrahlung radiation, where two charged particles pass near each other. This passage can cause one of the particles to decelerate, which releases photons and again reduces the total kinetic energy of the region. Both of these cooling processes happen more frequently when the density is higher, and are why cooling is more efficient in higher-density regions. As the gas that makes up the interstellar medium (ISM) gets denser and colder, it transitions from ionized to neutral atomic, and then into molecules (Kwok, 2007). Star formation occurs primarily in the densest regions of the ISM, which are collections of molecular gas known as molecular clouds. In particular, giant molecular clouds (GMCs) host a majority of observed star formation (e.g., Wong & Blitz, 2002; Kennicutt et al., 2007; Leroy et al., 2008; Bigiel et al., 2008).

Compared to most of the gas in galaxies, GMCs are extremely cold and dense. Molecular gas in the ISM is very sensitive to the ultraviolet (UV) radiation produced by high-mass stars. This sensitivity is why, even though gravity wants to increase the density of the gas, most of the gas in the ISM is in the neutral atomic state. Therefore, for GMCs to exist, they require many layers that shield the molecules from UV radiation. This multi-layering leads to reservoirs of molecular gas that are 10 – 100 pc across, have masses  $10^3 - 10^6 M_{\odot}$ , and are surrounded by atomic gas (Krumholz, 2015). Similar to the mass distribution of stars, there seem to be more low-mass than high-mass GMCs. However, the exact GMC mass function is difficult to obtain due to incomplete samples of lower mass GMCs (Roman-Duval et al., 2010). While these GMCs are cold, they are also turbulent with typical gas velocity dispersions of a few  $\text{km s}^{-1}$  (e.g., Rosolowsky et al., 2003; Gratier et al., 2014). Larson (1981) was the first to note that many GMC properties are related. Now known as Larson’s laws which are one, GMCs have typical surface densities of  $\sim 100 M_{\odot} \text{ pc}^{-2}$ . Two, the velocity dispersion is proportional to the radius of the GMC. Three, GMCs are roughly in virial equilibrium or  $\alpha_{\text{vir}} \approx 1$ , Section 1.1.2 will

define this quantity in more detail and Section 1.1.4 will discuss how these GMCs are detected and how some of these properties are measured.

### 1.1.2 From Gas to Clusters

Star formation occurs in the densest,  $100 - 1000 \text{ cm}^{-3}$ , and coldest,  $\sim 10 \text{ K}$ , regions of GMCs. This is because the main force driving star formation is gravitational collapse. As gravity tries to collapse the gas, thermal pressure, magnetic fields, and turbulence resist the collapse. The exact effects of magnetic fields and turbulence on star formation are difficult to quantify (Lada & Lada, 2003). However, just looking at the simplest case where there is only gravitational energy,  $\mathcal{W}$ , and thermal energy,  $\mathcal{T}$ , we can see how increasing density leads to collapse following a simplified derivation from Krumholz (2015). With only these two forces, the virial theorem can be written as

$$\frac{1}{2}\ddot{I} = 2\mathcal{T} + \mathcal{W}, \quad (1.1)$$

where  $\ddot{I}$  is the double time derivative of the moment of inertia, which determines if the cloud is expanding (positive value) or contracting (negative value). Therefore, the virial parameter  $\alpha_{\text{vir}}$ , where

$$\alpha_{\text{vir}} = \frac{2\mathcal{T}}{|\mathcal{W}|}, \quad (1.2)$$

is often used as a simple way of defining whether a system is bound by gravity. For a cloud with mass  $M$  and radius  $R$  that is governed only by thermal pressure and gravity

$$\mathcal{T} = \frac{3}{2}Mc_s^2 \quad (1.3)$$

and

$$\mathcal{W} = -a\frac{GM^2}{R}, \quad (1.4)$$

where  $c_s$  is the sound speed and  $a$  is a constant based on the density profile. Since a negative  $\ddot{I}$  means collapse, the condition for collapse is

$$\frac{GM^2}{R} \gtrsim Mc_s^2, \quad (1.5)$$



and with some manipulation and dropping factors of order unity, we get

$$R \gtrsim \sqrt{\frac{c_s^2}{G\rho}}. \quad (1.6)$$

This equation defines the scale of perturbations that would lead to collapse. Therefore, for increasing density, smaller perturbations can lead to collapse. Once the gas is dense enough that random perturbations from turbulence or outside forces lead to collapse, stars begin to form from this collapsing gas.

Only some of the gas that makes up GMCs will collapse into stars. Most of the gas that makes up GMCs is eventually dispersed back into the interstellar medium (Chevance et al., 2022a). By surveying large populations of GMCs and measuring the star formation rate, one can obtain  $\tau_d$ , the depletion time. The depletion time is the time it would take for all of the molecular gas in GMCs to be converted into stars and is on the order of a Gyr (Leroy et al., 2008). Individual stars form more quickly than this; on the order of  $\tau_{ff}$ , the free-fall time given by

$$\tau_{ff} = \sqrt{\frac{3\pi}{32G\rho}}. \quad (1.7)$$

The free-fall time determines the time it would take for a cloud to collapse with negligible resistance. Efficiency is found by comparing these two timescales

$$\epsilon_{ff} = \frac{\tau_{ff}}{\tau_d}. \quad (1.8)$$

Both simulations (e.g., Semenov et al., 2016; Grisdale et al., 2019; Kim et al., 2021b) and observations (e.g., Evans et al., 2014; Leroy et al., 2017; Schrubba et al., 2019) have converged on efficiencies of  $\approx 0.01$ , which means only a tiny fraction of the gas is converted into stars per free-fall time. The stars created by GMCs are responsible for dispersing the gas through stellar feedback and supernovae (Kim et al., 2018; Chevance et al., 2022b).

Since star formation occurs in limited environments, stars typically form very near to other stars. It is also important to note that fragmentation is common. When a

large mass of gas begins to collapse, the density will increase, which leads to perturbations on smaller scales being unstable to collapse (Equation 1.6), and the large-scale cloud is thus fragmented into smaller clouds. Therefore, 70 – 90% of stars form in clusters (Lada & Lada, 2003). These clusters can then begin dispersing the sensitive molecular gas, returning it to the ISM as atomic gas. By the time the gas is dispersed, only  $\sim 10\%$  of stars remain in gravitationally bound clusters, where Lada & Lada (2003) define a bound cluster as a cluster with  $\alpha_{\text{vir}} < 1$  or total energy that is negative. Most clusters lose members and eventually dissolve through outside and internal interactions (Krumholz et al., 2019). Therefore, the amount of stellar mass contained in bound clusters, known as the cluster formation efficiency  $\Gamma$ , changes with stellar age. However, the reason most clusters are quickly destroyed is not fully understood. One leading theory is that in the densest regions of GMCs, there is a local increase in efficiency leading to a bound cluster. In contrast, clusters formed in lower-density areas are primarily held together by the surrounding gas, so when the gas is cleared away, the cluster is no longer bound (Kruijssen, 2012; Grudić et al., 2021).

The clusters that remain gravitationally bound are important laboratories. The stars that make up clusters are formed nearly simultaneously, under the same conditions, and in the same location in a galaxy, which means their properties can be analyzed with minimal complications. With knowledge of stellar evolution, these simple stellar populations can be age dated. Over their lifetime, stars of different masses trace different paths on colour-magnitude diagrams. Therefore, the members of clusters trace unique distributions on colour-magnitude diagrams based on their age and mass. These unique distributions allow the age and mass of clusters with resolved members to be measured accurately (Krumholz, 2015). This makes clusters one of the most direct measurements of star formation history in a galaxy (Mo et al., 2010).

Clusters inherit many properties from the gas they are formed in. GMCs are

confined tightly to the planes of galaxies. Therefore, young clusters begin very close to the planes of galaxies and are distributed in the same way as GMCs. However, clusters also inherit the turbulent gas's random velocities, allowing them to drift. The turbulent velocities in a GMC are on the order  $10 \text{ km s}^{-1}$  (Rosolowsky et al., 2003), and this measurement is made using the methods discussed in Section 1.1.4. To see evidence of how clusters have inherited velocities Soubiran et al. (2018) found that young clusters in the Milky Way have similar velocity dispersions based on measurements of proper motion and radial velocity from *Gaia*. This random drift means that clusters move toward randomly distributed as they get older (Soubiran et al., 2018). Therefore, young clusters should be the closest to where they were created.

To understand how the stellar populations and the gas in galaxies evolve, it is essential to constrain the timescales involved in these processes. Some of the important timescales that characterize galaxy evolution are the lifetime of the GMCs, the cluster feedback time, and the timescales over which the cluster distributions evolve. The lifetime of the GMCs are vital for understanding the cycling of gas in galaxies and will add constraints on the physics involved in the star formation process. The feedback time (i.e., the measure of time it takes for clusters to clear away their natal gas) is required to understand why most clusters are disrupted by the end of this time. The rates at which the clusters integrate into the larger stellar distribution of the galaxy will add to the understanding of how stellar populations evolve.

### 1.1.3 Correlation Analysis

One way to begin characterizing these timescales is through spatial correlation analysis. Over time the spatial structure of clusters and clouds will change. This change is due to the random velocities of the clusters. Therefore, by using the cluster ages and analyzing the spatial structure, many timescales can be determined. The simplest way of analyzing spatial correlation is finding the typical separations between

objects. However, a more formal way to characterize spatial correlation is using the two-point correlation function, also known as the autocorrelation function. The two-point correlation function finds the excess probability,  $\xi(r)$ , that a distribution is more clustered than a random distribution at a scale  $r$ . Thus, the probability of finding an object at a distance  $r$  from another object in a volume  $dV$  is given by

$$dP = n[1 + \xi(r)]dV, \quad (1.9)$$

where  $n$  is the number density. For a distribution that is more clustered than random,  $1 + \xi(r) > 1$ , while  $1 + \xi(r) = 1$  and  $1 + \xi(r) < 1$  represent random and anticorrelated distributions, respectively. Typically the random distribution that is compared to is a Poisson distribution such as a radial exponential to represent the disk of a galaxy. However, in some cases, a different random distribution is used to represent a different structure. The two-point correlation function describes the clustering of a distribution made up of one type of object. To find the clustering between two types of objects, the cross-correlation function is used. The cross-correlation finds the probability that a distribution is clustered with another distribution compared to two random distributions (Peebles, 1980).

The two-point correlation function was first used in astronomy to understand how galaxies are distributed in the Universe. Galaxies are distributed in three-dimensional (3D) space. Therefore, the ideal correlation function would find the separations of galaxies,  $r$  in 3D space, which can then be compared to a random distribution (Peebles, 1980). Unfortunately, accurate distance information for galaxies is difficult to obtain. This difficulty has led to most measurements of the correlation of galaxies being estimations based on angular separations. The excess probability  $\omega(\theta)$  of finding objects separated by an angle  $\theta$  is

$$dP = N[1 + \omega(\theta)]d\Omega, \quad (1.10)$$

where  $N$  is the surface density and  $d\Omega$  is the solid angle. One popular estimator from Landy & Szalay (1993) finds the excess probability,  $\omega(\theta)$ , by comparing the number

of real object pairs to random object pairs in an angular distance bin  $\theta + d\theta$ . This estimator is given by

$$\omega(\theta) = \frac{1}{RR} \left[ DD \left( \frac{N_R}{N_D} \right)^2 - 2DR \left( \frac{N_R}{N_D} \right) + RR \right], \quad (1.11)$$

where  $N_D$  and  $N_R$  are the total numbers of real and random objects, and  $DD$ ,  $RR$ , and  $DR$  are the number of real, random, and real-random pairs in the bin.

More recently, various forms of spatial correlation analysis have been applied to clusters and GMCs in nearby galaxies. For example, in the Antennae galaxies, Zhang et al. (2001) used correlation analysis on a small number of clusters and the molecular gas. Zhang et al. (2001) found through the two-point correlation function that young clusters exist in a clustered distribution and through the cross-correlation function that young clusters are linked with the molecular gas. In the Large Magellanic Cloud (LMC), Kawamura et al. (2009) used the spatial overlap of clusters with GMCs to show that as clusters get older, the typical separation between clusters and GMCs increases. These two examples were among the first to find that young clusters correlate with each other and GMCs, but this correlation decreases with time. This correlation is expected to decrease because, as noted above, clusters have random motions and, therefore, will drift toward a random distribution over time.

With improvements in observations, more work has been done with larger numbers of resolved clusters and GMCs in galaxies more distant than the LMC. Grasha et al. (2018, 2019) used the two-point correlation function to understand the distribution of large numbers of clusters in NGC 7793 and M51, respectively. They found that young clusters  $\leq 10$  Myr are correlated with each other, while clusters  $> 10$  Myr are consistent with a random distribution. Turner et al. (2022) expanded this methodology to 11 nearby galaxies, finding similar results using the two-point correlation function. They also used the cross-correlation function to find that young clusters are correlated with GMCs, but older clusters are not. These initial results are promising, but they suffer from looking at a limited number of distant galaxies where star clusters and

GMCs are poorly resolved. By studying nearer galaxies, we have the opportunity to use correlation functions as a probe of GMC and cluster evolution.

### 1.1.4 Observational Methodology

**Giant Molecular Clouds** – Measuring the molecular gas is an essential aspect of understanding the star formation process. Ideally, we want to measure  $\text{H}_2$ , the most abundant molecule in the ISM. Unfortunately, the energy levels of  $\text{H}_2$  are quite far apart, with the lowest excited state at  $\approx 500$  K above the ground state. Therefore,  $\text{H}_2$  is not readily excited in the cold temperatures found in GMCs which makes  $\text{H}_2$  difficult to detect. However, CO, the second most abundant molecule, is readily excited in the conditions of GMCs. These excitations make CO much easier to detect, and it is the most common tracer of molecular gas in the ISM (Krumholz, 2015). The lowest energy transitions for molecules are rotational transitions indicated by a J. The most frequent CO transitions J=1-0 ( $\approx 115$  GHz), and J=2-1 ( $\approx 231$  GHz) are observable by some of the most powerful radio telescopes in operation (e.g., the Atacama Large Millimeter/submillimeter Array, ALMA).

Many of the properties of GMCs can be derived from observations of CO. The first step to deriving the properties of GMCs is to identify the boundaries of the GMCs. While the boundaries of GMCs are not always clear (Chevance et al., 2022a), many algorithms attempt identifying individual clouds from molecular gas emission. Examples of these algorithms are the CLUMPFIND algorithm (Williams et al., 1994), which identifies clouds using contours of constant flux, and the SCIMES (Colombo et al., 2015) algorithm, which identifies clouds using graph theory. Once individual GMCs are identified, the total amount of CO emission can be used to find the mass. However, since most of the mass is actually made up of  $\text{H}_2$  a conversion between CO brightness and  $\text{H}_2$  needs to be made, which is a function of metallicity and has been the subject of much research (e.g., Arimoto et al., 1996; Leroy et al., 2011; Bolatto et al., 2013). The typical conversion is done using  $X_{\text{CO}}$  with typical values in the

Milky Way of  $\approx 2 \times 10^{20} \text{ cm}^{-2} (\text{K km s}^{-1})^{-1}$ . Accurso et al. (2017) find that the conversion factor is related to metallicity by

$$X_{\text{CO}} \propto Z^{1.6}, \quad (1.12)$$

where  $Z$  is the metallicity. From the width of the CO emission line, the velocity dispersion of the gas can be found. This is due to the Doppler effect, as particles emitting photons at different speeds will emit slightly shifted frequencies (Kwok, 2007).

**Star Clusters** – Clusters are most easily identified in the optical. In nearby galaxies, clusters have primarily been identified using the *Hubble Space Telescope (HST)*. It is difficult to automate the process of identifying clusters in the crowded discs of galaxies. Therefore, most catalogues of clusters are identified by eye (e.g., Whitmore et al., 2021). One way to identify clusters in large optical surveys is to employ the power of citizen scientists. For example, Johnson et al. (2015) and Johnson et al. (2022) recruited members of the public to identify clusters in PHAT (M31) and PHATTER (M33) surveys. Then DOLPHOT is used to identify individual cluster members. DOLPHOT fits the unique point spread function of each filter to the images to extract photometry for each star. Then the age and the mass of the cluster can be found by plotting the photometry of the cluster members on a CMD and comparing with simulated clusters (Wainer et al., 2022).

## 1.2 Thesis Objectives

The galaxies M31 and M33 are ideal targets for extending work on the relationship between clusters and GMCs. M31 and M33 are the nearest true spiral galaxies, which allows for studies of star formation at excellent resolution while avoiding the problems present in Milky Way studies. These problems include distance ambiguity, line-of-sight blending, and very large surveys to get representative samples. M31 and M33

are both at a single known distance, so all clusters and GMCs within them are also at a single distance. M33 is especially desirable because its low inclination and smaller size allow for clear lines of sight and extensive surveys. M31's inclination and size make surveys more difficult compared to M33, but still more manageable than the Milky Way.

With the recent *HST* and molecular gas surveys in M31 and M33, we can analyze the distribution of clusters and GMCs at unprecedented resolutions. In M31, the high-resolution surveys of CO were obtained by the Combined Array for Research in Millimeter-wave Astronomy (CARMA) (Caldú-Primo & Schruba, 2016; Leroy et al., 2016; Schruba et al., 2019) along with clusters from the PHAT survey (Johnson et al., 2015). In M33, we have an ALMA survey of CO (Koch et al. in prep) and clusters from the PHATTER survey (Johnson et al., 2022). These cluster surveys have sufficient resolution to identify individual cluster members, yielding extremely accurate ages and masses. These accurate ages allow us to find the two-point correlation between the clusters and the cross-correlation between clusters and GMCs, in terms of age. Analyzing the correlations over time enables the characterization of many of the timescales involved in the star formation process.

### 1.3 Thesis Outline

The following chapter of this thesis will describe how I use the high-resolution surveys of GMCs and clusters to constrain several timescales in the star formation process. These timescales are found using the spatial correlations of these clusters and GMCs first in M33 and then in a more limited capacity in M31. Specifically, I begin by finding the separations between the clusters and GMCs. Then I measure the two-point correlation of the clusters. Finally, I calculate the cross-correlation between the clusters and GMCs. All of these correlations are made using a variety of similarly aged cluster samples across a wide range of cluster ages, which is then used to determine when the clusters and GMCs are no longer statistically correlated, the feedback



timescale, and the lifetime of the GMCs in M33. I then develop a simple drift model to find the drift velocity that erases the correlation over the correct timescale.

The final chapter of this thesis provides some concluding remarks and points to future work.

# Chapter 2

## Paper 1

This is a copy of the paper, published in *Monthly Notices of the Royal Astronomical Society*, vol. 522, issue 4, 6137–6149, which has been adapted for formatting, but otherwise remains unchanged.

### Abstract

We use young clusters and giant molecular clouds (GMCs) in the galaxies M33 and M31 to constrain temporal and spatial scales in the star formation process. In M33, we compare the Panchromatic Hubble Andromeda Treasury: Triangulum Extended Region (PHATTER) catalogue of 1214 clusters with ages measured via colour–magnitude diagram (CMD) fitting to 444 GMCs identified from a new 35 pc resolution Atacama Large Millimeter/submillimeter Array (ALMA)  $^{12}\text{CO}(2-1)$  survey. In M31, we compare the Panchromatic Hubble Andromeda Treasury (PHAT) catalogue of 1249 clusters to 251 GMCs measured from a Combined Array for Research in Millimeter-wave Astronomy (CARMA)  $^{12}\text{CO}(1-0)$  survey with 20 pc resolution. Through two-point correlation analysis, we find that young clusters have a high probability of being near other young clusters, but correlation between GMCs is suppressed by the cloud identification algorithm. By comparing the positions, we find that younger clusters are closer to GMCs than older clusters. Through cross-correlation analysis of the M33 cluster data, we find that clusters are statistically

associated when they are  $\leq 10$  Myr old. Utilizing the high precision ages of the clusters, we find that clusters older than  $\approx 18$  Myr are uncorrelated with the molecular interstellar medium (ISM). Using the spatial coincidence of the youngest clusters and GMCs in M33, we estimate that clusters spend  $\approx 4-6$  Myr inside their parent GMC. Through similar analysis, we find that the GMCs in M33 have a total lifetime of  $\approx 11-15$  Myr. We also develop a drift model and show that the above correlations can be explained if the clusters in M33 have a  $5-10 \text{ km s}^{-1}$  velocity dispersion relative to the molecular ISM.

## 2.1 Introduction

The essential heavy elements around us are evidence of the cycling process between gas and stars in the Universe. Molecular gas in the interstellar medium (ISM) forms high-mass stars and clusters that, in their evolution, will disrupt the surrounding gas. This star formation and resulting disruption rely on the interplay of gravity, turbulence, stellar feedback, magnetic fields, chemistry, and thermal regulation. The processes that play the most significant role in regulating star formation remain areas of active research (McKee & Ostriker, 2007; Chevance et al., 2022a).

Constraining the time-scales associated with the various phases of the star formation process can help constrain which physical processes are at play. This has motivated a number of studies that attempt to infer the evolutionary time-scales of the molecular ISM, in particular, lifetimes of giant molecular clouds (GMCs) or the time-scale over which feedback destroys the GMCs. GMCs are massive collections of molecular gas that are the primary sites of star formation. Early attempts to measure the lifetimes of GMCs led to a range of results with some estimates converging to long lifetimes ( $10^{7.5} - 10^8$  yr; Bash et al., 1977; Scoville & Hersch, 1979) and others to shorter lifetimes ( $< 10^{7.5}$  yr; Blitz & Shu, 1980; Blitz, 1993). Lifetimes that are  $10^{7.5}$  yr or longer require GMCs to survive much of a galactic rotation period with forces that prevent gravitational collapse. Some recent analyses of gas distribution still point to

long GMC lifetimes (Koda et al., 2012, 2016, 2020). However, there are also many approaches that point to short lifetimes. Kawamura et al. (2009) used the spatial coincidence of GMCs and young clusters to determine that GMCs live for 20-30 Myr in the Large Magellanic Cloud (LMC). Numerical simulations have also suggested a short molecular cloud lifetime, for example, 4-25 Myr in Dobbs & Pringle (2013) and 13-20 Myr in Jeffreson et al. (2021). Using the “tuning fork” measurements of the decorrelation between star formation tracers and molecular gas developed in Schruba et al. (2010, see also Onodera et al. 2010), Kruijssen & Longmore (2014) and Kruijssen et al. (2018) developed their “uncertainty principle” formalism to measure a short (10 Myr in NGC 300 Kruijssen et al., 2019) lifetime for molecular clouds. This methodology has now been replicated in simulations (Semenov et al., 2021) and applied to broader samples of galaxies (Chevance et al., 2022b; Kim et al., 2021a, 2022). Chevance et al. (2022b) argued for short lifetimes and contended that GMCs are dispersed quickly after the onset of star formation. They suggest that pre-supernova feedback mechanisms play a key role in disrupting clouds.

While the details of some of the measurements can be quite sophisticated, these measurements are usually framed around our naïve model for star formation in molecular clouds. The model starts with overdense turbulent complexes of molecular gas that can be divided into individual GMCs. However, the boundaries between GMCs and the outside ISM are not always clear (Chevance et al., 2022a). These GMCs then begin forming clusters that inherit the clustered structure of their progenitor clouds (Grasha et al., 2018, 2019; Turner et al., 2022). After the cluster spends some time inside their progenitor clouds they will disrupt the cloud through a combination of supernovae and stellar feedback (the significance of each effect is still debated; Kim et al., 2018; Chevance et al., 2022b). We define the lifetime of a GMC as the time between when the cloud can first be detected and when it can no longer be detected because of this disruption. The disruption of the dense gas occurs before most of the gas can be converted into stars, leading to very long depletion times and inefficient

star formation (Chevance et al., 2022a). After the clusters have dispersed the gas, the correlated structure of the clusters will be erased by random drift velocities inherited from the turbulent motions of the gas.

Many of the approaches to characterize the evolutionary time-scale of the molecular ISM rely on tracers of star formation ( $H\alpha$ , ultraviolet, and mid-infrared) that do not directly measure the ages of a stellar population. Instead, they trace the integrated radiation from a stellar population in a given waveband; for example,  $H\alpha$  traces the ionizing photon radiation that comes from short-lived O-stars. Using an assumed model of the initial mass function, these measurements can then be translated back into star formation rates and characterize the time-scale for evolution of molecular clouds, providing a narrow window into the star formation history (Kennicutt & Evans, 2012). These tracers are typically restricted to only a portion of the star formation process.  $H\alpha$ , for example, requires high-mass stars and is thus restricted to star formation events that host massive star formation.

Alternatively, the ages of simple stellar populations provide a robust method to establish evolutionary time-scales and provide a long view into the star formation history of a galaxy. Stellar clusters represent close approximations of simple stellar populations, which Kawamura et al. (2009) leveraged for their measurement of cloud evolution. Such studies cannot be carried out in the Milky Way for large samples because of line-of-sight blending and extinction. Therefore, comparing clusters and GMCs has been limited only to the nearest galaxies (e.g. the LMC) or limited to unresolved cluster candidates in galaxies like M33 (Corbelli et al., 2017). However, with recent *Hubble Space Telescope* (*HST*) surveys and interferometer observations, resolved clusters can be compared to GMCs in nearby galaxies like is done in Grasha et al. (2018, 2019) and Turner et al. (2022).

In this work, we take advantage of new, high-quality data to measure the relationship between cluster and molecular cloud populations in the two largest star-forming galaxies in the Local Group: M31 and M33. Clusters were identified in these galax-

ies using the wide-area *HST* surveys, Panchromatic Hubble Andromeda Treasury (PHAT; Dalcanton et al., 2012) in M31 and Panchromatic Hubble Andromeda Treasury: Triangulum Extended Region (PHATTER; Williams et al., 2021) in M33. In addition, high-resolution surveys of  $^{12}\text{CO}$ , obtained by the Combined Array for Research in Millimeter-wave Astronomy (CARMA) in M31 (Caldú-Primo & Schruba, 2016; Leroy et al., 2016; Schruba et al., 2019) and the Atacama Large Millimeter-/submillimeter Array (ALMA) in M33 (Koch et al. in prep), allow for the analysis of GMCs at sub-cloud scales. M31 is about three times the diameter (Vaucouleurs et al., 1991) and  $23^\circ$  more inclined (Koch et al., 2018; van der Marel et al., 2019) than M33. In M33, the surveys thoroughly cover the central part of the galaxy and a majority of the star-forming spiral arms (Williams et al., 2021). Because of its inclination and large area on the sky, the *HST* surveys in M31 are more limited and only cover a quadrant of the galaxy. Despite the limitations in M31, these high-resolution surveys and the relative proximity of M31 and M33 allow for colour–magnitude diagram (CMD) fitting of the clusters, which yields robust age and mass estimates. These accurate ages and the relatively deep mass completeness limits allow for an unprecedented study of molecular gas and star formation.

We focus on the correlation structure between the GMCs identified in the  $^{12}\text{CO}$  surveys and the clusters identified in the *HST* surveys. Previous studies have shown that very young clusters are typically near GMCs (Kawamura et al., 2009; Whitmore et al., 2014; Corbelli et al., 2017; Grasha et al., 2018, 2019). It has also been shown that the two-point correlation is stronger for younger clusters than older clusters (Grasha et al., 2015, 2017; Menon et al., 2021). Finally, the cross-correlation function has been used to show that young clusters are more correlated with GMCs than old clusters (Turner et al., 2022). However, Li & Barmby (2021) have noted that applying correlation functions to non-homogeneous populations may produce spurious correlation signal. While the more sophisticated ways of interpreting correlation, like spatial-point processes, have their benefits, they are difficult to interpret

Table 2.1: Adopted Parameters for M33 and M31.

	M33	M31
Distance (kpc)	859 <sup>a</sup>	776 <sup>b</sup>
Orientation	$i=55^\circ$ <sup>c</sup> , PA=201 <sup>o</sup> <sup>c</sup>	$i=78^\circ$ <sup>d</sup> , PA=38 <sup>o</sup> <sup>d</sup>
Central position <sup>e</sup>	RA=23.46204 <sup>o</sup> DEC=30.66022 <sup>o</sup>	RA=10.68479 <sup>o</sup> DEC=41.26907 <sup>o</sup>
SFR <sup>e</sup> ( $M_\odot$ /yr)	0.32	0.39
Stellar Mass <sup>e</sup> ( $M_\odot$ )	$2.63 \times 10^9$	$5.37 \times 10^{10}$

<sup>a</sup>de Grijs et al. (2017); <sup>b</sup>Savino et al. (2022); <sup>c</sup>Koch et al. (2018); <sup>d</sup>van der Marel et al. (2019); <sup>e</sup>Leroy et al. (2019)

and implement. Therefore, we have decided to address some of the concerns with correlation functions by building in the overall effects of galactic structure by using random cluster distributions that contain the same non-homogeneous structure.

The details of the M33 surveys and the more limited M31 surveys are presented in Section 2.2. Because of the limitations of the M31 data, Section 2.3 focuses on the main results found in M33. These results include an analysis of the correlation between clusters and GMCs, a comparison of cluster and GMC properties, and estimations of GMC timescales. Section 2.4 is a partial parallel analysis on M31. In Section 2.5, we develop a simple model to estimate the drift velocity of the clusters and discuss the effects of completeness. Finally, a summary of the work is presented in Section 2.6.

## 2.2 Observations

### 2.2.1 M33

We trace GMCs in M33 using a new CO survey carried out using the ALMA Atacama Compact Array (ACA). The ALMA ACA survey covers a section of M33 shown in Figure 2.1 as a dark blue outline and is described in more detail in Koch et al. (in prep). The ACA survey is centred on the  $^{12}\text{CO}$  J=2-1 transition at 230.538 GHz

with a bandwidth of 154 MHz  $\approx$  200 km s<sup>-1</sup>. The synthesized beam size of  $\approx$  8''.5 corresponds to a physical size of  $\approx$  35 pc at the distance of M33 from Table 2.1. The data have a noise level of 45 mK in a 0.7 km s<sup>-1</sup> channel. Koch et al. (in prep) applied the Spectral Clustering for Molecular Emission Segmentation (SCIMES) algorithm (Colombo et al., 2015) to the ACA data to obtain a catalogue of 444 GMCs. These GMCs are shown in Figure 2.1 as blue diamonds. An integrated intensity map for one of these GMCs is shown in Figure 2.1 in the bottom zoomed-in frame.

The PHATTER survey (Williams et al., 2021) is composed of observations in six *HST* filters and covers a region of M33 shown in Figure 2.1 as a red outline. The top zoomed-in panel of Figure 2.1 shows a portion of the PHATTER survey (in the F475W filter for *HST*), illustrating the resolution is sufficient to identify individual stars and clusters. Using a crowd-sourced visual search, Johnson et al. (2022) identified 1214 clusters from the PHATTER survey that are believed to be long-lived. These clusters are marked in Figure 2.1 as orange circles. The ages and masses of the clusters were identified using CMD fitting as described in Wainer et al. (2022). This catalogue of optically identified clusters is distinct from the catalogue of 630 young star cluster candidates (YSCCs) in Corbelli et al. (2017) found using *Spitzer Space Telescope* 24  $\mu$ m images (Verley et al., 2007). These infrared YSCCs, originally identified by Sharma et al. (2011), are thought to be young ( $\leq$ 10 Myr) embedded objects. However, Johnson et al. (2022) found that at least 30% of the YSCCs cannot be embedded young clusters because of the lack of visual extinction. Because of the potential contamination of non-cluster objects, we will only use this YSCC catalogue from Corbelli et al. (2017) for comparison in Section 2.3.2 and 2.3.3.

### 2.2.2 M31

A portion of M31 was surveyed using CARMA, PI A. Schruba, which includes short spacing data from Nieten et al. (2006). These data have appeared in Caldú-Primo & Schruba (2016), Leroy et al. (2016), and Schruba et al. (2019) with description and



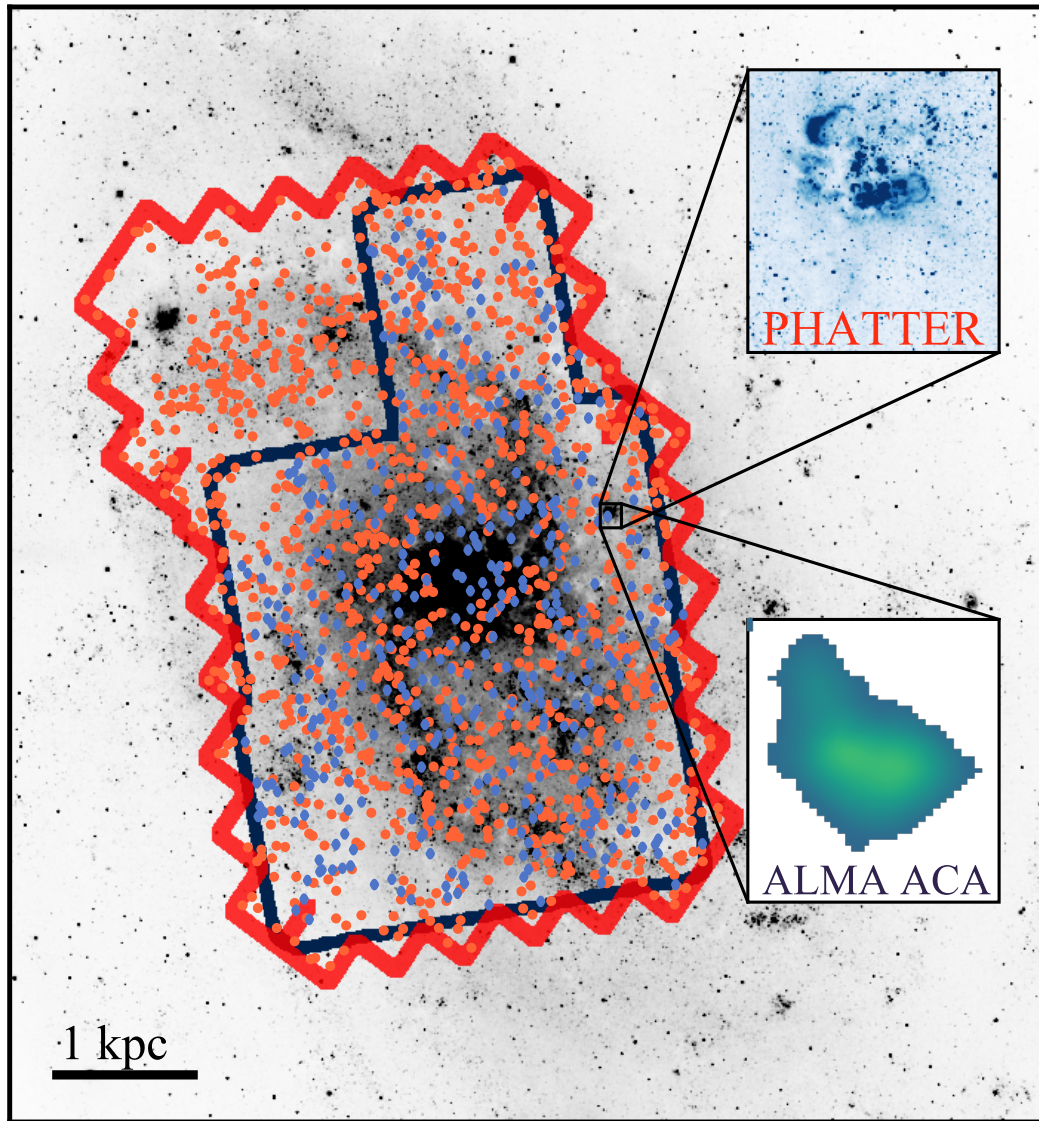


Figure 2.1: The data used to study the clusters and GMCs in M33. The background is a  $B$  band image from the 4-m Mayall telescope (Massey et al., 2006). The red outline shows the footprint of the PHATTER survey, while the dark blue outline shows the footprint of the ALMA ACA survey. The top zoomed-in frame shows the resolution of the PHATTER data. The bottom zoomed-in frame shows a GMC in the same area identified from the ALMA ACA data. The orange circles show the locations of the clusters identified from the PHATTER data. The GMCs identified from the ALMA ACA survey are shown as blue diamonds.

images of the data in those papers. Figure 2.2 shows the area surveyed in the CO survey. With a synthesized beam size of  $\approx 5''.5$  and at the distance of M31 from Table 2.1, the physical resolution is  $\approx 20$  pc. The data have a noise level of 190 mK in a  $2.5 \text{ km s}^{-1}$  channel. We do not match the CARMA resolution to the ALMA ACA resolution because we want to utilize the highest resolution possible, and most of our analysis will be done at scales larger than the resolution.

We then applied the SCIMES algorithm (Colombo et al., 2015) to the CARMA data, which yielded a catalogue of 251 GMCs. We use the same default algorithm parameters as were used in Koch et al. (in prep) for signal identification and cloud decomposition. These GMCs are shown in Figure 2.2 as blue diamonds. Figure 2.2 also shows an integrated intensity map of one of the clouds in the bottom zoomed-in frame.

We use the stellar cluster catalogue generated from the PHAT survey (Dalcanton et al., 2012), the predecessor to PHATTER. PHAT covers a quadrant of M31’s star-forming disc in six *HST* bands. Figure 2.2 shows an area that is smaller than the PHAT survey, with a small portion shown in the top zoomed-in panel. Using a crowd-sourced visual search, Johnson et al. (2015) identified 2753 clusters from the PHAT survey. Age and mass estimates are derived from CMD fitting, and a subsample of 1249 clusters with ages between 10-300 Myr were reported in Johnson et al. (2016). These clusters are shown in Figure 2.2 as orange circles.

### 2.2.3 Completeness

Throughout the analysis, we will use all of the clusters identified in Johnson et al. (2022) and Johnson et al. (2016) to use the greatest number of sources. However, our results will depend on the completeness properties of the cluster and GMC catalogues. For the clusters in both PHATTER and PHAT, the visual search for clusters included injected synthetic clusters of known mass and age (Johnson et al., 2015, 2022). These synthetic clusters allow for thorough completeness analysis. In M31, the cluster cat-

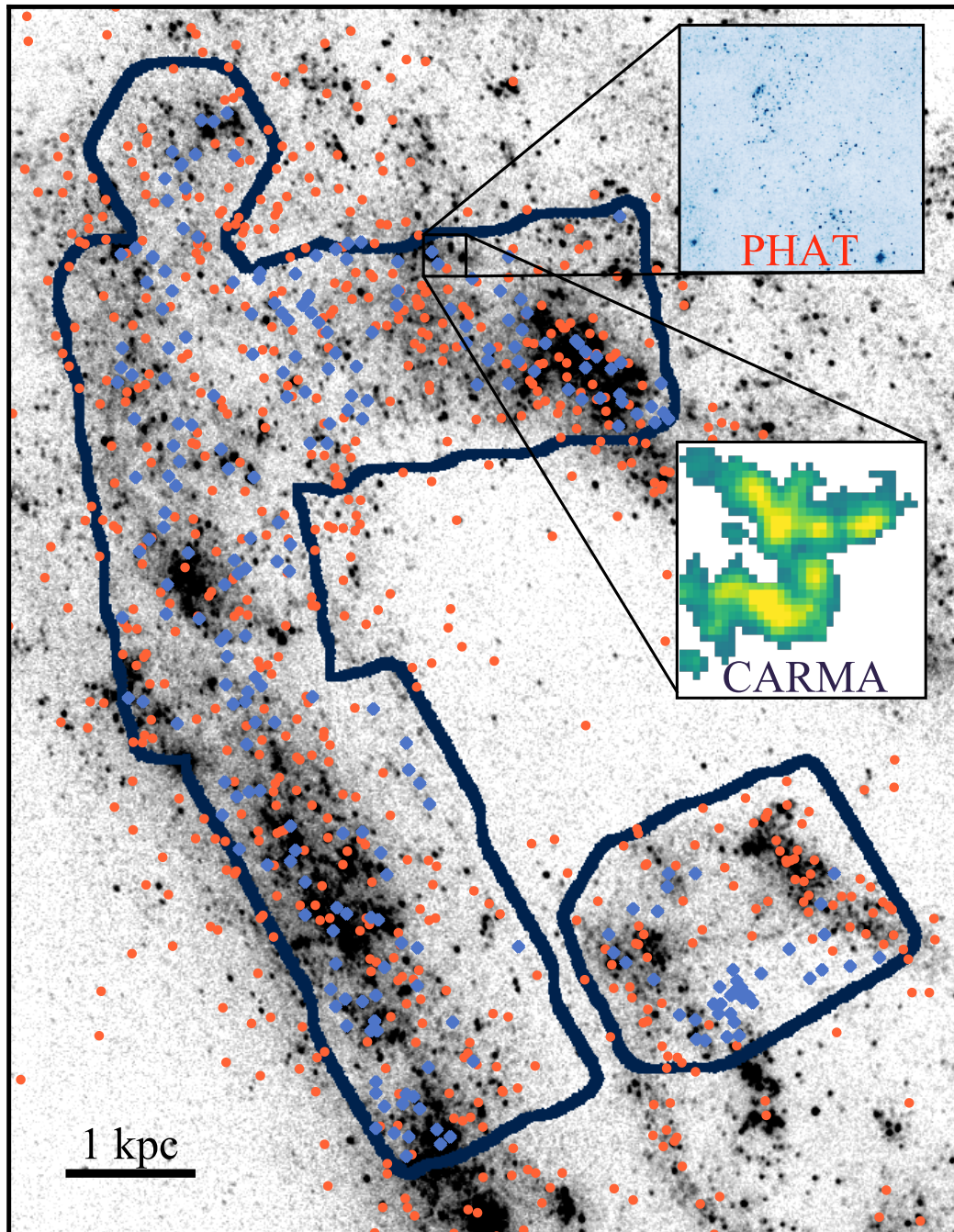


Figure 2.2: The data used to study the clusters and GMCs in M31. The background is a far-ultraviolet image from the *Galaxy Evolution Explorer* (Gil de Paz et al., 2007). The dark blue outline shows the footprint of the CARMA survey. The top zoomed-in frame shows the resolution of the PHAT data. The bottom zoomed-in frame shows a GMC in the same area identified from the CARMA data. The orange circles show the locations of the clusters identified from the PHAT data. The GMCs identified from the CARMA survey are shown as blue diamonds.

alogue at 100 Myr is 50% complete down to  $M_{c,50} \approx 500 M_{\odot}$  (Johnson et al., 2015). The completeness is better for younger clusters since they are more likely to have bright young stars. However, the oldest clusters in the catalogue (100-333 Myr) are 50% complete down to  $M_{c,50} \approx 1000 M_{\odot}$  (Johnson et al., 2015). The completeness in M33’s cluster catalogue is worse than in M31 with a 50% completeness limit of  $\approx 1000 M_{\odot}$  at an age of 100 Myr (Johnson et al., 2022). This difference in completeness in M31 and M33 comes from the higher degree of crowding in M33, which makes identifying clusters more difficult (Johnson et al., 2022). Therefore, more crowded regions of each galaxy will also lead to lower completeness in these regions. Another property that has an impact on cluster completeness is extinction. The very youngest (<3 Myr) deeply embedded clusters will be difficult to identify due to optical extinction (Johnson et al., 2022). The very young embedded stars should be visible with recent observations from the *James Webb Space Telescope* (JWST), which can see with eyes unclouded by the optical extinction of GMCs (Peltonen et al., in prep).

For M33 and M31, we measure the 50% completeness limit for molecular cloud identification as  $M_{\text{GMC},50} = 1.3 \times 10^4 M_{\odot}$  and  $M_{\text{GMC},50} = 3.0 \times 10^4 M_{\odot}$  respectively, based on the artificial cloud recovery test method presented in Rosolowsky et al. (2021). This approach inserts GMCs of known brightness and properties into the data and tests whether they are recovered in the cloud identification algorithm. This value is approximately  $40\times$  the  $1\sigma$  noise level in a single synthesized beam. Our estimate assumes a Galactic CO-to- $\text{H}_2$  conversion factor (Bolatto et al., 2013) and a CO(2-1)/CO(1-0) line ratio of  $R = 0.7$  (Leroy et al., 2022).

Another issue to consider is not the absolute mass limit of the cloud and cluster catalogues independently, but possible mismatches between these two limits. In other words, could small clusters be undetectable when hosted by the lowest mass of the GMC sample, or conversely, could the host GMC of a low-mass stellar cluster be undetectable in the CO observations? The mass of the GMC completeness limit is 15 to 30 times larger than the corresponding cluster mass completeness limit (typically

$10^3 M_\odot$  as above), suggesting these two catalogue limits are well-matched as long as the efficiency of a GMC forming a cluster (by mass) is  $\sim 0.03$ - $0.06$ . If the true efficiency is lower (i.e., a given GMC can only host a smaller maximum-mass cluster), then some catalogued GMCs may host clusters that couldn't actually be detected, and the cluster catalogue would be incomplete with respect to the GMC catalogue. If the efficiency is higher, then the converse is true and some low mass stellar clusters may actually live in GMCs that are too low mass to have been detected, and the GMC catalogue would be incomplete with respect to the stellar cluster catalogue.

We can compare the “matched” catalogue of efficiency of  $0.03$ - $0.06$  to current constraints on the efficiency of stellar clusters forming in GMCs. Estimates suggest that, on the scale of individual clouds, up to 30% of the cloud mass may be converted into stars (Krumholz et al., 2019). This is a much higher efficiency, suggesting that the lowest mass stellar clusters in our samples may in fact be hosted by undetected GMCs, but that every catalogued GMC that hosts a stellar cluster should have the cluster detected. In other words, the cluster catalogue is complete with respect to the GMC catalogue, but the GMCs are not complete with respect to the cluster catalogue. We can also look at global, ensemble estimates of the efficiency, by multiplying the efficiency of turning molecular gas into stars on large (kpc) scales ( $f_\star$ ) by the fraction of the overall star formation that produces stellar clusters ( $f_{\text{clust}}$ ). Current estimates of these quantities are  $f_\star \sim 0.03$  (Utomo et al., 2018; Chevance et al., 2022a; Kim et al., 2022) and  $f_{\text{clust}} \sim 0.1$  (Krumholz et al., 2019), giving an overall efficiency of  $f_\star f_{\text{clust}} \sim 0.003$ . This is much lower than our ideal “matched” catalogue efficiency, which would imply that some GMCs in our catalogue may actually host undetectable stellar clusters, but that every stellar cluster should have its host GMC detected. In this case, the strength of our cross-correlation signals would be lower limits since the presence of more clusters correctly matched near progenitor cloud structures should enhance the cross-correlation amplitude.

We can also assess whether there is additional low mass incompleteness in the

stellar cluster catalogue by constraining the population of embedded clusters using the 24  $\mu\text{m}$ -derived YSCC catalogue from Corbelli et al. (2017) to assess whether there are infrared sources without associated CO emission, which would imply our survey misses clouds hosting embedded cluster formation. We find that  $188/244 = 77\%$  of sources in the survey area overlap with CO clouds. Since up to 30% of these infrared sources could be interlopers (Johnson et al., 2022), the degree of spatial coincidence suggests our cloud catalogue is sufficiently deep to include most of the cluster-forming cloud population.

Finally, we note that both cloud catalogues rely on CO as a tracer of molecular gas and our mass estimates above rely on assumptions about the CO-to- $\text{H}_2$  conversion factor in these galaxies. Moreover, these two studies use two different line tracers ( $J = 2 - 1$  for M33 and  $J = 1 - 0$  for M31), so the catalogues inherit potential biases from relying on CO emission as a proxy for star forming gas. Our study, however, relies primarily on the locations of the CO emission and not on its brightness, so concerns about conversion factor variations (Leroy et al., 2011) will manifest as uncertainties in the mass completeness limit. Similarly, systematic variations in the CO line ratios with galaxy and environment (e.g. Koda et al., 2020; den Brok et al., 2021; Leroy et al., 2022) may represent a caution for comparing the M33 and M31 results. Our analysis relies on whether a given cloud is detected or not, so such concerns will manifest as changes in the true completeness limits relative to those determined by false source injection of CO sources. Given the magnitude of these variations observed in the literature, we expect that the typical change in completeness limit would be about 0.2 dex, though these variations are typically measured in more massive systems than M33. This effect is relatively small compared to the range of cloud masses probed ( $\sim 2$  dex) so is not likely to dramatically change the results.

Despite careful measures of completeness for both the clouds and clusters, the lack of constraints on the mass fraction of stars found in bound clusters ( $f_{\text{clust}}$ ) and the relationship between cloud mass and cluster mass (Krumholz et al., 2019) precludes

a clear answer to how tightly the two populations relate to each other. We proceed assuming the populations are comparable and our later results do not contradict this assumption.

## 2.3 The Cloud-Cluster Population in M33

Utilizing the high-quality data of clusters (PHATTER) and GMCs (ALMA ACA) in M33, we determine how these clusters and GMCs are correlated. First, we find the separations between the clusters and clouds and determine if they correlate with the cluster’s age, as would be expected in the standard model of cloud and cluster evolution. We then refine this measurement with a two-point correlation analysis of the individual objects and a cross-correlation analysis of the clusters and GMCs. This cross-correlation analysis indicates the degree to which clusters are statistically associated with GMCs. We then look at how the properties of the clusters are related to the properties of the associated GMC, with a particular focus on the angle at which the clusters leave their associated clouds. Finally, we determine the lifetime of the GMCs in M33 by using the method of Kawamura et al. (2009), which is based on the spatial overlap of young clusters. These results depend on the completeness limits of the contributing catalogues, which we discuss in more detail in Section 2.2.3.

### 2.3.1 Cloud-Cluster Spatial Offsets

To compare the properties of the GMCs and clusters, we must first find their locations in the plane of M33’s disc. As seen in Figure 2.1 the clusters (from PHATTER) and the GMCs (from ALMA ACA) cover different areas. Therefore, we only consider GMCs and clusters in the overlapping survey areas, resulting in a sample of 444 GMCs and 934 of 1214 clusters. Here we treat the SCIMES molecular clouds and PHATTER clusters as point sources. We use the centres of the clusters as the point source location. For the molecular clouds, we use the location of the brightest CO emission (CO peak) as the location of the point sources. Using the orientation parameters

from Table 2.1, we convert celestial coordinates into galactocentric coordinates for each object and measure distances in the plane of the galaxy.

We then compare the azimuthally averaged radial distributions and generate random distributions that match the radial distributions of the different objects. We use these random distributions to assess the significance of our results. Figure 2.3 shows the radial distribution for the GMCs and the clusters split into three age categories. We choose three age categories that have distinct relations to the GMCs. The 60 youngest clusters have ages  $\leq 10^7$  yr, the 93 medium-aged clusters have ages between  $10^7$  yr and  $10^{7.5}$  yr, and the 781 oldest clusters are  $> 10^{7.5}$  yr old.

We assume that GMCs and clusters both follow an exponential distribution for surface density, so that the number in a given radial bin is:

$$N \sim 2\pi R \Delta R \exp\left(-\frac{R}{R_d}\right), \quad (2.1)$$

where  $R$  is the radial distance from the galactic centre and  $\Delta R$  is the width of the radial bin. The different cluster age bins and the GMCs have different scale lengths,  $R_d$ , that can be seen from the shapes of each radial distribution. We then generate random exponential distributions, selecting the same number of sources (934) in the same overlapping survey area. We generate 100 of these random distributions at each scale length for a range of scale lengths. Then we compared these random distributions to our real distributions and found the real-random distribution pair with the lowest chi-square value. We find a best-fitting scale length of 1.6 kpc for the youngest clusters, 3.2 kpc for the medium-aged clusters, 5.8 kpc for the oldest clusters, and 2.5 kpc for the GMCs. These scale lengths differ from M33's molecular gas scale length of  $\approx 2.1$  kpc (Druard et al., 2014) and stellar scale length of  $\approx 1.55$  kpc (Verley et al., 2009). The averaged random distribution fitted to the oldest clusters (5.8 kpc) is shown in Figure 2.3 as a red line. This random distribution based on the oldest clusters is used as our standard reference, but we use the other distributions for our correlation analysis.



There are many possible explanations for the differing scale lengths. The young clusters have a scale length consistent with the stellar scale length (Verley et al., 2009), and the GMCs have a scale length reasonably close to M33’s molecular gas scale length (Druard et al., 2014). However, the medium-aged and oldest clusters have longer scale lengths than the stellar scale length indicating fewer clusters at smaller radii, which is clear from Figure 2.3. Therefore, there might be something preventing the older clusters from being identified in the central region of M33. One possible explanation that is discussed in more detail in Section 2.2.3 is that clusters with young bright stars are easier to visually identify in the crowded central region (Wainer et al., 2022). However, this trend has also been observed in the Milky Way (Soubiran et al., 2018) and in simulations (Roškar et al., 2008). Therefore, the more likely explanation is due to clusters being destroyed in the crowded central region and from clusters migrating to larger radii.

We expect that the youngest clusters will be closer to their parent GMCs (e.g. Kawamura et al., 2009; Corbelli et al., 2017; Grasha et al., 2018, 2019), which can be shown by comparing cluster age and cluster-cloud separation. We find the closest GMC to each cluster using the galactocentric coordinates. We then measure the physical separation between the centre of each cluster and the nearest molecular cloud CO peak. Figure 2.4 shows the result of creating a box plot with the separation versus cluster age. Clusters of all ages have typical separations far greater than the 35 pc resolution of the ALMA ACA survey. The youngest clusters have the lowest median separation of 90 pc and the smallest interquartile range (IQR) of 60 pc. With a median separation of 100 pc, the medium-aged clusters are further from GMCs and have an IQR of 80 pc, larger than the youngest clusters. The oldest clusters comprised of three bins in Figure 2.4 are all quite similar, with medians of 120 pc and IQRs of 100-120 pc consistent with random. The random median (120 pc) and IQR (120 pc) shown in Figure 2.4 is from the oldest random distribution ( $R_d=5.8$  kpc). Using the other random distributions shifts the median and IQR by  $\approx 10$  pc, which

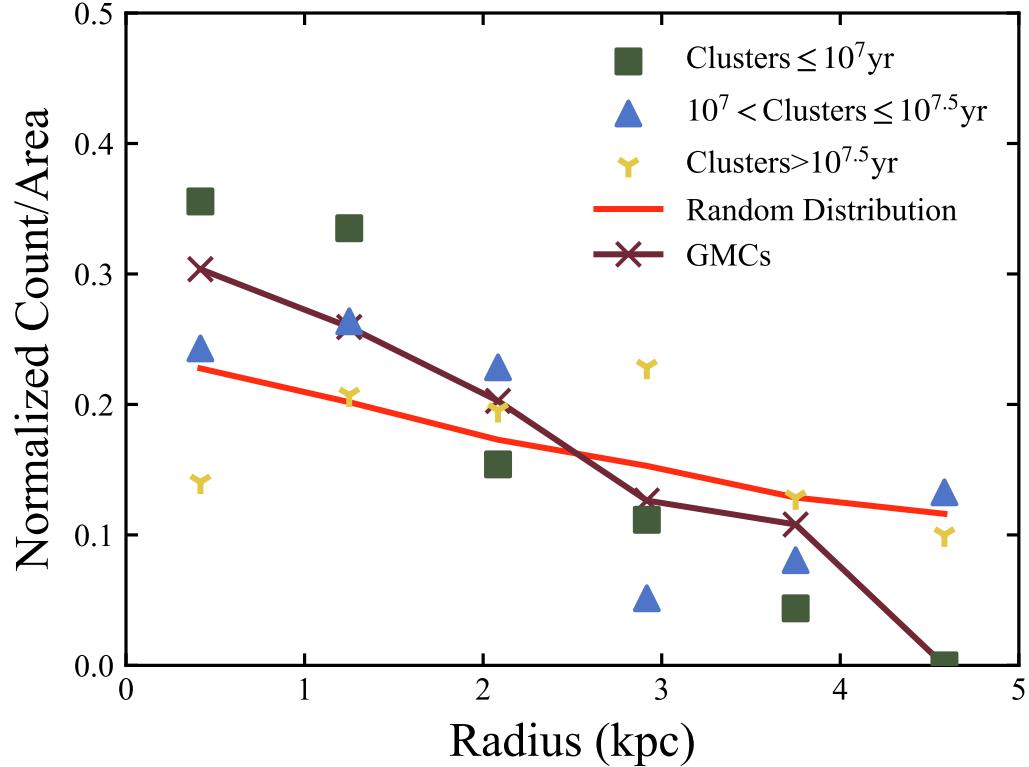


Figure 2.3: The normalized count per area of the clusters, GMCs, and the generated random clusters. The count is found in galactocentric radial bins and then divided by the area contained in the overlapping survey region in that radial bin. The clusters are split into the three age categories youngest (green squares), medium-aged (blue triangles), and oldest (yellow tri-points). The GMC radial distribution (dark pink line with crosses) is found from the coordinates of the CO peak of each cloud. The random cluster distribution (red line) results from averaging 100 exponential distributions with a scale length fitted from the oldest clusters. The random distribution traces a similar distribution to the oldest clusters in M33.

is still most consistent with the oldest clusters. As expected, the clusters start close to a GMC and drift towards randomly distributed as they age.

### 2.3.2 The Two-Point Correlation Function

We now further analyze the spatial properties of our distributions using the two-point correlation function that quantifies the amount of clustering on different spatial scales (Peebles, 1980). In general, the two-point correlation function describes the probability of finding an object in two volume elements separated by  $r$ . The two-point correlation of a real data set can then be compared to the two-point correlation of a random distribution to find the excess probability of spatial correlation (Peebles, 1980). We use a slight variation on the two-point correlation function since we are working in the plane of the disc. Therefore, we use  $\omega(r)$  to indicate the excess probability that an object will be found at a distance  $r$  from another object of the same type compared to a random distribution.  $1 + \omega = 1$  indicates a random uncorrelated distribution,  $1 + \omega > 1$  indicates a correlated distribution, and  $1 + \omega < 1$  indicates an anti-correlated distribution.  $\omega(r)$  is calculated in radial bins separated by  $\log_{10}(r) = 0.2$ . For each radial bin, we count the number of real catalogue pairs  $DD$ , random catalogue pairs  $RR$ , and pairs of one real and one random object  $DR$ . Then we use the Landy & Szalay (1993) estimator in the same form as Turner et al. (2022):

$$\omega(r) = \frac{1}{RR} \left[ DD \left( \frac{N_R}{N_D} \right)^2 - 2DR \left( \frac{N_R}{N_D} \right) + RR \right], \quad (2.2)$$

where  $N_R$  is the total number of objects in the random catalogue, and  $N_D$  is the total number of objects in the real catalogue. This process is repeated for each object type using their respective 100 random exponential distributions. Figure 2.5 shows the average two-point correlation function for the GMCs and the three cluster age categories. The standard deviation in two-point correlations from the random distributions is shown in Figure 2.5 as error bars.

The youngest and medium-aged clusters are correlated at separations smaller than

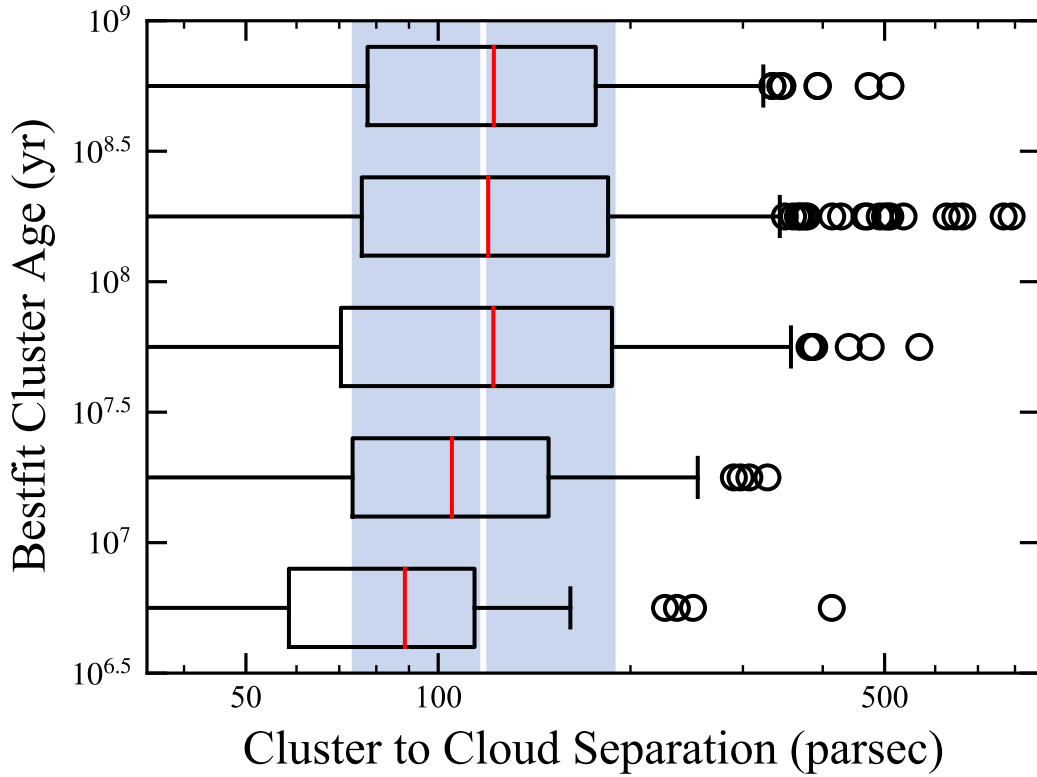


Figure 2.4: The separation between the clusters and the GMCs based on the age of the clusters in M33. The boxes show the span of the first and third quartiles (interquartile range or IQR) for each age bin that are  $10^{0.5}$  yr wide, with the medians marked with red lines. The error bars extending from the boxes indicate the minimum and maximum values in each bin, excluding the outliers. The outliers, marked with circles, are defined as points outside of 1.5 times the IQR. The separation between clusters with random positions and the GMCs have their median (white line), and IQR (blue shaded region) plotted. The youngest clusters have a shorter median and a smaller IQR than random clusters. Clusters older than  $10^{7.5}$  yr have medians and IQRs consistent with the random distribution.

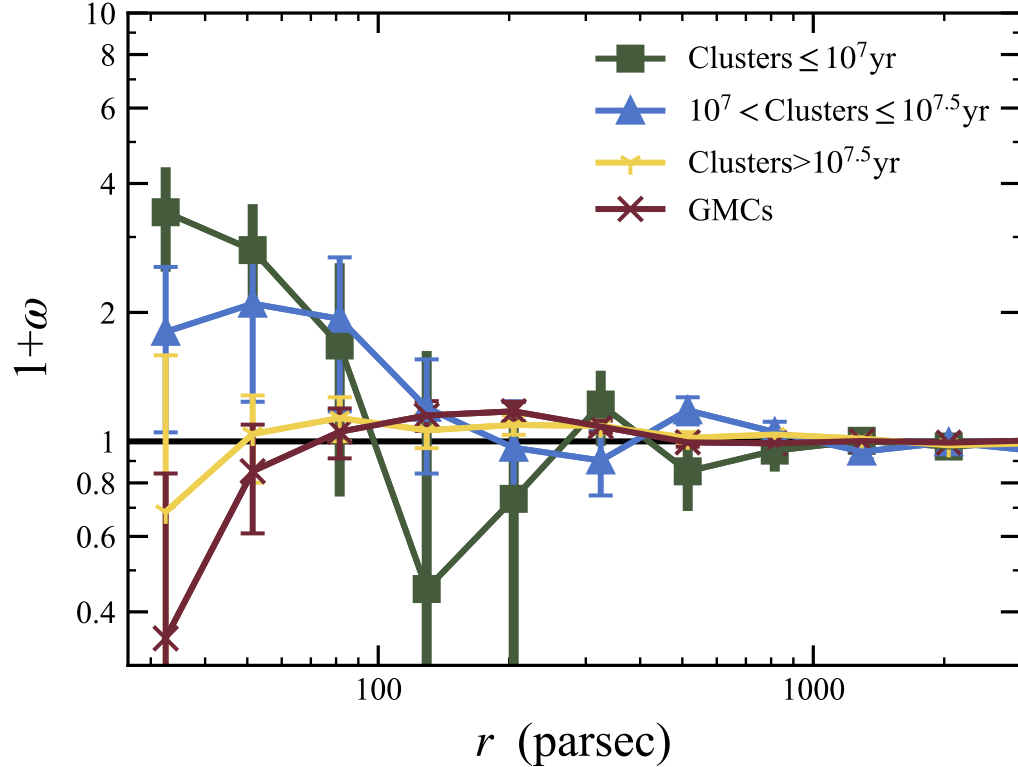


Figure 2.5: The two-point correlation function,  $1 + \omega$ , at radial separations,  $r$  in M33. The two-point correlation is shown for the GMCs as a dark pink line with crosses. The two-point correlation is also shown for the youngest clusters (green line with squares), medium-aged clusters (blue line with triangles), and the oldest clusters (yellow line with tri-points). These two-point correlations are calculated with respect to the 100 random cluster distributions and averaged. The error bars show the standard deviation of 100 two-point correlations. The black line marks uncorrelated. This plot shows that the youngest and medium-aged clusters are correlated, and the GMCs are anti-correlated at small radii. At larger radii, all the groups become uncorrelated.

$\sim 100$  pc. The oldest clusters are uncorrelated at all scales. GMCs show anticorrelation on small ( $< 50$  pc) scales, which we attribute to the object identification algorithm. When the SCIMES algorithm defines the local maxima of GMCs, it requires a minimum spatial separation between the maxima, which is set to 50 pc (Koch et al. in prep). As the radial separation increases beyond 100 pc, all catalogues tend toward  $1 + \omega = 1$ , which indicates they are uncorrelated on large scales. The youngest clusters show modest anticorrelation on medium scales (200 pc) and are uncorrelated at larger scales. This anticorrelation is likely the consequence of strong correlation at small scales ( $< 100$  pc). Not shown in Figure 2.5 is the two-point correlation of the YSCCs from Corbelli et al. (2017) because we find no significant correlation at any scale.

We also tried to test the effects of completeness by removing clusters below a certain mass. We removed the clusters below  $10^{3.5} M_{\odot}$  where the cluster catalogue is  $\approx 90\%$  complete. This high level of completion should eliminate the effects of difficulty identifying older clusters and in crowded regions. There are fewer clusters which makes the correlation structure less consistent. However, the trend for decreasing correlation strength with age is still apparent. Therefore, we assume the change in completeness due to crowding and cluster ageing will not significantly impact our results. Performing a similar test for the GMCs, removing the GMCs below  $3.6 \times 10^4 M_{\odot}$  leaving only the GMCs that are  $\approx 90\%$  complete. Removing these lower-mass clouds has very little effect on the two-point correlation of the GMCs. Grasha et al. (2018, 2019) found that removing lower mass clusters and GMCs resulted in higher correlation magnitudes. We see this effect with the clusters but not with the GMCs.

The results of our two-point correlation analysis are consistent with what other studies have found. Grasha et al. (2018, 2019) and Turner et al. (2022) all found that young clusters ( $< 10$  Myr) have stronger correlation than the older clusters. While the main results are consistent, there are two major differences between our results and previous studies. In NGC 7793, Grasha et al. (2018) found the two-point correlation

using a catalogue of 293 clusters using the same estimator. However, Grasha et al. (2018) find a larger magnitude of correlation and the clusters remain correlated until separations of approximately 1000 pc instead of 100 pc. Grasha et al. (2019) and Turner et al. (2022) found the same difference in magnitude and scale in the galaxy M51 and 11 PHANGS galaxies, respectively. This difference likely comes from the inherent clustering of stars in a galaxy that we have tried to account for by using exponential random distributions fitted to our clusters.

Another important result that is consistent between our results and Grasha et al. (2018, 2019) and Turner et al. (2022) is that the GMCs are much closer to a random distribution than the clusters. If each GMC produced only a single cluster it would be expected for the correlation structure of the GMCs and young clusters to match. The strong correlation seen in the youngest clusters and the lack of correlation in GMCs could suggest that GMCs produce multiple clusters. However, we find that none of our GMCs are overlapping with several of the youngest clusters, which does not fully rule out this possibility since clusters drift. Another possible solution is that enough GMCs are quickly destroyed by young clusters that the correlation structure of the GMCs is erased. Finally, a more mundane solution would be to attribute the lack of correlation to the cloud identification algorithm, which suppresses the correlation structure in the molecular clouds at short scales (Figure 2.5) by forcing them into discrete, well separated units. The resolution of the ALMA ACA survey could prevent the cloud identification algorithm from distinguishing between a complex of smaller clouds and one large cloud. We will address this possibility by analyzing the correlation of all of the  $^{12}\text{CO}$  emission without cloud decomposition in Peltonen et al. (prep).

### **2.3.3 The Cloud-Cluster Cross Correlation**

We want to understand how cluster distributions are related to the GMC distribution at different scales, which can be shown using the cross-correlation function. The cross-correlation function,  $\zeta(r)$  indicates the excess probability that two data sets

are jointly clustered more than two random data sets at a distance  $r$ . As with the two-point correlation function,  $1 + \zeta = 1$  marks the boundary between correlated and anti-correlated. We use the cross-correlation function to find the correlation between GMCs and clusters in the three age categories. To estimate this cross-correlation, we use the three random cluster distributions and the random GMC distributions. Then by repeating this process with all 100 respective random distributions, we find an average cross-correlation. For each radial bin, we find the number of real cluster GMC pairs  $D_{sc}D_{gmc}$ , real cluster random GMC pairs  $D_{sc}R_{gmc}$ , random cluster real GMC pairs  $R_{sc}D_{gmc}$ , and random cluster random GMC pairs  $R_{sc}R_{gmc}$ . Again, we use the Landy & Szalay (1993) estimator in the same form as Turner et al. (2022):

$$\zeta = \left( \frac{N_{R_{sc}} N_{R_{gmc}} D_{sc} D_{gmc}}{N_{D_{sc}} N_{D_{gmc}} R_{sc} R_{gmc}} \right) - \left( \frac{N_{R_{sc}} D_{sc} R_{gmc}}{N_{D_{sc}} R_{sc} R_{gmc}} \right) - \left( \frac{N_{R_{gmc}} R_{sc} D_{gmc}}{N_{D_{gmc}} R_{sc} R_{gmc}} \right) + 1, \quad (2.3)$$

where  $N_{D_{sc}}$  is the number of real clusters,  $N_{D_{gmc}}$  is the number of real GMCs,  $N_{R_{sc}}$  is the number of random clusters, and  $N_{R_{gmc}}$  is the number of random GMCs. Figure 2.6 shows the result of finding the average cross-correlation between the GMCs and the clusters of the three age categories. The standard deviations of the 100 cross-correlations are shown in Figure 2.6 as error bars. The youngest clusters are correlated with GMCs at small radial separations. The medium-aged and oldest clusters are uncorrelated at most separations. In the first bin, the medium-aged clusters are anti-correlated. However, this could be explained by the large standard deviations, smaller number of sources at small separations, and extinction from the GMCs. The older clusters would be less subject to the extinction since older clusters have larger scale heights. The youngest clusters are likely also affected by this extinction. The correlation at small scales is expected to be quite strong, which could be partially overcoming the extinction, and the missing embedded clusters would likely amplify the correlation strength at these small separations. Removing the clusters and GMCs below the 90% completeness limits slightly reduces the correlation magnitudes. How-



ever, the overall shape of the correlations remains unchanged and our conclusions are the same.

The correlation between the Corbelli et al. (2017) YSCCs in our survey area (244 out of 630) and the GMCs is similar to our youngest clusters. We have found this correlation with respect to a random distribution with a scale length of 2.9 kpc, which was found in the same way as the other random exponential distributions. The YSCCs become uncorrelated at a smaller scale than our youngest clusters. The YSCCs also have a greater magnitude of correlation at the smallest scale. This increased correlation could be due to the very young  $\leq 3$  Myr clusters that cannot be identified in the optical.

Fewer studies have analyzed cross-correlation than two-point correlation. However, our results appear consistent with what Turner et al. (2022) found for a sample of 11 PHANGS galaxies. There is some variation over the 11 galaxy sample, but Turner et al. (2022) find that clusters younger than 10 Myr have similar correlation magnitudes to what we have found. There is still a difference in random catalogues, but Turner et al. (2022) find that older clusters are typically much less correlated with GMCs, similar to our analysis. This correlation between young clusters and GMCs allows for two possibilities. The first possibility is that clusters are still nearby to their parent GMC, which means that the cluster has not had sufficient time to fully destroy the progenitor cloud. The second possibility is that the cluster has destroyed its true progenitor cloud but, is still associated with a nearby GMC. As discussed in Section 2.3.2 our cloud identification algorithm might be identifying complexes of multiple GMCs as a single larger GMC which would erase the two-point correlation of the GMCs. Therefore, it would be difficult to distinguish between these two possibilities using these methods. We note that, unlike the real GMC catalogue, objects in our random GMC catalogue are not required to be separated by 50 pc. Therefore, the random GMC catalogue is more correlated at small scales, which could decrease the cross-correlation magnitudes at these small scales.

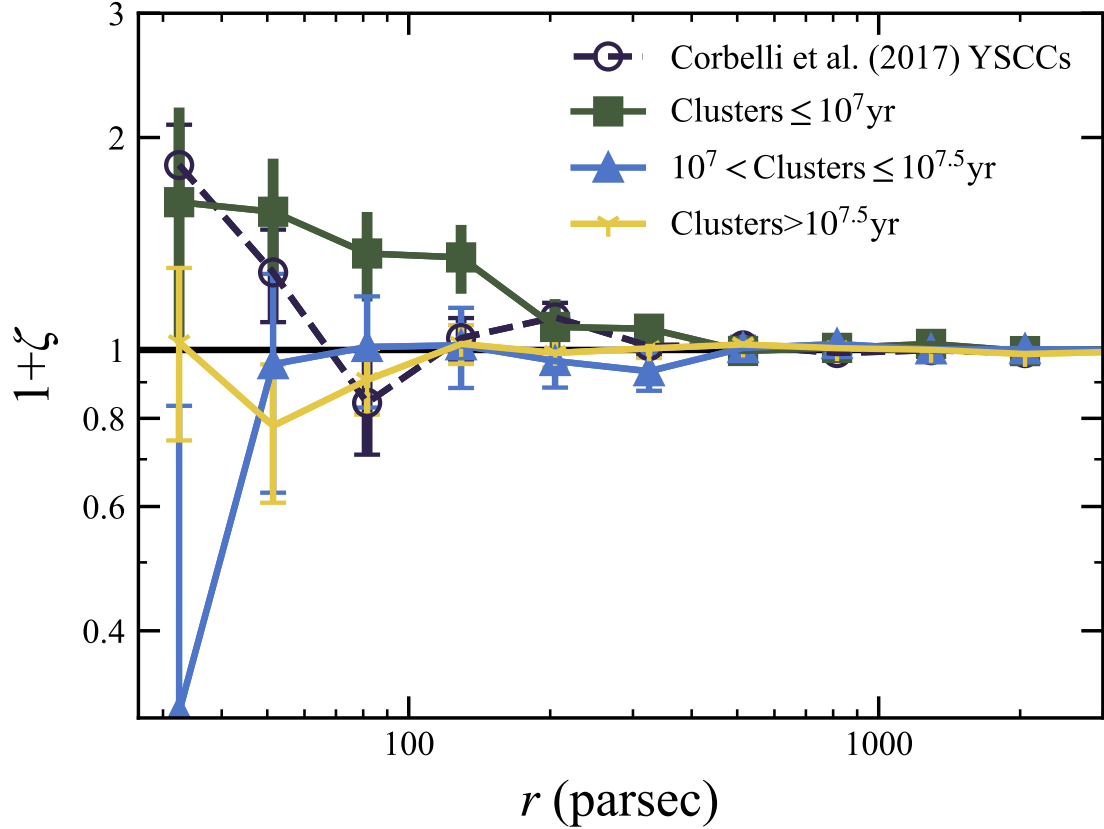


Figure 2.6: The cross-correlation function,  $1 + \zeta$ , at radial separations,  $r$  between the clusters and GMCs in M33. The cross-correlation is shown for GMCs correlated with the youngest clusters (green line with squares), medium-aged clusters (blue line with triangles), and the oldest clusters (yellow line with tri-points). The blue dotted line with circles shows the cross-correlation between the Corbelli et al. (2017) YSCCs and GMCs. These cross-correlations have been found with respect to the 100 random cluster and GMC distributions and averaged. The error bars show the standard deviation of 100 cross-correlations. The black line marks uncorrelated. This plot shows that the youngest clusters are correlated with the GMCs at small radii.

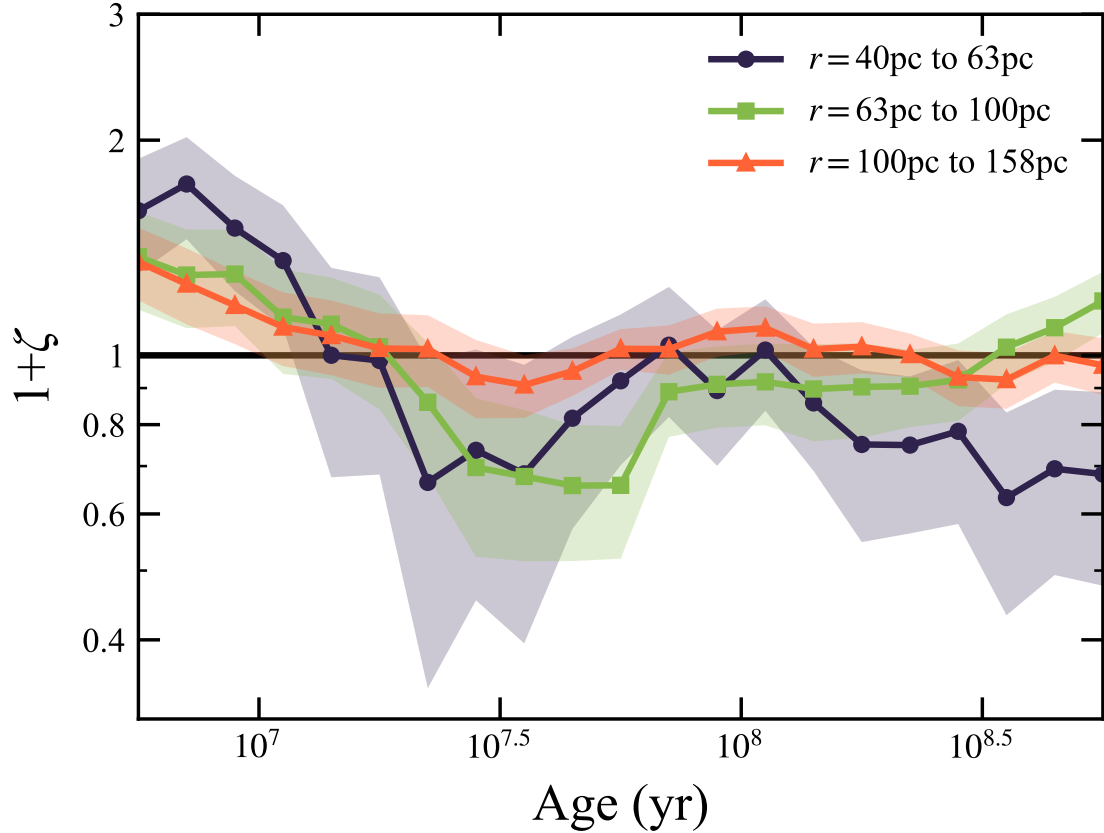


Figure 2.7: The cross-correlation function,  $1 + \zeta$ , at specific radial bins, versus cluster age in M33. The second (dark blue line with dots), third (light green line with squares), and fourth (orange line with triangles) bins from Figure 2.6 are plotted. Where the age bins are 0.5 dex wide. The cross-correlations have been found with respect to the 100 random cluster and GMC distributions and averaged. The shaded regions show the standard deviation of the 100 cross-correlations. The black line marks uncorrelated. The cross-correlation between clusters and GMCs decreases with age.

We now determine at what age clusters transition from correlated to uncorrelated with GMCs. Figure 2.7 shows the cross-correlation at specific radial bins broken into smaller, overlapping age groups. The age groups are logarithmically spaced with a width of 0.5 dex, starting with  $10^{6.5} - 10^7$  yr and increasing by a factor of 0.1 dex. We choose overlapping age bins to increase the number of sources and reduce noise. The radial bins used (second, third, and fourth bins from Figure 2.6) are where the youngest clusters are correlated with GMCs and the medium-aged clusters are not anti-correlated. To find the scale length of the random distribution for each overlapping bin we assume the scale length increases linearly with age. Figure 2.7 shows that the clusters begin with being correlated with GMCs at these small separations, and then at  $\approx 18$  Myr, all bins are consistent with being uncorrelated. The three radial bins in Figure 2.7 behave similarly before  $10^7 - 10^{7.5}$  yr, but with different amplitudes of correlation. The longest radial bin (100 pc to 158 pc) trends directly to uncorrelated at  $\approx 18$  Myr while the shorter bins become anti-correlated. This anti-correlation could be due to the larger correlation at younger ages or from extinction from GMCs.

### 2.3.4 Properties Of Associated Clusters

Since the molecular cloud population is the site of star formation, the resulting star formation properties and cluster population should depend on the properties of the progenitor clouds (Kruijssen, 2012; Krumholz et al., 2019). We look at the properties of the clusters younger than 10 Myr that are likely still correlated with GMCs. We find, similar to other studies (Corbelli et al., 2017; Grasha et al., 2018, 2019), that the properties of the nearest cloud (like mass, radius, and surface density) have no significant correlation with the mass of the clusters produced. This lack of correlation could mean cluster properties are not determined by the properties of the cloud. However, this lack of correlation can also be explained by the evolution of GMCs. It has been suggested that GMCs continually accrete additional gas throughout their

lifetimes (Fukui et al., 2009; Gratier et al., 2012). Then once the GMC has produced a sufficient number of stars, the GMC is dispersed through stellar feedback. Therefore, the properties of GMCs are not constant, and the properties of the nearest GMC to a cluster are unlikely to be the same as the progenitor cloud’s properties at the time of formation.

One property that will be more difficult to erase via feedback is the direction a cluster leaves its progenitor cloud. If a cloud forms multiple clusters or clouds exist in complexes, the direction of the correlated young cluster ( $\leq 10$  Myr) will be preserved. However, if the progenitor cloud is destroyed and there are no GMCs near the progenitor cloud then the direction will not be preserved. We select the nearest young cluster-cloud pairs and find the angle from the cloud to the cluster. We define the angle to start from  $0^\circ$  pointing from the molecular cloud to the galactic centre and with  $90^\circ$  pointing in the direction of galactic rotation. The histogram in Figure 2.8 shows this cloud to cluster angle broken into  $20^\circ$  bins. We then found the angles from our molecular clouds to the 100 random young cluster distributions ( $R_d=1.6$  kpc). The 100 angle distributions are then averaged, which is shown in Figure 2.8 as a red line. The red shaded region in Figure 2.8 shows the standard deviation in the 100 random angular distributions. Young clusters are marginally more likely to be at  $90^\circ$  (in the direction of galactic rotation) and  $270^\circ$  (in the opposite direction of rotation). We recognize that the significance of this result is weak, but we find it notable that the peaks in the real angular distributions lie in the  $90^\circ$  and  $270^\circ$  directions, though nothing in the analysis favours these directions. Even after varying the bin positions and widths, the peak at  $270^\circ$  is persistent. However, the strength of the peak at  $90^\circ$  is diminished in some variations of binning.

### 2.3.5 Molecular Cloud Lifetimes

Using our cluster ages, we make an estimate of the total molecular cloud lifetime following a similar procedure as used in Kawamura et al. (2009). If we assume clusters

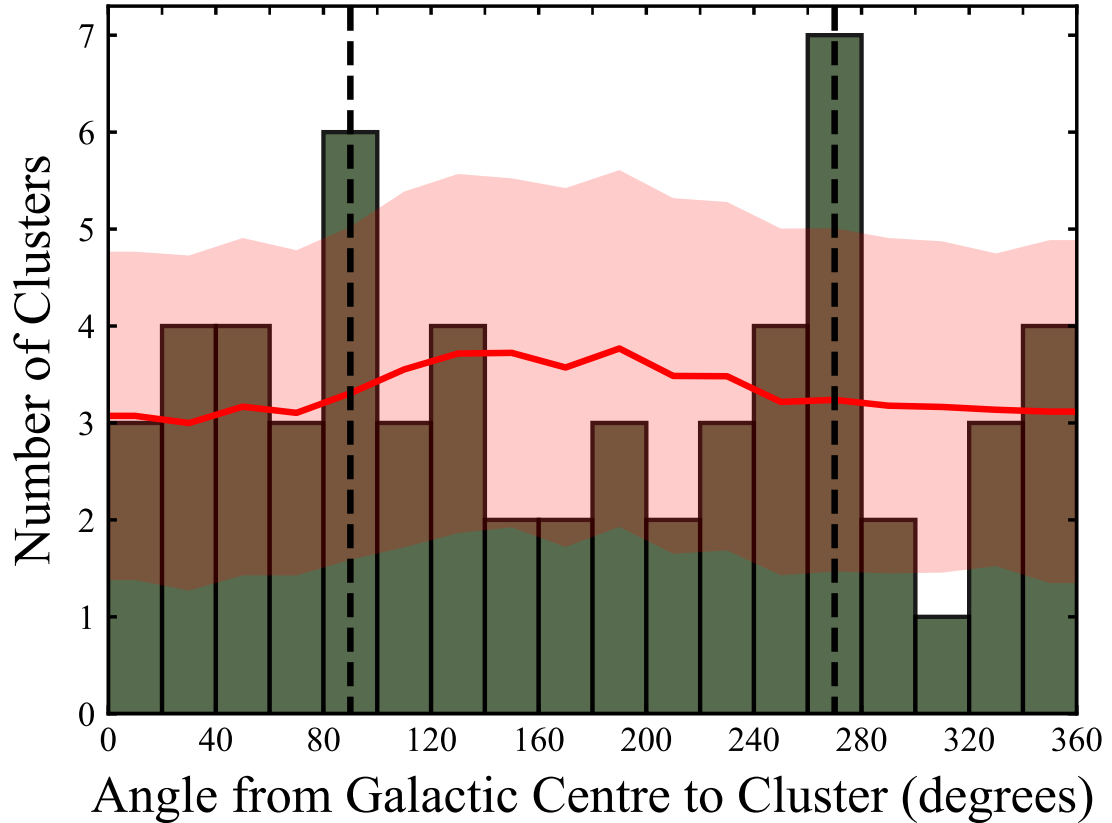


Figure 2.8: The angular distributions of real and random young clusters ( $\leq 10$  Myr) in M33. The green histogram shows the distribution of angles between pairs of young clusters and their nearest GMC. The angle is defined from the GMC to the young cluster with  $0^\circ$  pointing towards galactic centre. The mean random angular distribution from the 100 random young cluster positions is shown as a red line. The red shaded region shows the standard deviation in the 100 random angular distributions. The angular bins are  $20^\circ$  wide. The vertical dashed lines indicate  $90^\circ$  and  $270^\circ$ . The real angular distribution of the young clusters is slightly different from random clustering.

are formed at a constant rate, the fraction of clusters in a phase represents the time spent in that phase. We divide cloud lifetime into two stages:  $\tau_{\text{dark}}$  where a cloud shows no association with any potential disrupting cluster and  $\tau_{\text{fb}}$  where a stellar source could be providing feedback. Then,  $\tau_{\text{GMC}} = \tau_{\text{dark}} + \tau_{\text{fb}}$ .

To determine the length of  $\tau_{\text{fb}}$  we only use the young clusters ( $\leq 10$  Myr) from Johnson et al. (2022) that are believed to be long-lived. We find the fraction of our youngest clusters overlapping with their parent GMC. We use this fraction of association to estimate how long clusters spend with their GMC. This fraction of 10 Myr will be known as  $\tau_{fb}$ , the feedback phase. We find that  $N_{\text{overlap}} = 23$  out of  $N_{\text{total}} = 60$  young clusters are associated with molecular clouds. Based on the fraction of associated young clusters, we have

$$\tau_{\text{fb}} = 10 \text{ Myr} \left( \frac{N_{\text{overlap}}}{N_{\text{total}}} \right) \approx 4 \text{ Myr}. \quad (2.4)$$

However, due to visual extinction, we are likely missing some clusters younger than 3 Myr. If we assume we are missing every deeply embedded cluster  $< 3$  Myr ( $\approx 30\%$ ), then  $\tau_{\text{fb}} \approx 6$  Myr. Therefore, we estimate the feedback phase to be 4-6 Myr.

We can now estimate the total lifetime of the GMCs by finding the fraction of GMCs in the feedback phase. This fraction of GMCs must include not only the long-lived clusters but also other sites of high-mass star formation that are visible but not classified as clusters in Johnson et al. (2022) (e.g., OB associations). Then we have:

$$\frac{\tau_{\text{fb}}}{\tau_{\text{dark}} + \tau_{\text{fb}}} = \frac{N_{\text{GMC,fb}}}{N_{\text{GMC}}} \quad (2.5)$$

where  $N_{\text{GMC,fb}}$  is the total number of GMCs experiencing feedback. Koch et al. (in prep) identified 217 GMCs associated with recent high-mass star formation through a visual inspection which includes many of the clusters from Johnson et al. (2022). However, this inspection was done without the Wainer et al. (2022) ages for the long-lived clusters. We visually inspect the 217 GMCs with visible clusters and find that 56 contain only an old cluster ( $> 10^7$ yr). These clusters that are much older

than  $\tau_{\text{fb}}$  are likely not in the same plane as the GMC and do not represent a GMC in the feedback phase. The Koch et al. (in prep) GMCs without the old clusters leave  $N_{\text{GMC,fb}} = 161$  out of  $N_{\text{GMC}} = 444$  GMCs in the feedback phase. Therefore, the feedback phase represents  $\approx 35\%$  of the total lifetime of GMCs in M33, giving a lifetime of  $\tau_{\text{GMC}} = \tau_{\text{dark}} + \tau_{\text{fb}} = N_{\text{GMC}}\tau_{\text{fb}}/N_{\text{GMC,fb}} = 11\text{-}15$  Myr. However, this lifetime estimate is sensitive to the visual cluster identification and the removal of older clusters. If the old clusters are not removed, the lifetime estimate would be 8-12 Myr. Regardless of this sensitivity, our GMC lifetime estimate is consistent with the short lifetimes found by previous studies (e.g. Kawamura et al., 2009; Blitz et al., 2007; Miura et al., 2012; Corbelli et al., 2017; Kruijssen et al., 2019; Chevance et al., 2020; Kim et al., 2021a; Pan et al., 2022; Kim et al., 2022).

## 2.4 Cloud-Cluster Population in M31

We now analyze the clusters (PHAT) and GMCs (CARMA) in M31 and compare these results to those found in M33. However, the different conditions of M31 only allow for certain results to be compared. M31 has a greater inclination (Table 2.1) than M33, making photometry more complex and limiting the age estimates for clusters. In addition, the limited survey area and the more defined rings of M31 make creating random catalogues more difficult. Despite these differences, we find the separations between the clusters and GMCs in M31 and compare them to M33's separations, which depend on cluster age. We then find the two-point correlation of the individual objects, which can confirm some of the properties of our methods. Finally, we find the cross-correlation between the clusters and GMCs in M31 to confirm that cross-correlation depends on cluster age.

### 2.4.1 Cloud-Cluster Spatial Offsets

Similar to M33, in section 2.3.1, we want to find the positions of the clusters and GMCs on the disc where the survey areas overlap. Including objects where PHAT



and CARMA data overlap leaves 480 of 1249 clusters and all 251 GMCs. We can then use the orientation parameters of M31 from Table 2.1, to convert the celestial coordinates to galactocentric coordinates.

The radial distribution of M31 in Figure 2.9 looks very different than the radial distributions of M33 (Figure 2.3). The GMCs and the clusters split into age categories follow a double peak structure offset from zero. The double peak comes from the odd shape of the CARMA survey area and because the survey area is centred on the star-forming rings of M31. It is also notable that M31 is much larger than M33. Therefore, the clusters and GMCs span a much larger area. The analysis of cluster ages in M31 restricted age values to be between  $10^7$  yr and  $10^{8.5}$  yr. Therefore, we choose only two age groups the 96 medium-aged clusters have ages  $\leq 10^{7.5}$  yr and the 384 oldest clusters with ages  $> 10^{7.5}$  yr. This more limited age range prevents a completely parallel analysis of M33.

Generating random distributions in the same way as M33 by fitting an exponential distribution to the clusters in the survey area yields a flat distribution ( $R_d = \infty$ ). This is likely a consequence of the defined ring structure present in the survey area. A flat distribution does not replicate the observed radial distributions of our clusters. To create a random distribution that matches the defined rings, we must sample a distribution with a similar structure. Therefore, we use the full-galaxy  $22\mu\text{m}$  image constructed by Leroy et al. (2019) using data from the *Wide-field Infrared Survey Explorer* (WISE) satellite (Wright et al., 2010). We azimuthally average the emission in the WISE filters to create a distribution that can be sampled. The filter W4 ( $22\ \mu\text{m}$ ) produces random distributions that best match the radial distributions of the clusters and GMCs in M31. The mean of 100 distributions produced from W4 is shown in Figure 2.9 as a red line.

Now that we have the positions of the clusters and GMCs in M31, we find the typical separations between clusters and GMCs. Figure 2.10 shows the separations between the centre of clusters split by age and the CO peaks of the nearest GMC

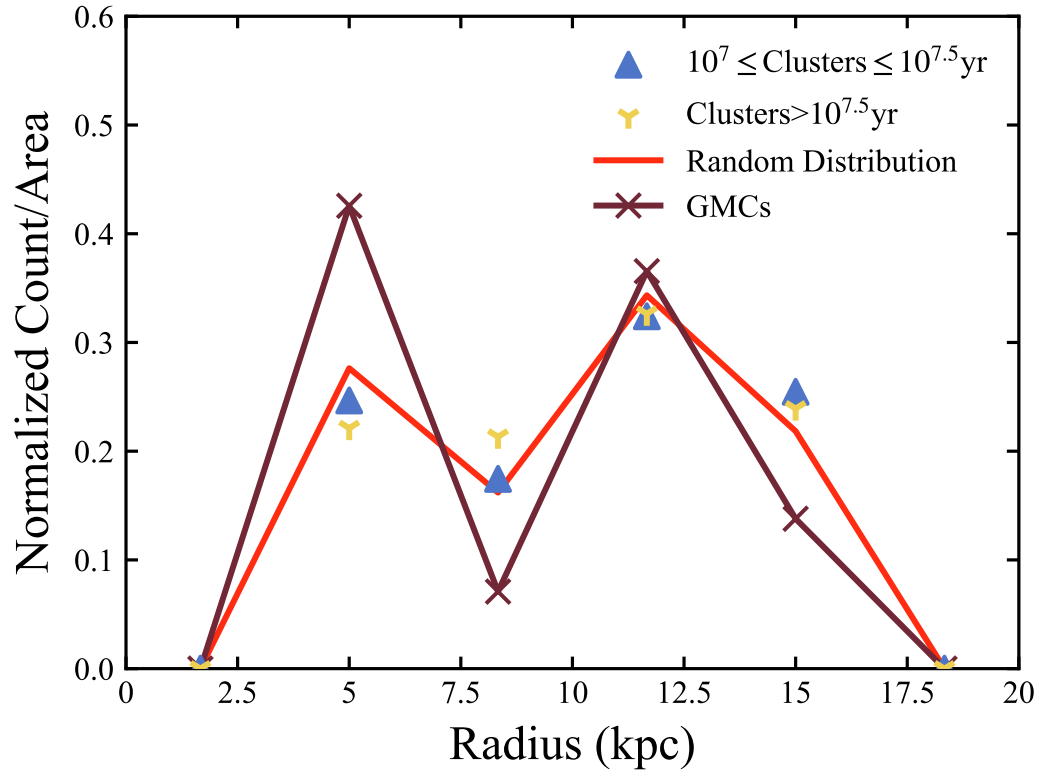


Figure 2.9: The normalized count per area of the clusters, GMCs, and the generated random clusters. The count is found in galactocentric radial bins and then divided by the area contained in the overlapping survey region in that radial bin. The clusters are split into the two age categories medium-aged (blue triangles) and oldest (yellow tri-points). The GMC radial distribution (dark pink line with crosses) is found from the coordinates of the CO peak of each cloud. The random cluster distribution (red line) results from averaging 100 distributions generated from WISE. The random distribution traces a similar distribution to the clusters in M31.

presented in the same way as Figure 2.4. In this section, we split the medium-aged clusters into two smaller age groups since this analysis is less sensitive to the number of objects than the two-point correlation and cross-correlation. The clusters  $\leq 10^{7.2}$  yr have the lowest median separation of 210 pc and the smallest IQR of 160 pc. The clusters  $> 10^{7.2}$  yr and  $\leq 10^{7.5}$  yr have a median separation of 280 pc closer to random (270 pc) but with a smaller IQR of 180 pc than random (280 pc). The oldest clusters have median separations of 270-290 pc and IQRs of 260-300 pc, similar to random. Even with a more limited cluster age range, the same trends are visible in M31 that are seen in M33. We find that the younger clusters are closer to GMCs and the oldest clusters are similar to a random distribution. However, the scale of the separations is larger in M31 because there are fewer objects in a larger area. In M31, the clusters have a density of  $\approx 3 \text{ kpc}^{-2}$  and the clouds have a density of  $\approx 1 \text{ kpc}^{-2}$ . While in M33, the density of clusters is  $\approx 36 \text{ kpc}^{-2}$  and the density of clouds is  $\approx 17 \text{ kpc}^{-2}$ .

### 2.4.2 Two-Point Correlation and Cross-Correlation in M31

We calculate the two-point correlation of the clusters and GMCs in M31. The two-point correlation is found in the same way described in section 2.3.2, showing the excess probability of clustering compared to 100 random distributions. The two-point correlation in different radial bins is shown in Figure 2.11 for the GMCs and the clusters in their age categories. The GMCs are anti-correlated at small scales ( $< 50$  pc) due to the decomposition algorithm. The GMCs are then correlated until  $\approx 1000$  pc. The oldest clusters have a small amount of correlation until they become uncorrelated at  $\approx 1000$  pc. The medium-aged clusters have large correlation magnitudes at small separations that decrease until becoming uncorrelated at  $\approx 1000$  pc. Similar to what we found in M33 we see that the younger clusters have greater correlation magnitudes than the older clusters. The two-point correlation in M31 differs from M33 because the clusters become uncorrelated at much larger scales ( $\approx 1000$  pc compared to  $\approx 100$  pc in M33). The two-point correlation in M31 looks more similar

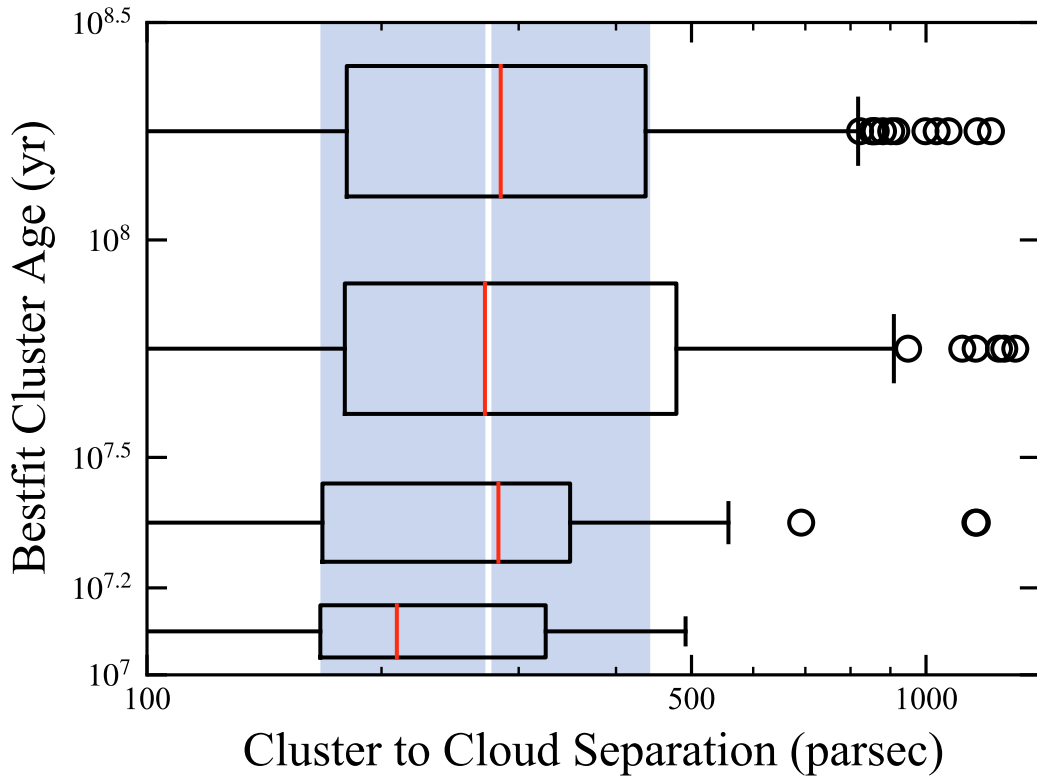


Figure 2.10: The separation between the clusters and the GMCs based on the age of the clusters in M31. The boxes show the IQR for each age bin, with the medians marked with red lines. The error bars extending from the boxes indicate each bin's minimum and maximum values, excluding the outliers. The outliers, marked with circles, are defined as points outside 1.5 times IQR. The age bins correspond to the youngest clusters ( $10^7$  yr- $10^{7.2}$  yr), medium-aged clusters ( $10^{7.3}$  yr- $10^{7.5}$  yr), and two bins for the oldest clusters both  $10^{0.5}$  yr wide. The separation between clusters with random positions and the GMCs have their median (white line) and IQR (blue shaded region) plotted. The youngest clusters have a shorter median and a smaller IQR than random clusters. Clusters older than  $10^{7.5}$  yr have medians and IQRs consistent with the random distribution.

to NGC 7793 (Grasha et al., 2018), where objects become uncorrelated at 1000 pc, and the GMCs have small magnitudes of correlation. These two differences could indicate that our random distribution generated from WISE has failed to account for additional large-scale correlation present in M31’s defined ring structure. Therefore, WISE W4 may not be the best representative of the full cluster population.

As we did for M33, we find the cross-correlation between clusters and GMCs in M31. The methods for finding cross-correlation are the same as presented in section 2.3.3, finding the excess probability that the clusters are associated with GMCs compared to 100 random distributions. The results of the cross-correlation analysis for M31 are shown in Figure 2.12 for the medium-aged clusters and the oldest clusters. There is no significant correlation for clusters at any age, which could be explained by the lack of clusters  $< 10^7$  yr. Similar to what we find in M33 and consistent with Turner et al. (2022), there is no correlation for clusters  $> 10^7$  yr.

## 2.5 Discussion

### 2.5.1 Cluster Drift Speed Estimates

Our timescale measurements support the idea that the star-formation process occurs rapidly. Once a cluster is formed, it will spend  $\approx 4-6$  Myr emitting radiation into its parent GMC. Our results also show that GMCs will only survive approximately 11-15 Myr. Therefore, a GMC will experience feedback from clusters for a significant fraction of its lifetime. Our correlation analysis shows that young clusters are correlated with GMCs for  $\approx 18$  Myr. This loss of correlation likely comes from the original cloud being dispersed and the clusters drifting.

Based on the timescales from our correlation analysis, we can estimate the drift speed, which represents the speed at which the young cluster population decouples from its birth molecular gas. From Figure 2.6, we see that the correlation of the youngest clusters and GMCs is no longer present at a scale of 200 pc. Then, from

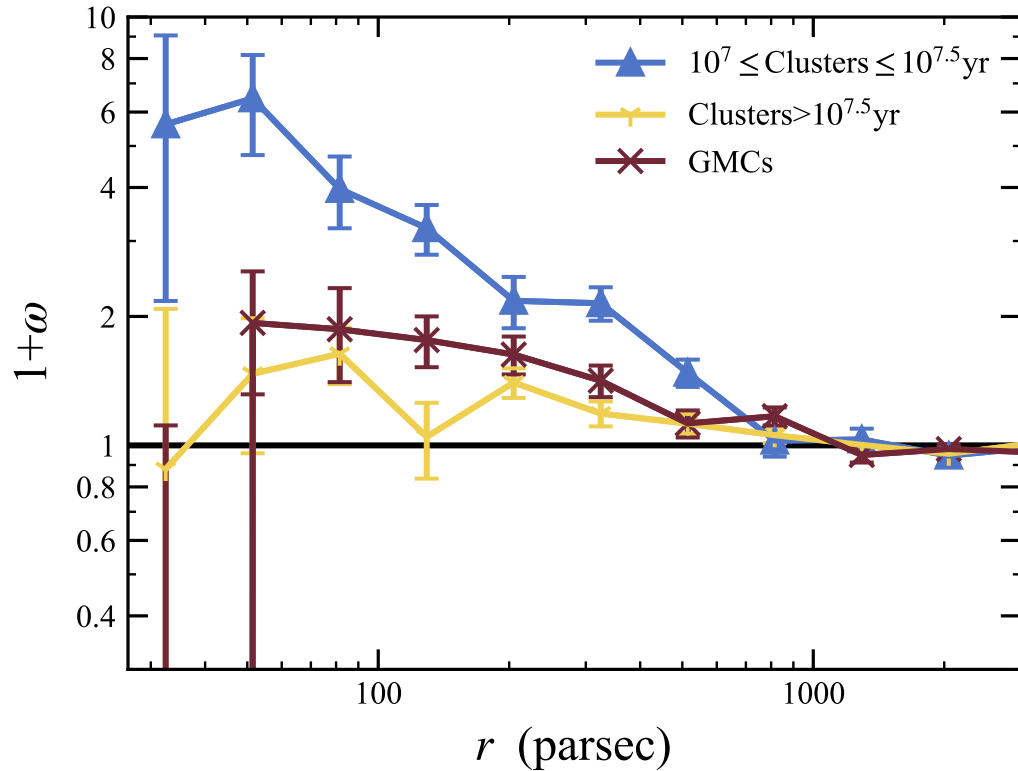


Figure 2.11: The two-point correlation function,  $1 + \omega$ , at radial separations,  $r$  for M31. The two-point correlation is shown for the GMCs as a dark pink line with crosses. The clusters have their two-point correlation shown for the medium-aged clusters (blue line with triangles) and the oldest clusters (yellow line with tri-points). These two-point correlations have been found with respect to the 100 random cluster distributions and averaged. The error bars show the standard deviation of 100 two-point correlations. The black line marks uncorrelated. This plot shows that medium-aged clusters are correlated at small radii. At larger radii, all the groups become uncorrelated.

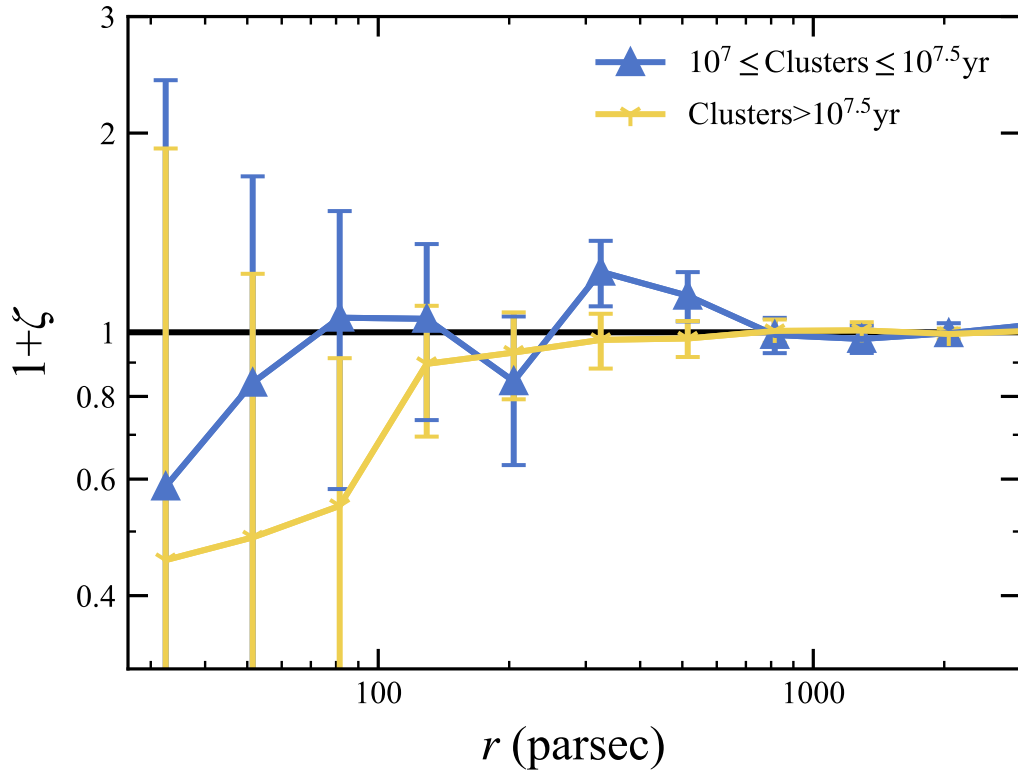


Figure 2.12: The cross-correlation function,  $1 + \zeta$ , at radial separations,  $r$  between the clusters and GMCs in M31. The cross-correlation is shown for GMCs correlated with the medium-aged clusters (blue line with triangles) and the oldest clusters (yellow line with tri-points). These cross-correlations have been found with respect to 100 random cluster and GMC distributions that are both from WISE W4. The error bars show the standard deviation of 100 cross-correlations. The black line marks uncorrelated. This plot shows that there is no significant correlation for clusters of any age in M31.

Figure 2.7, the correlation between clusters and GMCs disappears at an age of  $\approx 18$  Myr. Simply dividing this spatial scale by the temporal scale gives a velocity of  $\approx 10 \text{ km s}^{-1}$ . It's important to note that this estimate assumes that the lack of correlation at 200 pc is purely due to cluster drift. In reality, this lack of correlation is likely due to a combination of cluster drift and cloud destruction. If a cluster destroys its progenitor cloud, it may still be associated with a more distant GMC in the same complex, which would increase the spatial correlation scale. Therefore, this estimate of drift velocity should be seen as an upper limit.

We can refine our simple drift model by trying to match the two-point correlation structure seen in M33. We start with a desired correlation model matching the two-point correlation of the young clusters in Figure 2.5. A power spectrum can then be generated from Fourier transformation of the model correlation using the relation

$$P(k) \propto \int_{\mathbb{R}^2} \omega(\mathbf{r}) e^{i\mathbf{k}\cdot\mathbf{r}} d\mathbf{r}, \quad (2.6)$$

where  $k$  is the wave vector,  $r$  is the correlation scale, and  $\omega(r)$  is the model correlation at that scale. We then calculate the Fourier transform of the density field as

$$A(\mathbf{k}) \propto \sqrt{P(k)} e^{i\phi(\mathbf{k})} \quad (2.7)$$

where  $\phi(\mathbf{k})$  is a uniformly distributed random phase factor. The density field is then generated and normalized so that, when sampled, it produces a cluster population with a similar correlation structure to the observations. These clusters are then given drift velocities drawn from three different velocity dispersions. These three velocity dispersions (measured in the 2D plane of the galaxy) are  $20 \text{ km s}^{-1}$ ,  $10 \text{ km s}^{-1}$ , and  $5 \text{ km s}^{-1}$ . There is about 10 Myr between our youngest and medium-aged clusters and about 100 Myr between our youngest and oldest clusters. Therefore, we let the clusters drift for 10 Myr and 100 Myr and found the two-point correlation for each time and velocity dispersion. This process is repeated 100 times for each velocity dispersion. Figure 2.13 shows the average two-point correlation for each time and



velocity dispersion. The 10 Myr drift should have a similar two-point correlation to the medium-aged clusters from Figure 2.5, and the 100 Myr drift should be similar to the oldest clusters. All three velocity dispersions seem to replicate the correlation structure of the oldest cluster. However, the magnitude of the correlation of the medium-aged clusters in M33 falls between the  $5 \text{ km s}^{-1}$  and  $10 \text{ km s}^{-1}$  models. Therefore, this simple model predicts a dispersion velocity of 5-10  $\text{km s}^{-1}$  in M33, which is comparable to the upper limit from the simple model inferred from the age argument above.

In order to further test the effects of completeness, we assigned masses to our model clusters. We then removed clusters after 10 Myr and 100 Myr of drifting to simulate decreasing completeness for older clusters. The clusters were removed based on the completeness fits from Johnson et al. (2022). Including this simulated completeness has no noticeable effect on correlation magnitudes. Therefore, we conclude that it is cluster drift and not completeness that results in decreased correlation with cluster age.

Soubiran et al. (2018) measured the velocities of open clusters in the Milky Way and found that young clusters ( $\log(\text{age}) < 7.8$ ) have a velocity dispersion of  $10.6 \text{ km s}^{-1}$  among the clusters. This velocity dispersion in the Milky Way increases with age, so this dispersion is consistent with our drift velocity estimation. In M33, Chandar et al. (2002) find that the velocity dispersion of clusters increases with cluster age and for a small sample of clusters  $< 10^{8.1} \text{ yr}$  find a velocity dispersion of  $17 \text{ km s}^{-1}$ . Therefore, for our larger sample of younger clusters, a lower velocity dispersion than  $17 \text{ km s}^{-1}$  is expected.

Interpreting the physical meaning of this drift velocity can be complicated. Part of this motion represents the random motions of clusters as they are born in a turbulent medium. If the cluster velocity comes from the turbulent motions of the gas, then the dispersion velocity of the gas should be similar to the drift velocity. Koch et al. (in prep) find that the typical line of sight velocity dispersion of CO in molecular clouds

to be  $\sigma_{v,1D} \sim 3 \text{ km s}^{-1}$  (see also Rosolowsky et al., 2003; Gratier et al., 2014), which in 2D assuming isotropy would equate to  $\sigma_{v,2D} \sim 2\sigma_{v,1D} \sim 4.5 \text{ km s}^{-1}$ , lower than but of a consistent scale than the drift speed inferred from correlation structure.

## 2.6 Conclusion

In this paper, we analyze the correlation structure between clusters and GMCs in M33 and M31. We use the PHATTER cluster catalogue from Johnson et al. (2022) with measured CMD ages from Wainer et al. (2022) for M33. In M31, we have the PHAT cluster catalogue from Johnson et al. (2015) with measured CMD ages from Johnson et al. (2016). We use the SCIMES algorithm to find GMCs in the 35 pc resolution ALMA ACA  $^{12}\text{CO}$  survey (Koch et al. in prep) in M33 and the 20 pc resolution CARMA  $^{12}\text{CO}$  survey in M31 (Caldú-Primo & Schrubba, 2016; Leroy et al., 2016; Schrubba et al., 2019). We then generate random catalogues of clusters and GMCs that match the radial distributions of the real clusters and GMCs. Using a random catalogue that matches the radial distribution of clusters and clouds removes the spurious signal from large-scale correlations in the cloud-cluster analyses.

In M33 and M31, we find that younger clusters ( $\leq 10 \text{ Myr}$ ) are found at a shorter distance from the nearest GMCs when compared to older clusters (Figures 2.4 and 2.10). We also find that clusters older than  $\approx 30 \text{ Myr}$  have separations from GMCs that matches a random distribution similar to what was found by Grasha et al. (2018, 2019) and Turner et al. (2022). As seen in Figure 2.8, we find that clusters are marginally more likely to drift in the direction of galactic rotation and in the opposite direction of galactic rotation.

The two-point correlation analysis for M33 shows that younger clusters ( $< 10^7 \text{ yr}$ ) are correlated with each other at small scales ( $< 200 \text{ pc}$ ) but older clusters are not (Figure 2.5). The same correlation analysis in M31 shows similar but weaker trends due to the more limited cluster age information in M31 (Figure 2.11). The younger clusters likely show a stronger correlation because they have not had sufficient time

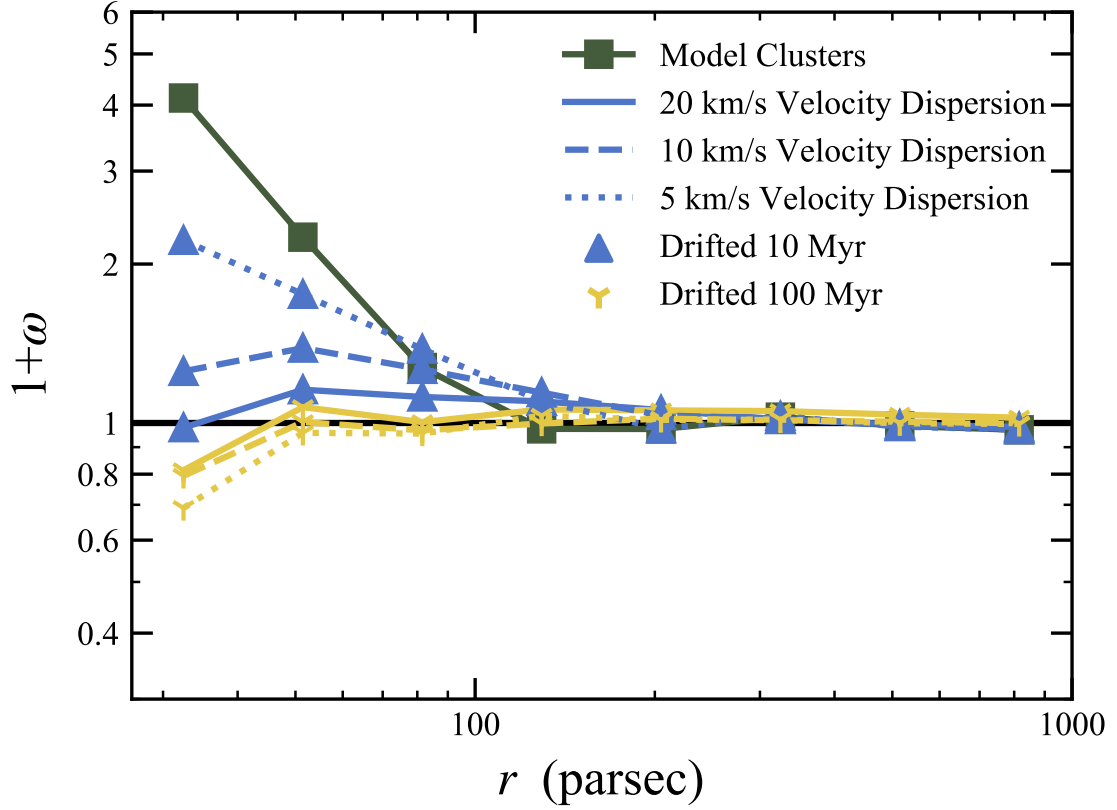


Figure 2.13: Model two-point correlation function,  $1+\omega$ , at radial separations,  $r$ . The two-point correlation of the model clusters to match the youngest clusters in Figure 2.5 is shown as a green line with squares. The different lines show the average two-point correlation of the different velocity dispersions given to the model clusters. The solid lines show a velocity dispersion of  $20 \text{ km s}^{-1}$ , the dashed lines show a velocity dispersion of  $10 \text{ km s}^{-1}$ , and the dotted lines show a velocity dispersion of  $5 \text{ km s}^{-1}$ . The different colours and symbols show the clusters after 10 Myr (blue triangles) and 100 Myr (yellow triangles) with their respective velocity dispersion. The black line marks uncorrelated. Our real two-point correlation in M33 falls between the model with a  $5 \text{ km s}^{-1}$  and a velocity  $10 \text{ km s}^{-1}$  dispersion.

to drift from other young clusters formed in nearby GMCs. Once the clusters are  $\approx 18$  Myr old, they no longer show correlation with GMCs. Because of the limited cluster ages, we find no cross-correlation between clusters and GMCs in M31. By comparing the 18 Myr temporal scale to our 200 pc spatial scale, we estimate the drift velocity of the clusters to be  $10 \text{ km s}^{-1}$ . This is consistent with the drift velocity of  $5\text{-}10 \text{ km s}^{-1}$  found from comparing our two-point correlation structure to a simple drift model applied to mock data.

We estimate the time the clusters spend overlapping with GMCs (feedback phase) to be 4-6 Myr in M33. This estimate comes from finding the fraction of the youngest clusters ( $\leq 10$  Myr) that overlap with GMCs. By finding the fraction of GMCs in the feedback phase, we estimate the total lifetime of GMCs to be 11-15 Myr in M33.

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## Data Availability

This paper makes use of the following ALMA data: ADS/JAO.ALMA#2017.1.00901.S and 2019.1.01182.S. ALMA is a partnership of ESO (representing its member states), NSF (USA) and NINS (Japan), together with NRC (Canada), MOST and ASIAA (Taiwan), and KASI (Republic of Korea), in cooperation with the Republic of Chile. The Joint ALMA Observatory is operated by ESO, AUI/NRAO and NAOJ. The National Radio Astronomy Observatory is a facility of the National Science Foundation operated under cooperative agreement by Associated Universities, Inc. These data are available through the ALMA archive.

Some of the data presented in this paper were obtained from the Mikulski Archive for Space Telescopes (MAST) <http://dx.doi.org/10.17909/t9-ksyp-na40>. All other analysis data are available from the authors upon receiving a reasonable request.

# Chapter 3

## Conclusions & Future Work

### 3.1 Conclusions

In this thesis, I measured the spatial and temporal relationship between star clusters and GMCs in the galaxies M33 and M31. To compare the distributions of clusters and GMCs, I first developed a series of random distributions that matched the radial distributions of the different objects. For M33, an exponential distribution with varying scale lengths produces a similar radial distribution to the GMCs and clusters of different ages. M31, however, has a much different radial distribution from its defined rings, and thus I azimuthally averaged WISE data to produce random distributions. Equipped with these random distributions, I showed that younger clusters ( $< 10$  Myr) are typically found and have smaller separations from GMCs than older clusters, which more closely resemble the random distributions. Using the two-point correlation function, I found that younger clusters are likely to be near other young clusters, while older clusters and GMCs more closely resemble random distributions. I also showed that this effect for clusters becoming uncorrelated as they age can be reproduced when they have velocities  $5\text{-}10 \text{ km s}^{-1}$  in a simple drift model. Then using the cross-correlation function, I found that younger clusters are more correlated with GMCs and that this correlation disappears at an age of  $\approx 18$  Myr, which is faster than the timescale over which two-point correlation functions of older clusters began to become uncorrelated. The cross-correlation being erased at younger ages is likely

a result of cluster drift and GMCs being destroyed by clusters. Finally, I determined two important GMC timescales in M33. The feedback time where clusters overlap with GMCs was found to be 4-6 Myr, and the typical lifetime of GMCs in M33 was found to be 11-15 Myr.

## 3.2 Future Work

The major limiting factor in this thesis was the available data. First, the resolution of the CO surveys limited the effectiveness of the cloud identification, which likely influenced the structure of the correlation functions at small scales. Utilizing data from all of ALMA instead of just data from the ACA could greatly improve the resolution of the CO data in M33. M31 is unfortunately not observable with ALMA, but a larger survey area that matches the PHAT footprint could greatly increase the quality of the statistics. Second, the lack of high-resolution infrared surveys resulted in possibly missing the youngest clusters. However, with new JWST data (GO program 2128) in M33, we can now identify star formation at earlier phases and determine if missing clusters are a concern. The high quality of the JWST data should also allow for individual stars to be observed forming, which will lead to even more direct measures of spatial and temporal timescales in star formation.

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