

23973

National Library
of CanadaBibliothèque nationale
du CanadaCANADIAN THESES
ON MICROFICHETHÈSES CANADIENNES
SUR MICROFICHENAME OF AUTHOR/NOM DE L'AUTEUR JAMAL HUSSEIN ABOU-KASSEMTITLE OF THESIS/TITRE DE LA THÈSE DETERMINATION OF BOTTOM-HOLE PRESSURE
IN FLOWING GAS WELLSUNIVERSITY/UNIVERSITÉ THE UNIVERSITY OF ALBERTADEGREE FOR WHICH THESIS WAS PRESENTED/
GRADE POUR LEQUEL CETTE THÈSE FUT PRÉSENTÉE M.Sc. IN PETROLEUM ENGINEERINGYEAR THIS DEGREE CONFERRED/ANNÉE D'OBTENTION DE CE GRADE 1975NAME OF SUPERVISOR/NOM DU DIRECTEUR DE THÈSE Peter M. DRANCHUK

Permission is hereby granted to the NATIONAL LIBRARY OF
CANADA to microfilm this thesis and to lend or sell copies
of the film.

The author reserves other publication rights, and neither the
thesis nor extensive extracts from it may be printed or other-
wise reproduced without the author's written permission.

L'autorisation est, par la présente, accordée à la BIBLIOTHÈ-
QUE NATIONALE DU CANADA de microfilmer cette thèse et
de prêter ou de vendre des exemplaires du film.

L'auteur se réserve les autres droits de publication; ni la
thèse ni de longs extraits de celle-ci ne doivent être imprimés
ou autrement reproduits sans l'autorisation écrite de l'auteur.

DATED/DATE April 17, 1975 SIGNED/SIGNÉ J. H. Abou KassemPERMANENT ADDRESS/RÉSIDENCE FIXE 5/6 Mineral Engineering Department
U. of A., Edmonton, ALTA. T6G 2E1
CANADA

THE UNIVERSITY OF ALBERTA

DETERMINATION OF BOTTOM-HOLE PRESSURE
IN FLOWING GAS WELLS

by



JAMAL HUSSEIN ABOU-KASSEM

A THESIS

SUBMITTED TO THE FACULTY OF GRADUATE STUDIES AND RESEARCH
IN PARTIAL FULFILMENT OF THE REQUIREMENTS FOR THE DEGREE
OF MASTER OF SCIENCE

in

PETROLEUM ENGINEERING

DEPARTMENT OF MINERAL ENGINEERING

EDMONTON, ALBERTA

SPRING, 1975

THE UNIVERSITY OF ALBERTA

FACULTY OF GRADUATE STUDIES AND RESEARCH

The undersigned certify that they have read, and recommend to the Faculty of Graduate Studies and Research, for acceptance, a thesis entitled "DETERMINATION OF BOTTOM-HOLE PRESSURE IN FLOWING GAS WELLS", submitted by Jamal Hussein Abou-Kassem, in partial fulfilment of the requirements for the degree of Master of Science in Petroleum Engineering.

Dr. D. J. ...

Supervisor

W. H. ...

A. J. ...

Date April 10, 1975

TO THE PEOPLE OF PALESTINE
WITH DEEP RESPECT

ABSTRACT

An existing mathematical model for calculating bottom-hole pressure in flowing single-phase gas wells has been modified. The modification improves the various correlations employed in the model. Furthermore, it extends the application to sour natural gas wells. The modified model was tested against field data. The discrepancies between the modified model predictions and the pressure data are traced to a variety of factors.

Several side benefits have resulted from this investigation. These are:

1. A new mathematical correlation was developed for the ideal isobaric heat capacity, for sweet natural gases over the range 32 - 2240°F and for sour natural gases over the range 32 - 300°F.
2. A new mathematical generalized correlation was developed for the heat capacity departure over a range of reduced pressure of 0 - 15 and reduced temperature of 1.05 - 3.0.
3. Mathematical expressions were developed for the partial derivatives of reduced densities and compressibility factors of natural gases with respect to reduced pressure and temperature.
4. A FORTRAN subroutine was developed that is capable of regenerating the Moody friction factor chart for

relative roughness up to 0.05 and Reynolds numbers greater than 4000.

A method is proposed for the estimation of absolute roughness of tubing. The method is essentially another modification of the existing model.

ACKNOWLEDGEMENT

The author wishes to express his gratitude and appreciation to the following persons and organizations:

Professor P.M. Dranchuk, for his invaluable source of guidance and constructive criticisms throughout this investigation,

Professors D. Quon, D.B. Robinson, A.E. Mather, and D.L. Flock, for their ideas and suggestions which contributed to the development of this work,

The DACS centre personnel of the Department of Chemical Engineering and the Computing Centre of the University of Alberta, for their assistance and co-operation,

The Alberta Energy Resources Conservation Board, for supplying field data,

The Department of Energy, Mines and Resources and the University of Alberta, for their financial support,

His family and his very close friend, Mr. Ahmed M. Khatib, for their financial support and continuous encouragement.

TABLE OF CONTENTS

	Page
LIST OF TABLES.....	xi
LIST OF FIGURES.....	xii
INTRODUCTION.....	1
LITERATURE REVIEW.....	3
Flowing Bottom-Hole Pressure.....	3
Compressibility Factor of Natural Gases.....	3
Isobaric Heat Capacity of Natural Gases.....	5
The Hydraulic Friction Factor.....	6
Absolute Roughness of Gas Well Tubing.....	7
Overall Heat-Transfer Coefficient.....	7
Thermal Conductivity of Gas Mixtures.....	8
Viscosity of Pure and Natural Gases.....	8
MODEL DEVELOPMENT.....	10
Main Model A.....	10
Main Model AA (Method 2).....	14
Main Model B.....	15
Main Model BB (Method 2).....	16
Special Cases of Main Model AA and BB.....	17
1. Adiabatic Flow (Method 1).....	17
2. Complete Solution of the Force- Momentum-Balance Differential Equation (Method 3).....	17
3. Method 4.....	18
Compressibility Factor of Natural Gases.....	18
Partial Derivatives of Compressibility Factor..	20
Isobaric Heat Capacity of Natural Gases.....	21
1. Ideal Isobaric Heat Capacity.....	21
a. Model $C_{p_m}^{\circ}$ 1.....	21
b. Model $C_{p_m}^{\circ}$ 2.....	22
2. Heat Capacity Departure.....	23
The Hydraulic Friction Factor.....	24
1. The Smooth Pipe Region.....	24
2. The Transition Region.....	25
3. The Rough Pipe Region.....	25

	Page
Overall Heat-Transfer Coefficient.....	27
1. Model U ₁	27
2. Model U ₂	27
3. Model U ₃	27
Gas-Film Heat-Transfer Coefficient.....	28
Thermal Conductivity of Gas Mixtures.....	29
Viscosity of Pure and Natural Gases.....	32
Pseudocritical Properties of Natural Gases.....	34
PRELIMINARY MODEL TESTING.....	36
Compressibility Factor of Natural Gases.....	36
Partial Derivatives of Compressibility Factor..	36
Isobaric Heat Capacity of Natural Gases.....	38
1. Ideal Isobaric Heat Capacity.....	38
2. Heat Capacity Departure.....	42
The Hydraulic Friction Factor.....	44
Viscosity of Natural Gases.....	44
DATA AND RESULTS.....	50
Back Pressure Data.....	50
Accuracy of Field Data.....	50
1. Flow Rate.....	50
2. Length of Tubing.....	50
3. Pressure Data.....	53
Results.....	53
DISCUSSION.....	60
Compressibility Factor of Natural Gases.....	60
Partial Derivatives of Compressibility Factor..	60
Isobaric Heat Capacity of Natural Gases.....	60
1. Ideal Isobaric Heat Capacity.....	60
2. Heat Capacity Departure.....	62
The Hydraulic Friction Factor.....	63
Overall Heat-Transfer Coefficient.....	64
1. Stanton Number.....	65
2. Prandtl Number of Natural Gas.....	66
Thermal Conductivity of Gas Mixtures.....	66
Viscosity of Pure and Natural Gases.....	67
Flowing Bottom-Hole Pressure.....	68
Reasons for Deviations of Calculated from	
Measured Bottom-Hole Pressures.....	71
1. The Existence of Two-Phase Flow.....	71
2. The Presence of Water Vapor in the	
Flowing Fluid.....	72

	Page
3. Well Bore Deviation from Vertical.....	72
4. Unreliable Estimates of the Viscosity of Sour Gases.....	73
5. Uncertainty of the Length of Flow String.....	73
6. Variation of the Overall Heat- Transfer Coefficient.....	75
7. Uncertainty of Absolute Roughness of Tubing.....	75
CONCLUSIONS.....	77
RECOMMENDATIONS.....	78
NOMENCLATURE.....	79
REFERENCES.....	85
APPENDIX	
A Development of a Mathematical Model for Steady State Compressible Flow of Gases in Pipes.....	97
B Development of Mathematical Expressions for the Partial Derivatives of Z with Respect to T_r and P_r	106
C Development of a Mathematical Expression for Heat Capacity Departure.....	112
D The Hydraulic Friction Factor in Both the Transition and the Rough Pipe Regions..	118
E An Algorithm for Solving the System of Differential Equations in Main Models AA and BB.....	121
F Source Data.....	126
G Copies of the FORTRAN Computer Programs for Main Models AA and BB.....	134

LIST OF TABLES

Table		Page
1	Comparison of Calculated Z-Factors with Experimental Data	37
2	Quality of Fit	43
3	Comparison of Equation 37 and Various ($C_p - C_p^\circ$) Generalized Correlations with the Experimental Data of Krase and Mackey for Nitrogen	45
4	Comparison of Equation 37 and Various ($C_p - C_p^\circ$) Generalized Correlations with the Experimental Data of Workman for Nitrogen and Oxygen	46
5	Comparison of Equation 37 and Other ($C_p - C_p^\circ$) Generalized Correlations with the Weiss and Joffe Correlation	47
6	Comparison of ($C_p - C_p^\circ$) Estimated by Equation 37 with the Sage, et al. Experimental Values for Two Natural Gases	48
7	Comparison Between the Moody and Cullender and Smith Friction Factors	49
8	Summary of Field Data	51
9	Pressure and Temperature Data	52
10	Accuracy of Pressure Measuring Devices	53
11	Accuracy of Pressure Data	54
12	Comparisons Between the Measured and the Calculated Flowing Bottom-Hole Pressures	56
13	Deviations of the Calculated Flowing Bottom-Hole Pressures from the Measured Values	57
14	Deviations of the Calculated Pressure Drops from the Measured Values	58
15	A Comparison of Calculated Bottom-Hole Pressures and Computational Costs for Main Models AA and BB	59

LIST OF FIGURES

Figure		Page
1	Ideal Isobaric Heat Capacity of Natural Gases.....	39
2	The Slopes and Intercepts in Equation 32 as Functions of Gas Gravity.....	40
3	Deviation of Pressure Drops Calculated by Method 2 from Measured Values for All Test Points.....	69
4	Fluid Temperature Profiles Calculated by Methods 1, 2, 3, and 4 for Test-Point C....	70
5	Plan View of 14 Vertical Wells Drilled to 6000 Feet (166).....	74

INTRODUCTION

Flowing bottom-hole pressure is of importance in studying reservoir behavior. Many methods have been developed for the calculation of flowing bottom-hole pressure from surface measurements for single-phase gas flow (1,2,3,4,5,6,7,8). These methods use the Force-Momentum-Balance differential equation and employ various simplifying assumptions.

In 1964, Dranchuk and Quon (9) proposed a general solution for the equations describing steady-state turbulent single-phase gas flow in circular conduits. This method has an advantage over the others in that it permits the simultaneous calculation of the pressure and temperature profiles in the flowing fluid.

Inspection of the Dranchuk and Quon method shows that it employs a constant step-size integrating scheme and approximate formulae for evaluating the various required terms such as: the gas compressibility factor, its partial derivatives with respect to reduced pressure and temperature, the Prandtl number, the gas-film heat-transfer coefficient, and friction factor. Furthermore it handles only the flow of sweet natural gas and requires that the gas composition be known.

The objectives of this study were to

1. Modify the Dranchuk and Quon mathematical model so as to

permit the estimation of absolute roughness when flowing bottom-hole pressure is provided in addition to various surface measurements.

2. Extend the applicability of both the original and modified models so as to handle sour natural gas and flow in the region of partially developed turbulence.
3. Study, test if possible, and choose more sophisticated computer adaptable methods for the evaluation of the gas compressibility factor, its partial derivatives, isobaric heat capacity, friction factor, overall heat-transfer coefficient, gas-film heat-transfer coefficient, Prandtl number, thermal conductivity, and viscosity.
4. Introduce necessary modifications to both the original and modified models so as to permit their use to flow problems where the gas gravity is the only known property of the gas and where gas composition is not available.
5. Compare the calculated bottom-hole pressure values by means of the extended model with field data.

LITERATURE REVIEW

Flowing Bottom-Hole Pressure

A variety of methods has been developed for calculating flowing bottom-hole pressure in dry gas wells (10,11,12, 13,14,15,16,17). Aziz (18) has reviewed most of these methods. These methods utilize the Force-Momentum-Balance differential equation and employ simplifying assumptions such as neglecting the kinetic energy term, assuming constant fluid temperature, etc. Dranchuk and Quon (19) suggested an alternative approach by simultaneously solving the Force-Momentum-Balance and the Total Energy Balance differential equations. Young (20) assumed a linear fluid temperature profile in order to numerically integrate the Force-Momentum-Balance differential equation. Recently Dranchuk and McFarland (21) have studied Young's method and as a result have proposed a modified solution for the same differential equation. The latter two methods assume a constant friction factor as suggested by Cullender and Smith (22).

Compressibility Factor of Natural Gases

The gas compressibility factor is a coefficient which corrects for the non-ideal behaviour of real gases. This factor is necessary in any calculation involving volumes of gaseous mixtures. Standing and Katz (23) presented a graphi-

cal correlation for the compressibility factor of dry and sweet natural gases. Elfrink, et al. (24) proposed another graphical correlation for the same factor for sweet and two-phase as well as one-phase natural gases. The former correlation has been widely accepted by the petroleum industry. Accordingly several fitting and interpolation techniques (25,26,27) have sought a mathematical representation for the Standing and Katz correlation.

In the last 35 years many equations of state have been developed to handle hydrocarbon gases and their mixtures (28,29,30,31). Wichert and Aziz (32,33) presented a brief review of the various equations of state used for calculating natural gas compressibility factors. Among these equations those of Redlich-Kwong (34) and Benedict-Webb-Rubin (35) have received wide interest. Several investigators (36,37,38) have introduced a few modifications to the B-W-R equation to make it applicable at very low temperatures and useful in handling complex hydrocarbon mixtures. Others (39,40,41) have fitted the generalized forms of various equations of state to the tabulated Standing and Katz correlation (42) in attempts to predict the volumetric behavior of sweet natural gases.

In order to predict the volumetric behavior of dry sour natural gases, Robinson, et al. (43) developed a correction factor correlation to be used with the Standing and Katz correlation. Wichert and Aziz (44,45,46) presented an alternative approach which employs modified pseudocriticals

of sour natural gas.

Isobaric Heat Capacity of Natural Gases

The isobaric heat capacity of gases is usually calculated by evaluating the ideal isobaric heat capacity and the heat capacity departure from the ideal state at high pressures.

Kobe, et al. (47) and Thinh, et al. (48) presented the ideal isobaric heat capacity of pure gases as a function of temperature by means of third degree polynomials. Kothari and Doraiswamy (49), through dimensional analysis, showed that the ideal isobaric heat capacity for pure gases could be expressed as a function of molecular weight and reduced temperature. On the other hand the ideal isobaric heat capacity of a gas mixture may be expressed by means of the molal average (50).

The heat capacity departure from the ideal state at high pressures has been studied by many investigators. Dodge (51) developed a generalized correlation for the ratio $\frac{C_p}{C_p^\circ}$. Watson and Smith (52), Lydersen, et al. (53), Edmister (54), Weiss and Joffe (55), and Sherwood (56) presented either graphical or tabulated generalized correlations for the difference $(C_p - C_p^\circ)$. Comings and Nathan (57) showed that the difference $(C_p - C_p^\circ)$ and not the ratio $\frac{C_p}{C_p^\circ}$ is a function of reduced pressure and temperature. As a result, they recommended the use of the difference $(C_p - C_p^\circ)$ as a generalized correlation. Furthermore they pointed out that the

Watson and Smith generalized correlation might give errors as high as 50 per cent. The Edmister generalized correlation may be in error by as much as 50 to 100 per cent (58). Sledjeski (59), Seifarth and Joffe (60) used the B-W-R equation of state for predicting ($\bar{C}_p = 20^\circ$) of propane and methane. They reported that satisfactory agreement with experimental data was obtained for temperatures greater than 150°F.

The Hydraulic Friction Factor

Under normal conditions of fluid flow in pipe (normal velocities, single-phase flow, confined conduits, no temperature difference between the fluid and the boundary wall) the friction factor is a function of Reynolds number and relative roughness (61,62). For commercial pipes, the friction factor decreases at a decreasing rate as Reynolds number increases, and once the boundary Reynolds number (Re_b) has been reached, the friction factor remains constant (63,64).

Several investigators (65,66) have developed different mathematical expressions for the friction factor in both the rough and smooth pipe regions. The friction factor in the transition region was controversial until Colebrook (67) bridged the gap between the smooth and rough pipe regions with his particular transition law. Later, Moody (68) constructed a friction factor chart in which he employed the Von Kármán formulae for the smooth and rough pipe regions and the Colebrook transition law for the transition region.

Cullender and Smith (69) assumed that gas flow through

tubing in gas wells is fully turbulent. Consequently they suggested a constant friction factor as a function of tubing diameter only.

Absolute Roughness of Gas Well Tubing

It is well known that absolute roughness changes with time due to pipe corrosion. Colebrook (70) found that the rate of deterioration of pipe depends on time, type of flow, lining material, size of pipe, and the character of the flowing fluid.

Cullender and Binckley (71) reported an absolute roughness of 0.0006 inch for gas well casing. Smith, et al. (72), from their own experimental data, reported a value of 0.00065 inch for the absolute roughness of gas well tubing. A range of absolute roughness of 0.00055-0.0019 inch for experimental gas pipe lines has been reported by others (73).

Overall Heat-Transfer Coefficient

Heat-transfer between a flowing fluid in a pipe and a solid separated by one or more solid layers can be conveniently handled by utilizing the overall heat-transfer coefficient concept. Expressions for the overall heat-transfer coefficient for uniform geometries may be found in any heat-transfer reference (74,75). Holst (76) studied this coefficient for the case of steam injection down a well assuming simple and uniform geometries. Dranchuk and Quon (77) suggested an expression which requires a trial-and-error

procedure for the same coefficient in a gas well system.

The various expressions for the overall heat-transfer coefficient require knowledge of the gas-film heat-transfer coefficient. The latter coefficient has been exhaustively investigated. Simple, semi-empirical, empirical and sophisticated models have been developed for the estimation of the gas-film heat-transfer coefficient (78,79,80,81,82,83).

Thermal Conductivity of Gas Mixtures

Reid and Sherwood (84) have reviewed most of the existing methods for the estimation of the thermal conductivity of pure gases and their mixtures. Among these methods, that of Misic and Thodos (85,86) for pure gases may be efficiently adapted to computer calculations. Lindsay and Bromley (87) proposed a generalized equation for the thermal conductivity of gas mixtures at atmospheric pressure. Stiel and Thodos (88) studied the effect of pressure on the thermal conductivity of pure gases.

Viscosity of Pure and Natural Gases

The viscosity of pure gases and their mixtures has been studied by numerous investigators. Reid and Sherwood (89) gave an excellent review on the subject. Herning and Zipperer (90) suggested an empirical formula for estimating the viscosity of a gas mixture from those of the individual components. Bicher and Katz (91) presented a graphical method for natural gas viscosity as a function of specific gravity.

temperature, and pressure. Carr, et al. (92) developed another graphical correlation, which utilized the law of corresponding state. Lee, et al. (93,94) studied and developed empirical formulae for the viscosity of sweet natural gases as well as pure hydrocarbons.

MODEL DEVELOPMENT

Four main models were developed to meet part of the objectives of this study. The objective, development, and requirements of each main model are as follows:

Main Model A

The objective of this model is to estimate the absolute roughness of tubing.

The development of this model is presented in detail in Appendix A. It can be briefly described as follows:

- The Force-Momentum-Balance differential equation for a fluid element contained in a pipe segment may be written as:

$$\frac{V}{M_w} dP + \frac{g}{g_c} dz + \frac{\langle u \rangle d\langle u \rangle}{\alpha g_c} + \frac{f \langle u \rangle^2}{2 g_c D} dL = 0 \quad (1)$$

where α is the inverse of the momentum correction factor.

(The value of the momentum correction factor for turbulent flow is closer to unity than for laminar flow (95).)

- The Total Energy Balance differential equation over the same pipe segment may be expressed as:

$$\delta q - \frac{dW}{J} = \frac{\pi D^2}{4} \left(\frac{G}{M_w} \right) \left[dh + \frac{\langle u \rangle d\langle u \rangle}{\alpha J g_c} + \frac{g}{g_c J} dz \right] \quad (2)$$

- The equation of state:

$$PV = ZRT \quad (3)$$

- Equations 1, 2, and 3 were rearranged to yield the following:

$$A^* \frac{dP_r}{dL} + B^* \frac{dT_r}{dL} = C^* \quad (4)$$

and

$$D^* \frac{dP_r}{dL} + E^* \frac{dT_r}{dL} = F^* \quad (5)$$

From which

$$\frac{dP_r}{dL} = \frac{C^* \cdot E^* - B^* \cdot F^*}{A^* \cdot E^* - B^* \cdot D^*} \quad (6)$$

and

$$\frac{dT_r}{dL} = \frac{A^* \cdot F^* - C^* \cdot D^*}{A^* \cdot E^* - B^* \cdot D^*} \quad (7)$$

where

$$\begin{aligned} A^* &= fcn(T_r, P_r, T_c, P_c, M_w, Q_0, D) \\ B^* &= fcn(T_r, P_r, T_c, P_c, M_w, Q_0, D) \\ C^* &= fcn(T_r, P_r, T_c, P_c, M_w, Q_0, D, \frac{E}{D}, S) \\ D^* &= fcn(T_r, P_r, T_c, P_c, M_w, Q_0, D) \\ E^* &= fcn(T_r, P_r, T_c, P_c, M_w, Q_0, D) \\ F^* &= fcn(T_r, P_r, T_c, P_c, M_w, Q_0, D, U) \end{aligned} \quad (8)$$

and

In order to obtain a solution, $\frac{dP_r}{dL}$ and $\frac{dT_r}{dL}$ must be finite.

Equations 6, 7, and 8 imply that

$$\frac{dP_r}{dL} = \text{fcn}(T_r, P_r, T_c, P_c, M_w, Q_o, D, \frac{\epsilon}{d}, S, U) \quad (9)$$

and

$$\frac{dT_r}{dL} = \text{fcn}(T_r, P_r, T_c, P_c, M_w, Q_o, D, \frac{\epsilon}{d}, S, U) \quad (10)$$

- For a specific flow problem, that is, when pipe dimensions and position in space (L^*, D, S), flow rate (Q_o), and gas properties (T_c, P_c, M_w) are given; Equations 9 and 10 reduce to

$$\frac{dP_r}{dL} = \text{fcn}(T_r, P_r, \frac{\epsilon}{d}, U) \quad (11)$$

and

$$\frac{dT_r}{dL} = \text{fcn}(T_r, P_r, \frac{\epsilon}{d}, U) \quad (12)$$

- The overall heat-transfer coefficient, U , as suggested by Dranchuk and Quon (96) is:

$$\frac{1}{U} = \frac{1}{h_o} + R^* \quad (13)$$

where h_o = gas-film heat-transfer coefficient

$$= \text{fon}(\bar{P}_r, R_e, \frac{\epsilon}{d})$$

$$= \text{fon}(T_r, P_r, \frac{\epsilon}{d})$$

(14)

$\bar{P}_r$ = Prandtl number

R_e = Reynolds number

and R^* = average thermal resistance of the surroundings, $\text{ft.}^2 \text{ sec. } ^\circ\text{R/Btu}$

- Combining Equations 11, 12, 13, and 14 yields,

$$\frac{dP_r}{dL} = \text{fon}(T_r, P_r, \frac{\epsilon}{d}, R^*)$$

(15)

and

$$\frac{dT_r}{dL} = \text{fon}(T_r, P_r, \frac{\epsilon}{d}, R^*)$$

(16)

- It is assumed that $\frac{\epsilon}{d}$ and R^* are constants but have unknown values which have to be determined.

Thus the main model A, which describes single-phase gas flow in pipes, consists of two non-linear differential equations, Equations 6 and 7 or 15 and 16, and two-parameter estimations namely $\frac{\epsilon}{d}$ and R^* . This system of differential equations has to be solved simultaneously for the boundary conditions:

$$\text{At } L = 0 \quad P(0) = P_1 \quad \text{and} \quad T(0) = T_1$$

$$\text{At } L = L^* \quad P(L^*) = P_2 \quad \text{and} \quad T(L^*) = T_2$$

This model requires algebraic expressions for the gas compressibility factor, its partial derivatives, friction factor, the Prandtl number, gas-film heat-transfer coefficient, isobaric heat capacity, viscosity and thermal conductivity.

Data required are: gas composition, flow rate, tubing head pressure and temperature, bottom-hole pressure and temperature, length and inner diameter of tubing, and ground surface temperature.

Main Model AA (Method 2)

The objective of this model is to calculate the flowing bottom-hole pressure.

This model is the original Dranchuk and Quon model. It is a special case of main model A where $\frac{\epsilon}{d}$ is assumed to be $\frac{0.0006}{d}$ as suggested by Cullender and Smith (97). Thus, this model consists of two non-linear differential equations, Equations 6 and 7, and one parameter estimation namely R^* . This system of differential equations has to be solved simultaneously for the boundary conditions:

$$\text{At } L = 0 \quad T(0) = T_1$$

$$\text{At } L = L^* \quad P(L^*) = P_2 \quad \text{and} \quad T(L^*) = T_2$$

This model requires algebraic expressions for the gas compressibility factor, its partial derivatives, friction factor, the Prandtl number, gas-film heat-transfer coefficient, isobaric heat capacity, viscosity and thermal conductivity.

vity,

Data required are: gas composition, flow rate, tubing head pressure and temperature, bottom-hole temperature, length and inner diameter of tubing, and ground surface temperature.

Main Model B

The objective of this model is to estimate the absolute roughness of tubing.

This model differs from main model A in that it assumes the overall heat-transfer coefficient constant at some average value; hence

$$\frac{1}{U} \approx R^{**} = \text{constant} \quad (17)$$

Consequently when Equations 11, 12, and 17 are combined,

$$\frac{dP_r}{dL} = \text{fon}(T_r, P_r, \frac{\epsilon}{d}, R^{**}) \quad (18)$$

and

$$\frac{dT_r}{dL} = \text{fon}(T_r, P_r, \frac{\epsilon}{d}, R^{**}) \quad (19)$$

The same boundary conditions as those in main model A are employed.

This model requires algebraic expressions for the gas

compressibility factor, its partial derivatives, friction factor, isobaric heat capacity, and viscosity.

Data required are: flow rate, tubing head pressure and temperature, bottom-hole pressure and temperature, length and inner diameter of tubing, ground surface temperature, and either the gas composition or the gas gravity along with the mole fractions of N_2 , CO_2 , and H_2S .

Main Model BB (Method 2)

The objective of this model is to calculate the flowing bottom-hole pressure.

This model is a special case of main model B where $\frac{\epsilon}{d}$ is assumed to be $\frac{0.0006}{d}$. Thus this model consists of two non-linear differential equations, Equations 6 and 7, and one parameter estimation namely R^{**} .

The same boundary conditions as those in main model AA are employed.

This model requires algebraic expressions for the gas compressibility factor, its partial derivatives, friction factor, isobaric heat capacity, and viscosity.

Data required are: flow rate, tubing head pressure and temperature, bottom-hole temperature, length and inner diameter of tubing, ground surface temperature, and either the gas composition or the gas gravity along with the mole fractions of N_2 , CO_2 , and H_2S .

Special Cases of Main Models AA and BB

In the following special cases both main models AA and BB reduce to one model:

1. Adiabatic Flow (Method 1):

There is no heat-transfer between the flowing fluid and the surrounding; that is,

$$\dot{U} = 0 \quad (20)$$

This method requires solving simultaneously Equations 6 and 7 for the boundary conditions.

$$\text{At } L = 0 \quad T(0) = T_1$$

and

$$\text{At } L = L^* \quad P(L^*) = P_2$$

2. Complete Solution of the Force-Momentum-Balance Differential Equation (Method 3):

This method assumes a linear temperature profile in the flowing fluid. It requires solving Equation 4 for the boundary condition:

$$L = L^* \quad P(L^*) = P_2$$

Equation 4 may be written as:

$$\frac{dP_r}{dL} = (C^* - B^* \frac{dT_r}{dL}) / D^* \quad (21)$$

where

$$\frac{dT_r}{dL} = \left(\frac{T(L^*) - T(0)}{L^*} \right) / p T_c \quad (22)$$

3. Method 4:

This method is a special case of Method 3 where the friction factor is assumed constant as suggested by Cullender and Smith (98).

Method 4 is equivalent to that of Dranchuk and McFarland (99).

The main models require algebraic expressions (models) for the gas compressibility factor, its partial derivatives with respect to reduced temperature and pressure, isobaric heat capacity, friction factor, overall heat-transfer coefficient, gas-film heat-transfer coefficient, the Prandtl number, thermal conductivity, and viscosity.

The various mathematical expressions for these terms are the subject of the rest of this section.

Compressibility Factor of Natural Gases

The compressibility factor of dry sweet natural gas may be expressed analytically by many methods. In this study the reduced form of the B-W-R equation of state as suggested by Purvis (100) and Dranchuk, et al. (101) was chosen to represent the compressibility factor, where

$$Z = \frac{Z_c P_r}{\rho_r T_r} = 1 + (A_1 + \frac{A_2}{T_r} + \frac{A_3}{T_r^3}) \rho_r + (A_4 + \frac{A_5}{T_r}) \rho_r^2 + \frac{A_5 A_6}{T_r} \rho_r^5 + \frac{A_7}{T_r^3} \rho_r^2 (1 + A_8 \rho_r^2) \exp(-A_8 \rho_r^2) \quad (23)$$

and

$$\begin{aligned} A_1 &= 0.31506237 & A_2 &= -1.0467099 & A_3 &= -0.57832727 \\ A_4 &= 0.53530771 & A_5 &= -0.61232032 & A_6 &= -0.10488813 \\ A_7 &= 0.68157001 & A_8 &= 0.68446549 \end{aligned}$$

Equation 23 may be extended to sour gas by applying modified pseudocritical pressure and temperature, when calculating the pseudoreduced pressure and temperature, as suggested by Wichert and Aziz (102), where

$$\begin{aligned} \epsilon_3 &= 120 \left[(Y_{CO_2} + Y_{H_2S})^{0.9} - (Y_{CO_2} + Y_{H_2S})^{1.6} \right] \\ &+ 15 \left[Y_{H_2S}^{0.5} - Y_{H_2S}^4 \right] \end{aligned} \quad (24)$$

$$T'_c = T_c - \epsilon_3$$

and

$$P'_c = \frac{P_c T'_c}{\left[T_c + Y_{H_2S} (1 - Y_{H_2S}) \epsilon_3 \right]}$$

Partial Derivatives of Compressibility Factor

The partial derivatives of Z with respect to reduced temperature and pressure may be calculated by employing a numerical technique. However mathematical expressions for these derivatives may be derived from Equation 23 as shown in Appendix B. These mathematical expressions may be expressed as:

$$\left(\frac{\partial Z}{\partial T_r}\right)_{P_r} = Z \left[-\frac{1}{T_r} - \frac{1}{\rho_r} \left(\frac{\partial \rho_r}{\partial T_r}\right)_{P_r} \right] \quad (25)$$

and

$$\left(\frac{\partial Z}{\partial P_r}\right)_{T_r} = Z \left[\frac{1}{P_r} - \frac{1}{\rho_r} \left(\frac{\partial \rho_r}{\partial P_r}\right)_{T_r} \right] \quad (26)$$

The partial derivatives of reduced density with respect to reduced temperature and pressure, $\left(\frac{\partial \rho_r}{\partial T_r}\right)_{P_r}$ and $\left(\frac{\partial \rho_r}{\partial P_r}\right)_{T_r}$, may be written as:

$$\left(\frac{\partial \rho_r}{\partial T_r}\right)_{P_r} = \frac{- \left[\rho_r + \left(A_1 - \frac{2A_3}{T_r}\right) \rho_r^2 + A_4 \rho_r^3 - \frac{2A_7}{T_r} \rho_r^3 (1 + A_8 \rho_r^2) \exp(-A_8 \rho_r^2) \right]}{\left[T_r + 2(A_1 T_r + A_2 + \frac{A_3}{T_r}) \rho_r + 3(A_4 T_r + A_5) \rho_r^2 + 6A_5 A_6 \rho_r^5 + \frac{A_7}{T_r} \rho_r^2 (3 + 3A_8 \rho_r^2 - 2A_8^2 \rho_r^4) \exp(-A_8 \rho_r^2) \right]} \quad (27)$$

and

$$\left(\frac{\partial \rho_r}{\partial P_r}\right)_{T_r} = \frac{Z_c}{\left[T_r + 2\left(A_1 T_r + A_2 + \frac{A_3}{T_r}\right)\rho_r + 3\left(A_4 T_r + A_5\right)\rho_r^2 + 6A_5 A_6 \rho_r^5 + \frac{A_7}{T_r^2} \rho_r^2 \left(3 + 3A_8 \rho_r^2 - 2A_8^2 \rho_r^4\right) \exp(-A_8 \rho_r^2) \right]} \quad (28)$$

Isobaric Heat Capacity of Natural Gases

The isobaric heat capacity of natural gas may be represented by:

$$C_{p_m} = C_{p_m}^{\circ} + (C_{p_m} - C_{p_m}^{\circ}) \quad (29)$$

where C_{p_m} = isobaric heat capacity of natural gases, Btu/lb mole. °R.

$C_{p_m}^{\circ}$ = ideal isobaric heat capacity of natural gases, Btu/lb mole. °R.

$(C_{p_m} - C_{p_m}^{\circ})$ = heat capacity departure from the ideal state of natural gases, Btu/lb mole. °R.

1. Ideal Isobaric Heat Capacity:

Two models may be used to represent the ideal heat capacity of natural gas.

a. Model $C_{p_m}^{\circ}$ 1:

In this model, $C_{p_m}^{\circ}$ is taken as the molal average; that is,

$$C_{p_m}^{\circ} = \sum_{i=1}^n C_{p_i}^{\circ} Y_i \quad (30)$$

where

$$C_{p_i}^{\circ} = \sum_{j=0}^k a_j T^j \quad (31)$$

This model requires that the gas composition be known.

b. Model $C_{p_m}^{\circ}$ 2:

This model was developed in this study to represent the ideal isobaric heat capacity of sweet natural gas as a function of gas gravity (G) and pseudoreduced temperature (T_r); that is,

$$C_{p_{m_{H.C.}}}^{\circ} = A(G) + B(G)T_r \quad (32)$$

where the subscript $m_{H.C.}$ means a mixture of hydrocarbons. This equation can be extended to sour natural gas by applying corrections to account for the presence of N_2 , CO_2 , and H_2S as follows:

$$C_{p_m}^{\circ} = \frac{C_{p_{m_{H.C.}}}^{\circ}}{F_{CO_2} \cdot F_{N_2} \cdot F_{H_2S}} \quad (33)$$

$$F_{N_2} = fcn(Y_{N_2}, T) \quad (34)$$

$$F_{CO_2} = f_{on}(Y_{CO_2}, T) \quad (35)$$

and

$$F_{H_2S} = f_{on}(Y_{H_2S}, T) \quad (36)$$

where T = temperature of gas, °R.

2. Heat Capacity Departure:

A heat capacity departure model was developed by employing the reduced form of the B-W-R equation of state and utilizing thermodynamic properties as shown in Appendix C. The heat capacity departure may be expressed as:

$$\begin{aligned} \frac{C_{P_m} - C_{P_m}^0}{R} = & \left[-1 - \frac{6A_3}{T_r^3} \rho_r - \frac{6A_7}{A_8 T_r^3} + \left(\frac{6A_7}{A_8 T_r^3} + \frac{3A_7}{T_r^3} \rho_r^2 \right) \exp(-A_8 \rho_r^2) \right] \\ & + \frac{T_r \left[1 + \left(A_1 - \frac{2A_3}{T_r^3} \right) \rho_r + A_4 \rho_r^2 - \frac{2A_7}{T_r^3} \rho_r^2 (1 + A_8 \rho_r^2) \exp(-A_8 \rho_r^2) \right]^2}{\left[T_r + 2(A_1 T_r + A_2 + \frac{A_3}{T_r^2}) \rho_r + 3(A_4 T_r + A_5) \rho_r^2 + \right.} \\ & \left. 6A_5 A_6 \rho_r^5 + \frac{A_7}{T_r^2} \rho_r^2 (3 + 3A_8 \rho_r^2 - 2A_8^2 \rho_r^4) \exp(-A_8 \rho_r^2) \right] \quad (37) \end{aligned}$$

Although this equation was developed for sweet natural

gases, it may be applied to sour natural gases as will be shown later.

The Hydraulic Friction Factor

The hydraulic friction factor as presented by Moody (103) in chart form has been widely accepted and used for estimating pressure drop due to friction in both gas and oil pipe lines. This chart was developed by using the following equations for the various regions:

1. The Smooth Pipe Region:

$$\frac{1}{\sqrt{f}} = 2 \log \left(\frac{R_e \sqrt{f}}{2.51} \right) ; \quad \epsilon^* < 3 \quad (38)$$

where f = friction factor.

R_e = Reynolds number

$$= 20014.855 \frac{Q_0 G}{\mu d} \text{ (field units)}$$

G = gas gravity, (air = 1).

$$= \frac{M_w}{M_{air}} = \frac{M_w}{28.97}$$

μ = gas viscosity, c.p.

Q_0 = gas flow rate, MMSCF/Day.

$$\epsilon^* = \sqrt{\frac{f}{8}} R_e \frac{\epsilon}{d}$$

ϵ = absolute roughness, inch.

and d = inner diameter of tubing, inch.

2. The Transition Region:

This region is sometimes referred to as the region of partially developed turbulence. The friction factor in this region was represented by the Colebrook transition law (104), where

$$\frac{1}{\sqrt{f}} = -2 \log \left(\frac{\epsilon/d}{3.7} + \frac{2.51}{Re \sqrt{f}} \right) ; \quad 3 \leq \epsilon^* \leq 60 \quad (39)$$

3. The Rough Pipe Region:

This region is sometimes referred to as the fully developed turbulent region. The Von Kármán rough pipe law was employed to represent the friction factor in this region, where

$$\frac{1}{\sqrt{f}} = 2 \log \left(\frac{3.7}{\epsilon/d} \right) ; \quad \epsilon^* > 60 \quad (40)$$

Rather than using ϵ^* to express the boundary between the transition and rough pipe regions one may use the boundary Reynolds number (Re_b) technique proposed by Moody (105), where

$$Re_b = \frac{3200}{\epsilon/d} \quad (41)$$

This method is simpler and easier to use than ϵ^* and was, therefore, chosen in this study.

Equation 39 is too cumbersome for hand calculation since

it requires trial-and-error. For this reason Cullender and Binckley (106) suggest a simplified friction factor formula as a function of flow rate, viscosity, gas gravity, and pipe diameter. On the other hand Cullender and Smith (107) suggested two friction factor formulae for fully developed turbulent flow (rough pipe region), both of which are functions of pipe diameter only. Both groups assumed an absolute roughness of 0.0006 inch in the development of their approximate formulae.

Moody's chart was employed (in this investigation) to represent the hydraulic friction factor for the following reasons:

1. Main models A and B require that the friction factor be a function of relative roughness in order to give an estimate of absolute roughness of the tubing.
2. Flow of natural gas in gas wells is usually turbulent and falls in either the transition or rough pipe region. However, during "Back Pressure Tests", flow is more likely to be in the transition region, especially at low flow rates. Therefore, the Cullender and Smith friction factor equations may not hold.
3. Moody's chart has been accepted as applicable to pipes with relative roughness of 0.0 - 0.05, while the other formulae are applicable only for pipes with absolute roughness of 0.0006 inch.

Overall Heat-Transfer Coefficient

The heat-transfer from a flowing fluid through pipe to the surroundings over a pipe segment (dL) may be expressed as:

$$- \delta q = - \pi D dL U(T_s - T) \quad (43)$$

For a composite cylindrical wall, the overall heat-transfer coefficient (U) may be represented by three models:

1. Model U_1 :

This model can be found in any heat-transfer reference (108). The model that applies to gas wells may be written as:

$$\frac{1}{U} = \frac{1}{h_o} + \frac{D_o}{2} \sum_{i=1}^n \frac{\ln(D_i/D_{i-1})}{k_{i-1,i}} \quad (44)$$

2. Model U_2 :

This model was suggested by Dranchuk and Quon (109) for gas well systems. It assumes a constant thermal resistivity of the surroundings, where

$$\frac{1}{U} = \frac{1}{h_o} + R^* \quad (45)$$

3. Model U_3 :

This model is proposed in this study for gas well systems. It assumes that the overall heat-transfer coefficient

ent is constant at some average value; that is,

$$\frac{1}{U} = R^{**} = \text{constant} \quad (46)$$

Gas-Film Heat-Transfer Coefficient

As can be seen from the previous section, models U_1 and U_2 require an estimate for the gas-film heat-transfer coefficient, h_o . Such an estimate, in turn, requires an estimate of the Stanton number.

Many investigators have developed various models for the Stanton number, which is related to the gas-film heat-transfer coefficient by (110):

$$S_t = \frac{h_o}{\rho \langle u \rangle c_p} \quad (47)$$

where S_t = Stanton number.

ρ = gas density, lbm/ft³.

$\langle u \rangle$ = mean gas velocity, ft/sec.

and c_p = isobaric heat capacity, Btu/lbm.°R.

The Gowen and Smith (111) model is considered the most sophisticated in addition to being the best of all in reproducing experimental data. This model may be expressed as:

$$S_t = \frac{\sqrt{\frac{f}{8}}}{X - 4.5} \quad (48)$$

where

$$X = \bar{P}_r^{0.5} \left[0.155 \left(\frac{f}{8} \right)^{0.27} R_e^{0.54} + \sqrt{\frac{8}{7}} \right] \quad (49)$$

R_e = Reynolds number.

$\bar{P}_r$ = Prandtl number.

$$= \frac{C_p \mu}{M_w k} \quad (50)$$

f = friction factor.

k = thermal conductivity, cal/cm. sec. K.

μ = viscosity, gm/sec. cm or poise.

C_p = isobaric heat capacity, cal/gm mole. K.

M_w = molecular weight, gm/gm mole.

Thermal Conductivity of Gas Mixtures

Thermal conductivity is required for the calculation of Prandtl number, which, in turn, is necessary for the calculation of the Stanton number. Thermal conductivity of gas mixtures may be expressed as:

$$k_m = k_m^{\circ} + (k_m - k_m^{\circ}) \quad (51)$$

where

k_m = thermal conductivity of gas mixtures.

k_m° = thermal conductivity of gas mixtures at one atmosphere.

$(k_m - k_m^{\circ})$ = thermal conductivity departure of gas mixtures.

The Lindsay and Bromley generalized equation (112) can be used for the estimation of the thermal conductivity of gas mixtures at one atmosphere, where

$$k_m = \sum_{i=1}^n \frac{k_i^{\circ}}{\left[1 + \sum_{\substack{j=1 \\ j \neq i}}^n A_{ij} \frac{Y_j}{Y_i} \right]} \quad (52)$$

$$A_{ij} = 0.25 \left\{ 1 + \left[\frac{\mu_j^{\circ}}{\mu_i^{\circ}} \left(\frac{M_{wj}}{M_{wi}} \right)^{0.75} \frac{(1 + \frac{S_i}{T})}{(1 + \frac{S_j}{T})} \right]^{0.5} \right\}^2 \left[\frac{\frac{S_{ij}}{1 + \frac{S_j}{T}}}{\frac{S_i}{1 + \frac{S_i}{T}}} \right] \quad (53)$$

$$S_i = 1.5 T_{bi} \quad (54)$$

$$S_{ij} = \sqrt{S_i S_j} \quad (55)$$

S_i = the Sutherland constant for i th component, K.

μ_i° = low pressure viscosity of i th component, c.p.

M_{wi} = molecular weight of i th component, gm/gm mole

T_{bi} = normal boiling point temperature of i th component, K.

and k_i° = thermal conductivity of i th component at one atmosphere, cal/sec. cm, K.

The thermal conductivity of pure components at one

atmosphere can be represented as suggested by Misic and Thodos (113, 114) where

For Hydrocarbons:

$$k_i^o = \frac{C_{p_i}^o}{10^6 \gamma_i} (14.52 T_{r_i} - 5.14)^{\frac{2}{3}} ; T_{r_i} \geq 1 \quad (56)$$

For Non-hydrocarbons:

$$k_i^o = \frac{C_{p_i}^o}{10^6 \gamma_i} \left[(195 Z_{c_i} - 31.94) T_{r_i} + (16.83 - 82.5 Z_{c_i}) \right]^{(1.524 - 2.8 Z_{c_i})} ;$$

$$1 \leq T_{r_i} \leq 3 \quad (57)$$

and

$$\gamma_i = \frac{T_{c_i}^{\frac{1}{6}} M_{w_i}^{\frac{1}{2}}}{P_{c_i}^{\frac{2}{3}}} \quad (58)$$

with T_{c_i} and P_{c_i} in degrees Kelvin and atmospheres respectively.

The thermal conductivity departure can be represented as suggested by Stiel and Thodos (115), where

$$(k_m - k_m^0) = \frac{14.0 \left[\exp(0.534 \rho_r) - 1 \right]}{10^8 \gamma Z_c}; \quad \rho_r \leq 0.5 \quad (59)$$

$$= \frac{13.1 \left[\exp(0.67 \rho_r) - 1.069 \right]}{10^8 \gamma Z_c}; \quad 0.5 < \rho_r \leq 2.0 \quad (60)$$

$$= \frac{2.976 \left[\exp(1.155 \rho_r) - 2.016 \right]}{10^8 \gamma Z_c}; \quad 2.0 < \rho_r \leq 2.8 \quad (61)$$

Viscosity of Pure and Natural Gases

The viscosity of natural gases is required for the calculation of the Reynolds and Prandtl numbers. It can be evaluated by means of the Lee, et al. correlation (116), where

$$\mu_m = K \exp(X \rho^Y)$$

$$K = \frac{(9.4 + 0.02 M_w) T^{1.5}}{(209 + 19 M_w + T)}$$

(62)

$$X = 3.5 + \frac{986}{T} + 0.01 M_w$$

$$Y = 2.4 - 0.2 X$$

where μ_m = natural gas viscosity, micropoise.

ρ = natural gas density, gm/cm³.

T = temperature, °R.

and M_w = average molecular weight, gm/gm mole.

They recommend that the density be calculated using the Standing and Katz compressibility factor correlation.

The viscosity of pure components at low pressure is needed for the calculation of thermal conductivity at low pressure. It can be expressed by using the modified rigid interacting sphere model (117), where

$$\mu_i^o = 0.0026693 \frac{\sqrt{M_{wi} T}}{\sigma_i^2 \Omega_{vi}} \quad (63)$$

$$\frac{1}{\Omega_{vi}} = 0.697 (1 + 0.323 \ln T_i^*) \quad (118) \quad (64)$$

μ_i^o = viscosity of i th component, c.p.

$$T_i^* = \frac{T}{\left(\frac{\epsilon_0}{K}\right)_i}$$

T = temperature, K.

ϵ_0 = maximum energy of attraction for Lennard-Jones potential, ergs.

K = Boltzmann constant, 1.3805×10^{-16} ergs/ K.

and σ = collision diameter for Lennard-Jones potential, Å.

Values of σ and $\left(\frac{\epsilon_0}{K}\right)$ are available in the literature (119,120). Alternatively, the Stiel and Thodos empirical

equations (121) may be used.

If $T_{r_i} < 1.5$ or 2.0 , Reid and Sherwood (122) recommend the modification of the value of μ_i^o according to the following relation:

$$\mu_{1i}^o(T_{r_i} < 2) = \mu_{2i}^o(T_{r_i} = 2) \frac{\left(\frac{T_1^*}{T_2^*}\right)_i}{\left(\frac{\Omega_{v1}}{\Omega_{v2}}\right)_i} \quad (65)$$

where μ_2^o is chosen arbitrarily to correspond to that at $T_{r_i} = 2$.

Pseudocritical Properties of Natural Gases

The gas compressibility factor, its partial derivatives, heat capacity and thermal conductivity departure models employ the law of corresponding states. Consequently an estimate of the pseudocritical properties is necessary to calculate the reduced pressure and temperature. These pseudocriticals may be taken as the molal average as suggested by Kay (123) when the gas composition is available. However when the gas composition is not available and where the gas gravity (G) is the only known property of the gas, the pseudocriticals may be estimated from the Brown, et al. correlation (124). Flores (125) presented the correlation in an algebraic form as

$$p^T_{cH.C.} = 171.137 + 313.725 G \quad (66)$$

$$p^P_{cH.C.} = 695.100 - 40.000 G ; \quad G \leq 0.85 \quad (67)$$

$$= 704.396 - 51.724 G ; \quad G > 0.85 \quad (68)$$

Carr, et al. (126) have introduced corrections in order to make this correlation applicable to sour natural gases if the mole fractions of N_2 , CO_2 and H_2S are known. Hence the pseudocriticals of sour natural gases may be estimated from:

$$p^T_c = p^T_{cH.C.} - 250 Y_{N_2} - 80 Y_{CO_2} + 130 Y_{H_2S} \quad (69)$$

and

$$p^P_c = p^P_{cH.C.} - 170 Y_{N_2} + 440 Y_{CO_2} + 600 Y_{H_2S} \quad (70)$$

PRELIMINARY MODEL TESTING

The various algebraic expressions selected in the previous section to represent compressibility factor, its partial derivatives, isobaric heat capacity, friction factor, and viscosity were tested by comparing with experimental data or correlations available in the literature. The results of these tests are presented in this section.

Compressibility Factor of Natural Gases

Table 1 presents a comparison of the compressibility factor model as expressed by Equation 23 with 1350 experimental data points (127,128,129). These experimental data are for sweet as well as sour natural gases with up to 73.85% H_2S and 54.46% CO_2 .

Partial Derivatives of Compressibility Factor

The partial derivatives of the compressibility factor with respect to reduced pressure and temperature were compared to values determined graphically. This was achieved by using experimental compressibility factors (130,131). First, a few isobars and isotherms were plotted, then the slopes of these curves at specific P_r and T_r were calculated. Finally these results were compared to the partial derivatives estimated by the algebraic expressions. It was found

TABLE 1
COMPARISON OF CALCULATED Z-FACTORS
WITH EXPERIMENTAL DATA

Source of Experimental Data	No. of Points	Average Absolute Deviation from Experimental Values
Reference (127)	58	1.030%
Reference (128)	207	1.202%
Reference (129)	1085	0.976%
All Data Points	1350	1.013%

that Equation 25 gives an average absolute deviation of 5.84 per cent when compared with 5 points, while Equation 26 deviates by 1.96 per cent when compared with 16 points.

Isobaric Heat Capacity of Natural Gases

1. Ideal Isobaric Heat Capacity

The isobaric heat capacity of sweet natural gases which have no nitrogen may be represented by Equation 32 as suggested in Model Development. This is an equation of a family of straight lines, the slopes and intercepts of which are functions of the gas gravity.

To test the model represented by Equation 32, the compositions of 72 gases were collected from the literature and the equivalent pure hydrocarbon compositions were calculated. Equations 30 and 31 were then employed to calculate the ideal isobaric heat capacities over a temperature range of 30 to 300°F at 10°F intervals, thus producing 2016 data points. Data points for each gas mixture were plotted as shown in Figure 1. Each gas mixture showed a straight line relationship expressed by Equation 32:

$$C_{p_{H.C.}}^{\circ} = A(G) + B(G)T_r \quad (32)$$

$A(G)$ and $B(G)$ were found to be adequately represented by means of second degree polynomials in G as shown in Figure 2; that is,

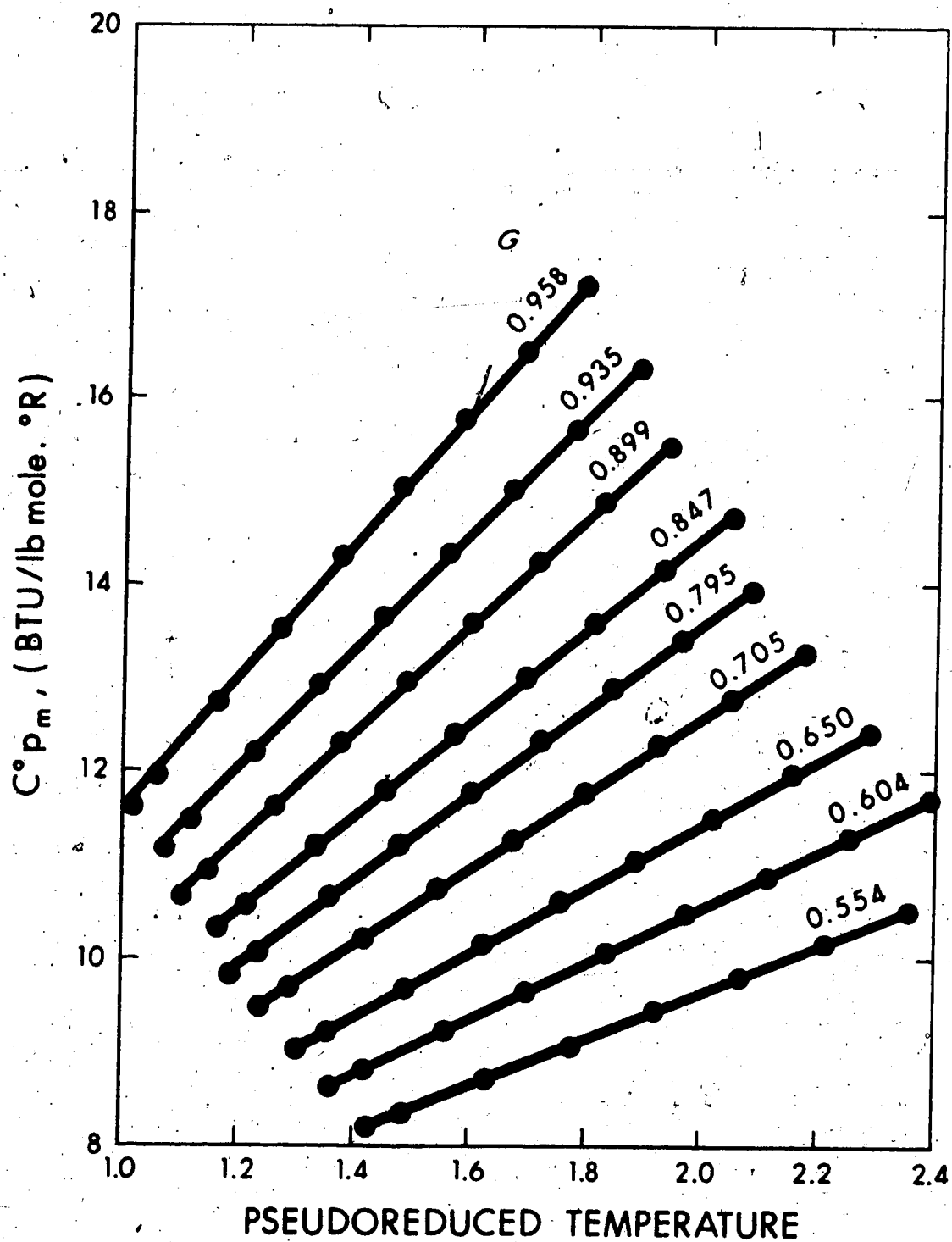


FIGURE 1. IDEAL ISOBARIC HEAT CAPACITY
OF NATURAL GASES

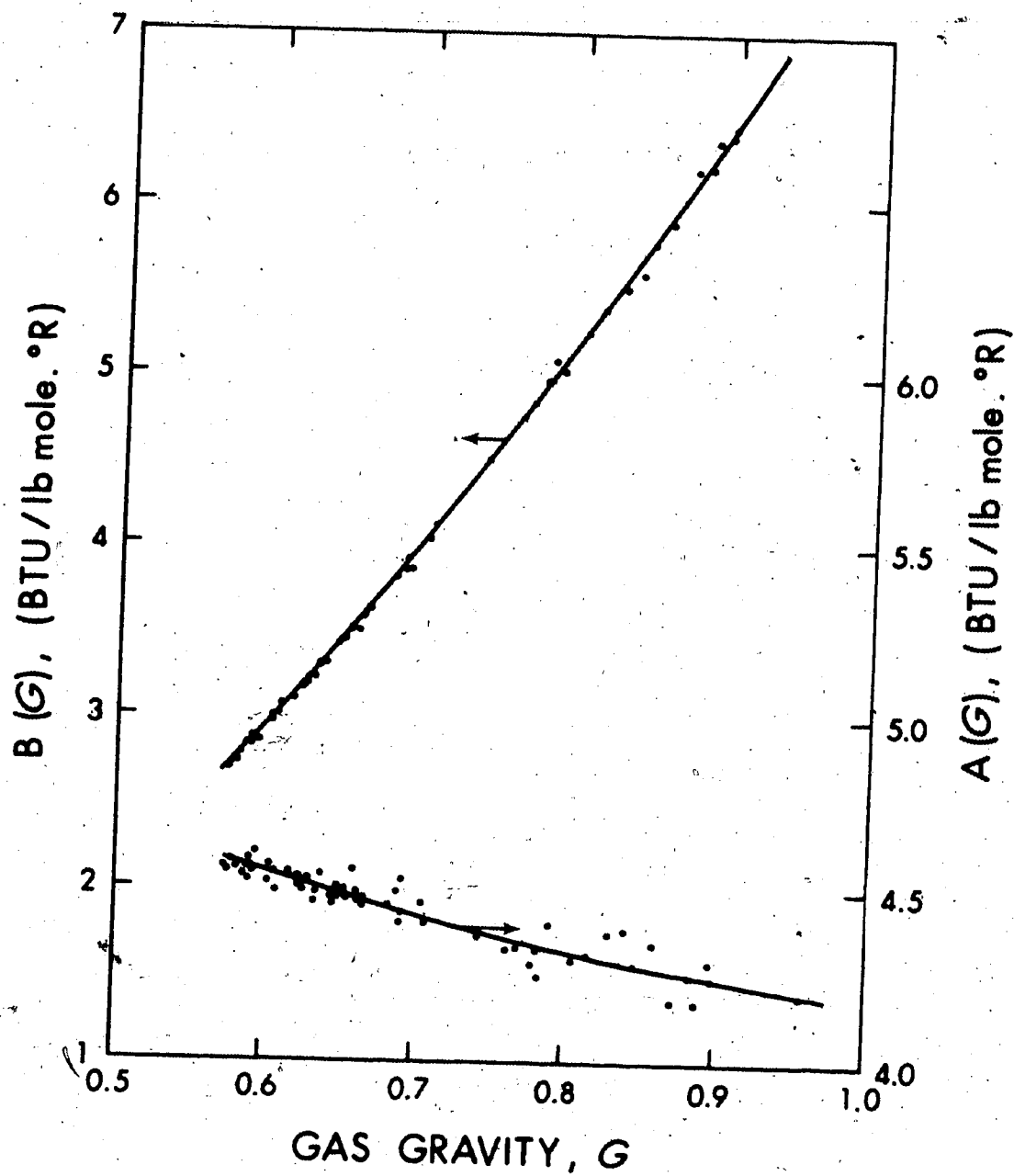


FIGURE 2. THE SLOPES AND INTERCEPTS IN EQUATION 32 AS FUNCTIONS OF GAS GRAVITY.

$$A(G) = B_1 + B_2 G + B_3 G^2 \quad (71)$$

and

$$B(G) = B_4 + B_5 G + B_6 G^2 \quad (72)$$

The coefficients in Equations 71 and 72 were obtained by minimizing the sum of squares of errors in $A(G)$ and $B(G)$ respectively. These coefficients (will be referred to as the set of coefficients No. 0) are:

$$B_1 = 5.596695 \quad B_2 = -2.233480 \quad B_3 = 0.807265$$

$$B_4 = -1.003900 \quad B_5 = 3.141600 \quad B_6 = 5.758700$$

Equations 32, 71 and 72 give an average absolute deviation of 0.079 per cent and a standard deviation of 0.00119 with respect to the 2016 data points used to determine the coefficients.

Studying the functional forms of Equations 34 through 36 individually showed that they can be expressed as:

$$F_{CO_2} = 1 + Y_{CO_2} (B_7 + B_8 T) \quad (73)$$

$$F_{H_2S} = 1 + Y_{H_2S} (B_9 + B_{10} T) \quad (74)$$

and

$$F_{N_2} = 1 + (B_{11} Y_{N_2} + B_{12} Y_{N_2}^2) + (B_{13} Y_{N_2} + B_{14} Y_{N_2}^2) T \quad (75)$$

where T = temperature, °R.

The coefficients B_7 through B_{14} were obtained by minimizing the sum of squares of errors in $C_{p,m}^\circ$ (as expressed by Equation 33) and were found to be;

$$\begin{aligned} B_7 &= 0.4258600 & B_8 &= 1.24323 \times 10^{-3} & B_9 &= -0.0405600 \\ B_{10} &= 1.00889 \times 10^{-3} & B_{11} &= 0.3623622 & B_{12} &= -0.4661581 \\ B_{13} &= 0.97570 \times 10^{-3} & B_{14} &= 2.708217 \times 10^{-3} \end{aligned}$$

Equations 33, 73, 74, and 75 give an average absolute deviation of 0.629 per cent and standard deviation of 0.00989 with respect to 2016 data points used to determine the coefficients B_7 through B_{14} .

Equations 32, 71, and 72 are applicable over a temperature range of 30-300°F to sweet natural gas which has no nitrogen. The same equations can be used over a broader temperature range; however, it is recommended that other sets of coefficients depending on temperature range of operation be used. The quality of fit of the new sets of coefficients is shown in Table 2.

2. Heat Capacity Departure:

The heat capacity departure model as expressed by Equation 37 was compared with the tabulated Weiss and Joffe generalized correlation (132). The model gives an average absolute deviation of 5.91 per cent when compared with 150 points.

TABLE 2

QUALITY OF FIT

Set of Coefficients No. 1		Set of Coefficients No. 2	
Temperature Range	80 ≤ T ≤ 1160°F	1160 ≤ T ≤ 2240°F	
No. of Test Points	2848 at 10°F intervals		
Average Deviation	-0.0103%	-0.00873%	
Average Absolute Deviation	0.7709%	0.8024%	
Standard Deviation	0.009586	0.009417	
Maximum Deviation	4.097%	2.430%	
No. of Points Having Deviations > 2%	229	160	
Set of Coefficients No. 1			
B ₁ = 4.72288190	B ₂ = -1.72132560	B ₃ = 3.21713790	
B ₄ = -0.91474246	B ₅ = 4.37040250	B ₆ = 3.04884500	
Set of Coefficients No. 2			
B ₁ = 6.31632700	B ₂ = -0.063316262	B ₃ = 11.54552600	
B ₄ = -0.97057393	B ₅ = 4.75787240	B ₆ = -0.80656836	

Tables 3 and 4 present comparisons of the various heat capacity departure correlations and the proposed generalized correlation as expressed by Equation 37 with the experimental data of Workman (133), Krase and Mackey (134). Table 5 presents comparisons of Equation 37 and the various $(C_p - C_p^\circ)$ generalized correlations with that of Weiss and Joffe.

Table 6 present comparisons of Equation 37 with the experimental data of Sage, et al. (135) for two natural gases. It also presents the effect of any errors in $(C_p - C_p^\circ)$ on the isobaric heat capacity (C_p) .

The Hydraulic Friction Factor

Table 7 presents comparisons between the values of friction factors as suggested by Moody (136) and Cullender and Smith (137). This comparison is carried out for the same 12 test-points used to test main models AA and BB.

Viscosity of Natural Gases

The Lee, et al. viscosity correlation (138) was tested with their 111 viscosity data points. This correlation gives an average absolute deviation of 3.71 per cent from experimental data and a standard deviation of 0.052527. The maximum deviation encountered is 16.38 per cent.

TABLE 3

COMPARISON OF EQUATION 37 AND VARIOUS ($C_p - C_p^\circ$) GENERALIZED CORRELATIONS
WITH THE EXPERIMENTAL DATA OF KRASE AND MACKEY FOR NITROGEN

T_r	P_r	$(C_p - C_p^\circ)_{\text{exp.}}$ Btu/lb mole. $^\circ R$	% Deviation from Experimental Values			
			Equation 37	Weiss and Joffe	Edmister	Hougen and Watson
2.37	5.97	1.47	14	21	-1	77
2.40	14.91	2.30	-6	9	*	*
2.56	14.91	2.00	-3	13	*	*
2.96	14.91	1.46	1	10	*	*
3.36	14.91	1.12	*	13	*	*

* Outside the range of correlation.

TABLE 4

COMPARISON OF EQUATION 37 AND VARIOUS $(C_p - C_p^\circ)$ GENERALIZED CORRELATIONS
WITH THE EXPERIMENTAL DATA OF WORKMAN FOR NITROGEN AND OXYGEN

P_r	$(C_p - C_p^{\circ})_{\text{exp.}}$ Btu/lb mole. °R	% Deviations from Experimental Values			
		Equation 37	Weiss and Joffe	Edmister	Hougen and Watson Lydersen, et al.
Oxygen ($T_r = 1.938$)					
0.119	0.1116	2	*	-1	59
0.597	0.3630	3	-1	-4	46
0.996	0.6140	-2	1	-2	45
1.392	0.8750	-1	-1	-4	39
1.792	0.1410	-1	-2	-4	38
2.390	0.1410	8	9	3	49
2.580	1.6800	-2	2	-6	34
Nitrogen ($T_r = 2.64$)					
0.295	0.0671	5	13	-15	49
0.895	0.2230	-3	2	-23	75
1.478	0.3820	-6	-1	-26	68
2.060	0.5400	-8	-4	-30	54
2.660	0.6860	-8	-2	-30	53
3.240	0.8260	-8	-2	-30	46
3.840	0.9620	-8	-1	-28	40
					*
					*
					-48
					-41
					-37
					-37
					-39

* Outside the range of correlation.

TABLE 5
COMPARISON OF EQUATION 37 AND OTHER ($C_p - C_p^\circ$) GENERALIZED
CORRELATIONS WITH THE WEISS AND JOFFE CORRELATION

Correlation	% Deviation from the Weiss and Joffe Values			
	$T_r = 1.0$	1.2	1.5	2.5
Hougen and Watson	59 to 62	26 to 55	33 to 64	43 to 63
Edmister	11 to 37	-1 to -11	-4 to -8	-22 to -23
Lydersen, et al.	-9 to 35	5 to 29	-61 to 20	-30 to -108
Sherwood	13 to 35	1 to 35	-25 to 5	8 to 66
Equation 37	*	-6 to 6	-17 to 0	-4 to -13

* Outside the range of correlation.

TABLE 5

COMPARISON OF $(C_p - C_p^0)$ ESTIMATED BY EQUATION 37 WITH THE SAGE, et al.
EXPERIMENTAL VALUES FOR TWO NATURAL GASES

Pressure (psia)	Temperature, °F											
	70	130	190	250	310	70	130	190	250	310	70	310
	$(C_p - C_p^0)_{exp.}$	$(C_p - C_p^0)_{exp.}$	$(C_p - C_p^0)_{exp.}$	$(C_p - C_p^0)_{exp.}$	$(C_p - C_p^0)_{exp.}$	$(C_p - C_p^0)_{calc.}^*$	$(C_p - C_p^0)_{calc.}^*$	$(C_p - C_p^0)_{calc.}^*$	$(C_p - C_p^0)_{calc.}^*$	$(C_p - C_p^0)_{calc.}^*$	$(C_p - C_p^0)_{calc.}^*$	$(C_p - C_p^0)_{calc.}^*$
	% Error in $(C_p - C_p^0)_{calc.}^*$	% Error in $(C_p - C_p^0)_{calc.}^*$	% Error in $(C_p - C_p^0)_{calc.}^*$	% Error in $(C_p - C_p^0)_{calc.}^*$	% Error in $(C_p - C_p^0)_{calc.}^*$	% Error in $(C_p - C_p^0)_{calc.}^*$	% Error in $(C_p - C_p^0)_{calc.}^*$	% Error in $(C_p - C_p^0)_{calc.}^*$	% Error in $(C_p - C_p^0)_{calc.}^*$	% Error in $(C_p - C_p^0)_{calc.}^*$	% Error in $(C_p - C_p^0)_{calc.}^*$	% Error in $(C_p - C_p^0)_{calc.}^*$
Natural Gas A												
200	0.3932	41.0	1.6	0.3324	15.0	0.5	0.2918	-4.8	-0.1	0.2594	-19.0	-0.4
400	0.8025	51.0	3.9	0.6769	19.8	1.2	0.5958	-2.9	-0.2	0.5391	-19.7	-0.9
500	0.0113	57.0	5.3	0.8572	21.8	1.7	0.7529	-2.1	-0.1	0.6809	-19.5	-1.1
Natural Gas B												
200	0.3847	29.1	1.1	0.3326	3.5	0.1	0.2805	-10.5	-0.3	0.2445	-22.0	-0.5
400	0.8055	33.4	2.5	0.6853	5.9	0.4	0.5731	-9.1	-0.5	0.4929	-20.5	-0.9
600	1.2854	35.0	4.2	1.0820	5.9	0.6	0.8816	-4.2	-0.6	0.7414	-18.9	-1.2

* $(C_p - C_p^0)$ calculated by using Equation 37, Btu/lb mole.°R.

TABLE 7

- COMPARISONS BETWEEN THE MOODY AND CULLENDER AND SMITH FRICTION FACTORS

Test- Point	Re _b (10 ⁶)	Re (10 ⁶)		Friction Factor			% Deviation of Cullender and Smith Values from Moody's	
		W	S	Moody ⁺		Cullender and Smith	W	S
				W	S			
A ₁	13.018	0.280	0.275	0.016719	0.016749	0.014319	-14.36	-14.51
A ₂	13.018	0.563	0.543	0.015693	0.015735	0.014319	-8.76	-9.00
A ₃	13.018	0.798	0.761	0.015340	0.015383	0.014319	-6.66	-6.92
A ₄	13.018	1.100	1.032	0.015091	0.015133	0.014319	-5.12	-5.38
A ₅	13.018	1.744	1.608	0.014825	0.014864	0.014319	-3.41	-3.67
A ₆	13.018	1.816	1.673	0.014807	0.014845	0.014319	-3.30	-3.54
B ₁	15.952	1.326	1.286	0.014479	0.014500	0.013681	-5.51	-5.65
B ₂	15.952	2.016	1.906	0.014243	0.014270	0.013681	-3.95	-4.13
B ₃	15.952	2.635	2.414	0.014131	0.014165	0.013681	-3.18	-3.42
B ₄	15.952	3.203	2.880	0.014065	0.014099	0.013681	-2.73	-2.97
C	15.952	4.076	3.515	0.013997	0.014037	0.013681	-2.26	-2.54
D	15.952	4.279	3.677	0.013985	0.014024	0.013681	-2.17	-2.45

where W indicates well head conditions

S indicates bottom-hole conditions

+ indicates absolute roughness of 0.0006 inch.

DATA AND RESULTS

Back Pressure Data

Back pressure data were used to test main models AA and BB. The test consists of 12 test-points taken from four back pressure tests for three natural gas wells. Table 8 presents a summary of these data. The pressure and temperature data are shown in Table 9. Appendix F is a tabulation of the source data and gas properties upon which this study was based.

Accuracy of Field Data

1. Flow Rate:

The flow rate was measured using either critical flow provers or orifice meters, and was reported at the standard conditions of 14.65 psia and 60°F. The accuracy of the flow measuring devices is not available.

2. Length of Tubing:

The reported values of length of tubing are those from Kelly Bushing (K.B.) to the mid-point of perforations (M.P.P.). The depth at which the pressure measuring device was located is not stated. The length of tubing of interest in this study is that between the well head and the bottom-hole pressure measuring device. This may be different from

TABLE 8
SUMMARY OF FIELD DATA

Length of Flow String, ft.	4753 - 11029
Flow Rate, MMSCF/Day	0.833 - 17.359
Gas/Condensate Ratio, MMSCF/bbl.	0.208 - ∞
Tubing-Head Temperature, °F	47 - 170
Bottom-Hole Temperature, °F	115 - 246
Tubing-Head Pressure, PSIA	985 - 3250
Bottom-Hole Pressure, PSIA	1104 - 4250
Water/Gas Ratio, bbl/MMSCF	0.0 - 5.2
y_{CO_2} , % mole	0.56 - 2.88
y_{H_2S} , % mole	0.0 - 18.89
y_{N_2} , % mole	1.05 - 4.56

TABLE 9

PRESSURE AND TEMPERATURE DATA

Test-Point	Q _o (MMSCF/Day)	Tubing-Head		Bottom-Hole		$\Delta P_{meas.}$ (PSIA)	W./G. Ratio (bbt/MMSCF)	G./Cond. Ratio (MMSCF/bbt)
		Temp. (°F)	Press. (PSIA)	Temp. (°F)	Press. (PSIA)			
A ₁	0.883	47	1660.0	115	1928.7	268.7	0.0	0.2388
A ₂	1.606	53	1566.0	115	1808.4	242.4	0.0	0.3071
A ₃	2.186	57	1473.0	115	1701.8	228.8	0.0	0.2081
A ₄	2.910	58	1360.0	115	1567.5	207.5	0.0	0.5569
A ₅	4.043	69	862.0	115	1048.0	186.0	0.0	0.3111
A ₆	4.291	68	895.0	115	1103.7	208.7	0.0	2.4655
B ₁	6.527	132	3249.4	246	4249.6	1000.2	1.5	∞
B ₂	9.598	140	3168.4	246	4205.8	1037.4	1.5	∞
B ₃	12.048	153	3078.4	246	4163.2	1084.8	1.5	∞
B ₄	14.277	158	2990.4	246	4123.2	1132.8	1.5	∞
C	15.606	155	2314.5	243	3295.5	981.0	5.2	∞
D	17.359	170	2599.5	243	3690.4	1090.9	1.3	∞

the values reported in Appendix F for length of tubing.

3. Pressure Data:

The tubing head pressure was measured with a dead weight gauge or a Barton gauge. The bottom-hole pressure was measured with an Amerada, tester type gauge, or the Sperry Sun gauge. The accuracy of these pressure devices is shown in Table 10 (139). The accuracy of the pressure data, is presented in Table 11.

TABLE 10
ACCURACY OF PRESSURE MEASURING DEVICES

Measuring Devices	Accuracy, % of Rated Pressure
Dead Weight Gauge	± 0.0
Barton Gauge	$\pm 0.5^*$ to $\pm 0.75^{**}$
Sperry Sun Gauge	± 0.1
Amerada or Tester Type Gauge	$\pm 0.2^*$ to $\pm 0.25^{**}$

* For rated pressure less than or equal to 2000 psi.

** For rated pressure greater than 2000 psi.

Results

Computer FORTRAN programs were written for main models AA and BB (method 2) and methods 1, 3, and 4.

Table 12 presents comparisons between the measured and

TABLE 11

ACCURACY OF PRESSURE DATA

Test-Point	Tubing-Head Press.		Bottom-Hole Press.	
	Rated Press. of Gauge (PSI)	Accuracy (PSI)	Rated Press. of Gauge (PSI)	Accuracy (PSI)
A ₁	2000	+10.0	2000	+4.0
A ₂	2000	+10.0	2000	+4.0
A ₃	2000	+10.0	2000	+4.0
A ₄	2000	+10.0	2000	+4.0
A ₅	2000	+10.0	2000	+4.0
A ₆	2000	+10.0	2000	+4.0
B ₁	4000	+30.0	5650	+11.3
B ₂	4000	+30.0	5650	+11.3
B ₃	4000	+30.0	5650	+11.3
B ₄	4000	+30.0	5650	+11.3
C	*	+0.0	5650 ⁺	+11.3
D	*	+0.0	5650 ⁺	+11.3

* Dead Weight Gauge was used.
+ Assumed.

calculated flowing bottom-hole pressures for all test-points. The test-points are presented in order of increasing flow rate. Table 13 presents the deviations of the calculated from the measured flowing bottom-hole pressures, while Table 14 exhibits the per cent deviations of the calculated from the measured pressure drops. Table 15 presents the various flowing bottom-hole pressures calculated by main models AA and BB.

TABLE 12

COMPARISONS BETWEEN THE MEASURED AND THE CALCULATED FLOWING BOTTOM-HOLE PRESSURES

Test-Point	B.H.P. calc. (PSIA)			
	B.H.P. meas. (PSIA)	Method 1	Method 2 (Main Model BB)*	Method 3
A ₁	1928.7	1888.4	1911.8	1916.0
A ₂	1808.4	1782.5	1799.5	1803.6
A ₃	1701.8	1682.5	1695.8	1699.6
A ₄	1567.5	1557.5	1567.7	1570.7
A ₅	1048.0	1029.6	1031.1	1031.9
A ₆	1103.7	1069.4	1071.1	1071.8
B ₁	4249.6	4138.4	4210.5	4235.4
B ₂	4205.8	4095.1	4152.2	4173.6
B ₃	4163.2	4046.2	4085.3	4101.8
B ₄	4123.2	4007.6	4037.5	4051.2
C	3295.5	3277.9	3293.0	3301.6
D	3690.4	3675.2	3686.1	3692.8
				3684.9

* G, T_c, p_c, p_c, Y_{N₂}, Y_{CO₂}, and Y_{H₂S} were supplied to the computer program as data.

TABLE 13

DEVIATIONS OF THE CALCULATED FLOWING BOTTOM-HOLE PRESSURES FROM THE MEASURED VALUES

Test- Point	Accuracy of B.H.P. calc. (PSI)	Accuracy of B.H.P. meas. (PSI)	Deviation of B.H.P. calc. from B.H.P. meas., (PSI)			
			Method 1	Method 2 (Main Model BB)	Method 3	Method 4
A1	+10.0	+ 4.0	- 40.3	-16.9	-12.7	-12.9
A2	+10.0	+ 4.0	- 25.9	- 8.9	- 4.8	- 5.3
A3	+10.0	+ 4.0	- 19.3	- 6.0	- 2.2	- 2.9
A4	+10.0	+ 4.0	- 10.0	+ 0.2	+ 3.2	+ 2.3
A5	+10.0	+ 4.0	- 18.4	-16.9	-16.1	-18.1
A6	+10.0	+ 4.0	- 34.3	-32.6	-31.9	-33.9
B1	+30.0	+11.3	-111.2	-39.1	-14.2	-16.5
B2	+30.0	+11.3	-110.7	-53.6	-32.2	-35.9
B3	+30.0	+11.3	-117.0	-77.9	-61.4	-66.3
B4	+30.0	+11.3	-115.6	-85.7	-72.0	-78.0
C	+ 0.0	+11.3	-17.6	- 2.5	+ 6.1	- 1.1
D	+ 0.0	+11.3	-15.2	- 4.3	+ 2.4	- 5.5

TABLE 14

DEVIATIONS OF THE CALCULATED PRESSURE DROPS FROM THE MEASURED VALUES

Test- point	Accuracy of $\Delta P_{meas.}$ in (% of $\Delta P_{meas.}$)	% Deviation of $\Delta P_{calc.}$ from $\Delta P_{meas.}$			
		Method 1	Method 2 (Main Model BB)*	Method 3	Method 4
A1	+4.9	-15.0	-6.3	-4.7	-4.8
A2	+6.1	-10.7	-3.7	-2.0	-2.2
A3	+6.7	-8.4	-2.6	-1.0	-1.3
A4	+7.5	-4.8	+0.1	+1.5	+1.0
A5	+6.7	-9.9	-9.1	-8.7	-9.7
A6	+7.5	-16.4	-15.6	-15.3	-16.2
B1	+4.4	-11.1	-3.9	-1.4	-1.7
B2	+4.3	-10.7	-5.2	-3.1	-3.5
B3	+4.1	-11.2	-7.4	-5.9	-6.3
B4	+3.9	-10.2	-7.6	-6.4	-6.9
C	+1.0	-1.8	-0.3	+0.6	-0.1
D	+0.9	-1.4	-0.4	+0.2	-0.5

* G , T_c , p_c , y_{N_2} , y_{CO_2} , and y_{H_2S} were supplied to the computer program as data.

TABLE 15
A COMPARISON OF CALCULATED BOTTOM-HOLE PRESSURES AND COMPUTATIONAL COSTS
FOR MAIN MODELS AA AND BB

Test-Point	Main Model BB			Main Model AA		
	a	b	c	c		
	B.H.P.cal. (PSIA)	Cost of Comput. (\$)	B.H.P.cal. (PSIA)	Cost of Comput. (\$)	B.H.P.cal. (PSIA)	Cost of Comput. (\$)
A ₁	1911.8	0.55	1911.8	0.55	1911.8	10.83
A ₂	1799.5	0.46	1799.5	0.46	1799.5	7.34
A ₃	1695.8	0.40	1695.8	0.40	1695.8	5.56
A ₄	1567.7	0.42	1567.7	0.42	1567.7	4.19
A ₅	1031.1	0.35	1031.1	0.35	1031.1	3.18
A ₆	1071.0	0.31	1071.1	0.31	1071.1	3.24
B ₁	4192.4	0.41	4210.5	0.41	*	*
B ₂	4135.9	0.43	4152.2	0.40	*	*
B ₃	4071.3	0.40	4085.3	0.36	*	*
B ₄	4025.2	0.35	4037.5	0.42	*	*
C	3289.1	0.42	3293.0	0.39	*	*
D	3678.3	0.32	3686.1	0.36	*	*

* Gas analysis is not available.

a G, Y_{CO₂}, Y_{N₂}, and Y_{H₂S} were supplied to the program as data.

b G, p, c, p_c, Y_{CO₂}, Y_{N₂}, and Y_{H₂S} were supplied to the program as data.

c The gas composition was supplied to the program as data.

DISCUSSION

Compressibility Factor of Natural Gases

The natural gas compressibility factor model as expressed by Equation 23, which is an equation of state, can predict Z with an average deviation of 1.013 per cent from experimental data as shown in Table 1. Such accuracy is considered acceptable for engineering applications.

Partial Derivatives of Compressibility Factor

The partial derivatives of the gas compressibility factor may be calculated by employing a numerical scheme or by using Equations 25 and 26. However, a numerical technique is cumbersome and costly. These equations are expected to work well since they were derived from an equation of state, Equation 23. Robinson (140) states that the B-W-R equation of state gives good predictions for the partial derivatives of compressibility factor.

Isobaric Heat Capacity of Natural Gases

1. Ideal Isobaric Heat Capacity:

The proposed model for ideal isobaric heat capacity of natural gases as expressed by Equations 32, 71, and 72 gives excellent results as was shown earlier in Preliminary Model Testing. However, it is limited in its application to

a specific temperature range and natural gases which consist of pure hydrocarbons. Furthermore, it requires that the hydrocarbon gas mixtures be rich in methane.

The ideal isobaric heat capacity of sweet natural gas may be expressed by the same equations over a wider range of temperature; however, new sets of coefficients, B_1 through B_6 , should be used. The set of coefficients No. 1 covers the range of temperature of $80 \leq T < 1160^\circ\text{F}$, while the set of coefficients No. 2 covers the range of temperature of $1160 \leq T \leq 2240^\circ\text{F}$. The temperature of 1160°F was chosen in a way that the standard deviation given by the model is almost the same when either set is used over its range as may be seen from Table 2.

It should be noted that the pseudoreduced temperature term appearing in Equation 32 is defined as $\frac{T}{p_c T_c}$. The pseudocritical temperature, $p_c T_c$, should be calculated using the molal average not Equations 66 and 69. However if the gas composition is not available Equations 66 and 69 may be used but higher errors should be expected.

Both sets of coefficients No. 1 and 2 were obtained by employing the Thinh, et al. equations (141) in obtaining the data points. These sets of coefficients could be used with Equations 32, 71, and 72 for natural gases operating at high temperatures.

Equations 33, 73, 74 and 75 were developed to take into account the effect of presence of N_2 , CO_2 , and H_2S . These correction equations give good results as was indicated

earlier in Preliminary Model Testing; however, their use is subject to the following restrictions.

- i. $32 \leq T \leq 300^{\circ}\text{F}$
- ii. $0.0 \leq Y_{\text{H}_2\text{S}} \leq 0.40$ mole fraction
- iii. $0.0 \leq Y_{\text{CO}_2} \leq 0.20$ mole fraction
- iv. $0.0 \leq Y_{\text{N}_2} \leq 0.25$ mole fraction
- v. They may only be applied when the set of coefficients No. 0 is used in Equations, 32, 71, and 72.

2. Heat Capacity Departure:

The proposed heat capacity departure model as expressed by Equation 37 was compared with the Weiss and Joffe (142) generalized correlation as shown in Table 5. When the various $(C_p - C_p^{\circ})$ generalized correlations (143,144) were compared with the Weiss and Joffe correlation, the proposed model gave the least deviation as shown in Table 5. This is probably due to the fact that both correlations use the B-W-R equation of state.

The Weiss and Joffe correlation was developed for pure hydrocarbons, then used to estimate the heat capacity departure of oxygen and nitrogen. It gave acceptable results and as a result was considered as a generalized correlation (145). In a similar manner the proposed heat capacity departure model was developed from an equation of state which describes the volumetric behavior of sweet natural gases. Equation 23, then was tested by comparing with the experimental data of Workman (146), Krase and Mackey (147) for oxygen and nitrogen. This model gives satisfactory agreement

with experimental data as shown in Tables 3 and 4. This check was performed in order to extend the applicability of the proposed model to non-hydrocarbon gases and thus to sour natural gases.

The various comparisons and tests performed on the proposed model suggest that it could be used as a generalized correlation for heat capacity departure. This is in spite of the fact that the coefficients of Equation 37 were originally developed using only sweet natural gas compressibility factor data.

It should be mentioned that although errors in $(C_p - C_p^\circ)$ may at times be large, the effect on overall accuracy of C_p is usually much smaller. Since the value of C_p is the quantity desired, it is the error in $\left(\frac{C_p - C_p^\circ}{C_p}\right)$ which is of importance not the error in $(C_p - C_p^\circ)$ (148) as shown in Table 6.

The proposed $(C_p - C_p^\circ)$ generalized correlation has an advantage over the others in that it is adaptable to computer usage.

The Hydraulic Friction Factor

The friction factor in the rough pipe region is represented by the Von Kármán formula. Colebrook (149) derived an approximation for the friction factor in this region when the absolute roughness of pipe is known. The friction factor equations suggested by Cullender and Smith (150) are but the Colebrook approximation when the absolute roughness is 0.0006 inch.

The assumption of fully developed turbulent flow of gas through tubing may be in error. As for example, for the 12 test-points used in testing main models AA and BB, flow was in the transition zone as the comparison between Reynolds number (R_e) and the boundary Reynolds number (R_{e_b}) in Table 7 may indicate. As a result the Cullender and Smith equations, which employ the assumption of fully turbulent flow, may give errors in friction factors estimated ranging from -2.17 to -14.51 per cent as shown in Table 7. This order of error may produce errors of -7.9 to -0.2 psi in the calculated flowing bottom-hole pressure as the comparison between methods 3 and 4 in Table 12 may indicate.

Overall Heat-Transfer Coefficient

In the development of the U_1 model it was assumed that the various solid layers were tightly fitted together with no intervening "air spaces". If the layers touch each other only at several points, the resistance to heat flow will increase remarkably (151). This suggests that the thermal resistance of the surroundings in gas well systems will have a comparatively large value due to the presence of the annular space.

In shallow gas wells where there exists one set of casing only, the composite cylindrical wall consists of tubing, annular space, casing, and cement. In deep wells the geometry of the composite wall is so complicated due to several factors such as: the presence of more than one set

of casing, eccentricity of both tubing and casing sets, lack of cement at some sections, inevitable reduction in well diameter as drilling advances, caving at some shale sections and enlargement at others, presence of fluid such as water in contact with either casing or cement at some unknown sections along the well course, corrosion, etc., Due to such complexity and lack of information, it is reasonable to assume that the system has an average thermal resistance, R^* , which will be determined by a trial-and-error procedure. Furthermore, this complexity and lack of information show it is not possible to employ model U_1 in gas well systems.

In model U_2 , the term $\frac{1}{h_0}$ is usually small, costly to compute, and of the order of 1 to 5 per cent of R^* . This suggests that model U_3 is preferable to U_2 since it requires considerably less computing time and does not affect the final results.

1. Stanton Number:

Comparison of the values of the Stanton number estimated using the various models indicates that there is little difference for low Prandtl Number fluids like air (152). However, the best expression that reproduces experimental data is that of Gowen and Smith (153). The reliability of this model in predicting the Stanton number lies in an implicit assumption that the friction factor and hence the equivalent sand relative roughness ($\frac{\epsilon_s}{d}$) completely characterizes the effect of roughness on temperature profile.

2. Prandtl Number of Natural Gases:

The Prandtl number, which is necessary to calculate the Stanton number, is defined as in Equation 50. For a gas mixture, it is recommended that the Prandtl number be calculated using Equation 50; however, the various terms appearing in the definition should be estimated separately

(154).

Thermal Conductivity of Gas Mixtures

Equation 56 applies to paraffins giving an average deviation of 2.6 per cent from the experimental data (155), while Equation 57 applies to non-hydrocarbons giving an average deviation of 2.2 per cent from the experimental data (156). The Sutherland constant estimated from boiling point temperature as expressed by Equation 54 may not be reliable; however, an error of 20 per cent in the Sutherland constant produces only an error of 1 per cent in the estimated gas mixture thermal conductivity at atmospheric pressure (157). Hence the simplification introduced by Equation 54 (due to necessity) is justified.

Equations 59, 60, and 61 for the estimation of the thermal conductivity departure, are usually applicable to pure non-polar gases. However, Reid and Sherwood (158) have recommended the use of these equations for gas mixtures. Their recommendation was based on an assumption that a gas mixture can be considered as a hypothetical pure gas with pseudocritical properties. This suggestion has been accepted

because few experimental data for thermal conductivity of high-pressure gas mixtures are available.

The gas mixture thermal conductivity model is expected to give errors of less than 5 per cent for simple non-polar gas mixtures for $T_r < 1.3$. The same model may be used for polar and non-polar polyatomic gas mixtures, but larger errors may occur (159).

Viscosity of Pure and Natural Gases

Equation 63 was originally developed for monatomic gases. However, it is found to be remarkably good for polyatomic gases as well (160).

Equation 64 is an approximation for the collision integral function, Ω_v , for non-polar gases. It deviates 1 to 2 per cent from the tabulated function (161). It may also be applied to polar gases; however, errors up to 2 per cent may occur (162). The collision integral function is required for the estimation of pure gas viscosities at atmospheric pressure as given by Equation 53.

The Lee, et al. correlation (163) for calculating the viscosity of sweet natural gases gives errors ranging from -11.05 to +16.38 per cent from their experimental data. The presence of N_2 , CO_2 , and H_2S tends to increase the viscosity of natural gases (164). As a result the Lee, et al. correlation may give errors larger or smaller than those mentioned when applied to sour natural gases.

Flowing Bottom-Hole Pressure

As may be seen from Tables 12 and 13, method 1 always underestimates the bottom-hole pressure. Methods 2 and 4 underestimate 11 of 12 test-points, while method 3 underestimates 9 of 12 test-points. It can also be seen that method 3 deviates the least from the measured bottom-hole pressure for 10 of 12 test-points.

Table 14 indicates that methods 2, 3, and 4 estimate pressure drops for test-points A_2 , A_3 , A_4 , B_1 , C, and D with accuracy comparable to those of the measured values. The various methods estimate flowing bottom-hole pressures in excellent agreement with the measured values for test-points A_4 , C, and D as shown in Table 14. This is probably due to using accurate surface pressure gauges (Dead Weight Gauges) in test-points C and D.

The analysis of Table 14 exhibits what may be a systematic error in the calculated bottom-hole pressure as the flow rate increases as is shown in Figure 3. This behavior may explain the excellent agreement of the various methods with test-point A_4 .

Figure 4 shows examples of estimated (methods 1 and 2) and assumed (methods 3 and 4) temperature profiles. It shows that the fluid temperature at every depth is decreasing in the following order: methods 1, 2, and either 3 or 4. It can also be seen that the fluid temperature is always greater than or equal to the earth temperature. This suggests

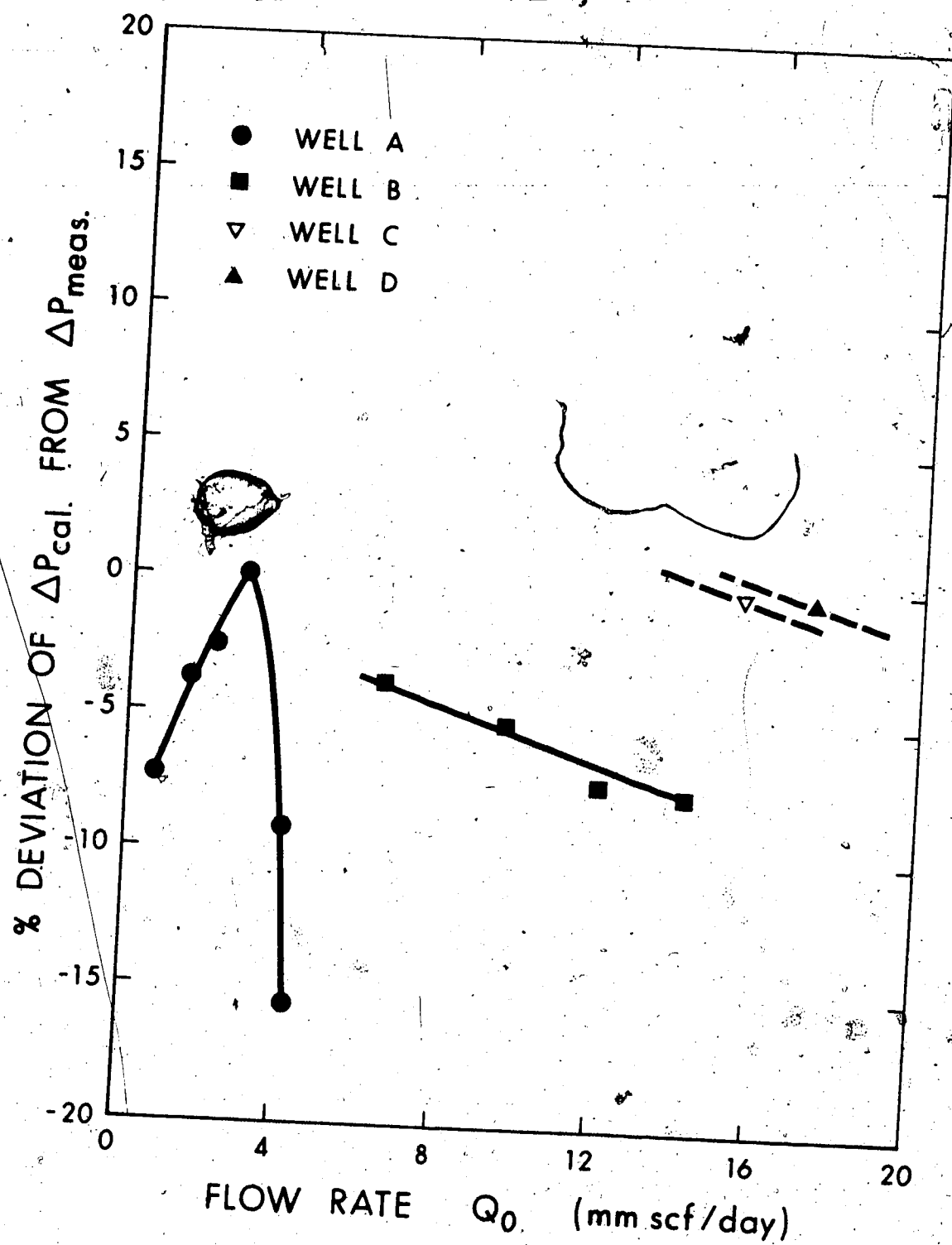


FIGURE-3. DEVIATION OF PRESSURE DROPS CALCULATED BY METHOD 2 FROM MEASURED VALUES FOR ALL TEST POINTS

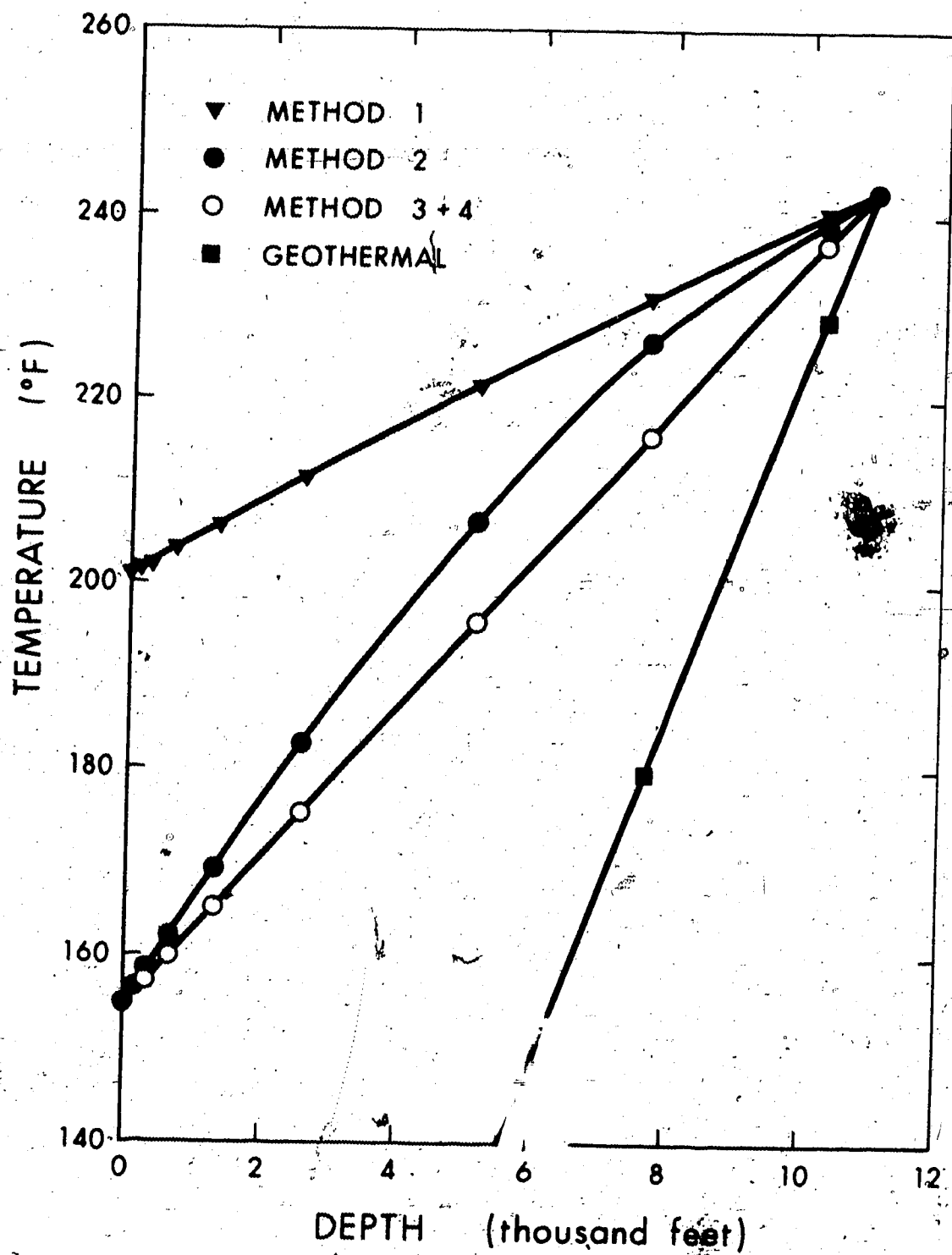


FIGURE 4. FLUID TEMPERATURE PROFILES CALCULATED BY METHODS 1, 2, 3, AND 4 FOR TEST-POINT C

that in these examples heat is being transferred from the flowing fluid to the surroundings. Furthermore, it suggests that methods 3 and 4 are equivalent to a hypothetical case where heat-transfer is such as to bring the fluid temperature profile to a straight line.

The analysis of Table 12 indicates that for any test-point, method 1 predicts the lowest flowing bottom-hole pressure followed by methods 2, 4, and finally method 3. This behavior is consistent in all test-points for methods 1, 2, and 3. However at three test-points, namely A_5 , A_6 , and D, methods 2 and 4 switch positions. This may be due to the use of the Cullender and Smith friction factor equations (which deviate from the Von Kármán rough pipe law by -0.04%, -0.04%, -0.41% for test-points A_5 , A_6 , and D respectively) in method 4.

Reasons for Deviations of Calculated from Measured Bottom-Hole Pressures

This discrepancy between measured flowing bottom-hole pressures and those calculated by method 2 may be attributed to a combination of the following factors:

1. The Existence of Two-Phase Flow:

The presence of condensate in the flowing gas has been dealt with according to Standing (165) when the gas/condensate ratio is greater than 40,000 SCF/bbl. He suggests

that the gravity of the flowing stream (G_f) be used instead of that of dry gas (G) according to:

$$G_f = \frac{G + 4.591 \frac{\rho_{\text{cond.}}}{R}}{1 + \frac{1123}{R}} \quad (71)$$

where R = Gas/Condens Ratio, SCF/bbl.

and $\rho_{\text{cond.}}$ = density of condensate, gm/cm³.

The pseudocritical pressure and temperature of the flowing stream may then be estimated using G_f .

This method of dealing with two-phase flow does not account for the effect of the liquid phase presence on the isobaric heat capacity, viscosity, thermal conductivity, and friction factor.

2. The Presence of Water Vapor in the Flowing Fluid:

The presence of water vapor in the flowing stream has not been accounted for. Although the water/gas ratio is small (1.5 bbl/MMSCF) except for test-point C (5.2 bbl/MMSCF), water is highly polar and as a result the behavior of the water-gas mixture system may deviate from the law of corresponding states. This behaviour may not be estimated accurately by the correlations employed and as a result errors in the estimated pressure may be introduced. However, it is surprising to find that test-point C is in very good agreement for the various methods.

3. Well Bore Deviation from Vertical:

Due to lack of information to the contrary all calcula-

that the gravity of the flowing stream (G_f) be used instead of that of dry gas (G) according to:

$$G_f = \frac{G + 4.591 \frac{\rho_{\text{cond.}}}{R}}{1 + \frac{1123}{R}} \quad (71)$$

where R = Gas/Condensate Ratio, SCF/bbl.

and $\rho_{\text{cond.}}$ = density of condensate, gm/cm³.

The pseudocritical pressure and temperature of the flowing stream may then be estimated using G_f .

This method of dealing with two-phase flow does not account for the effect of the liquid phase presence on the isobaric heat capacity, viscosity, thermal conductivity, and friction factor.

2. The Presence of Water Vapor in the Flowing Fluid:

The presence of water vapor in the flowing stream has not been accounted for. Although the water/gas ratio is small (1.5 bbl/MMSCF) except for test-point C (5.2 bbl/MMSCF), water is highly polar and as a result the behavior of the water-gas mixture system may deviate from the law of corresponding states. This behaviour may not be estimated accurately by the correlations employed and as a result errors in the estimated pressure may be introduced. However, it is surprising to find that test-point C is in very good agreement for the various methods.

3. Well Bore Deviation from Vertical:

Due to lack of information to the contrary all calcula-

tions of flowing bottom-hole pressure were made assuming that the tubing was situated in the vertical position. However, in practice well bore and therefore tubing segments may deviate from the vertical as is shown in Figure 5. Such deviations may result in a decrease in pressure drop, in which case the calculated flowing bottom-hole pressure may show positive deviations.

4. Unreliable Estimates of the Viscosity of Sour Gases:

The presence of appreciable amounts of CO_2 , N_2 , and H_2S tends to increase the viscosity of hydrocarbon mixtures (167). The Lee, et al. (168) viscosity correlation, therefore, may underestimate or overestimate the viscosity of sour natural gases depending on the sign and magnitude of errors produced by this correlation when applied to hydrocarbon mixtures. Consequently Reynolds number and hence the friction factor may be in error. Such errors may be significant especially at relatively low Reynolds number.

5. Uncertainty of the Length of Flow String:

The measured flowing bottom-hole pressures were reported at the mid-point-of-perforations (MPP). It is not known from the available data whether the measuring device was located at that point or whether in fact the tubing extended to the mid-point of perforations. If the tubing does not extend to the mid-point-of-perforations, method 2 tends to under-

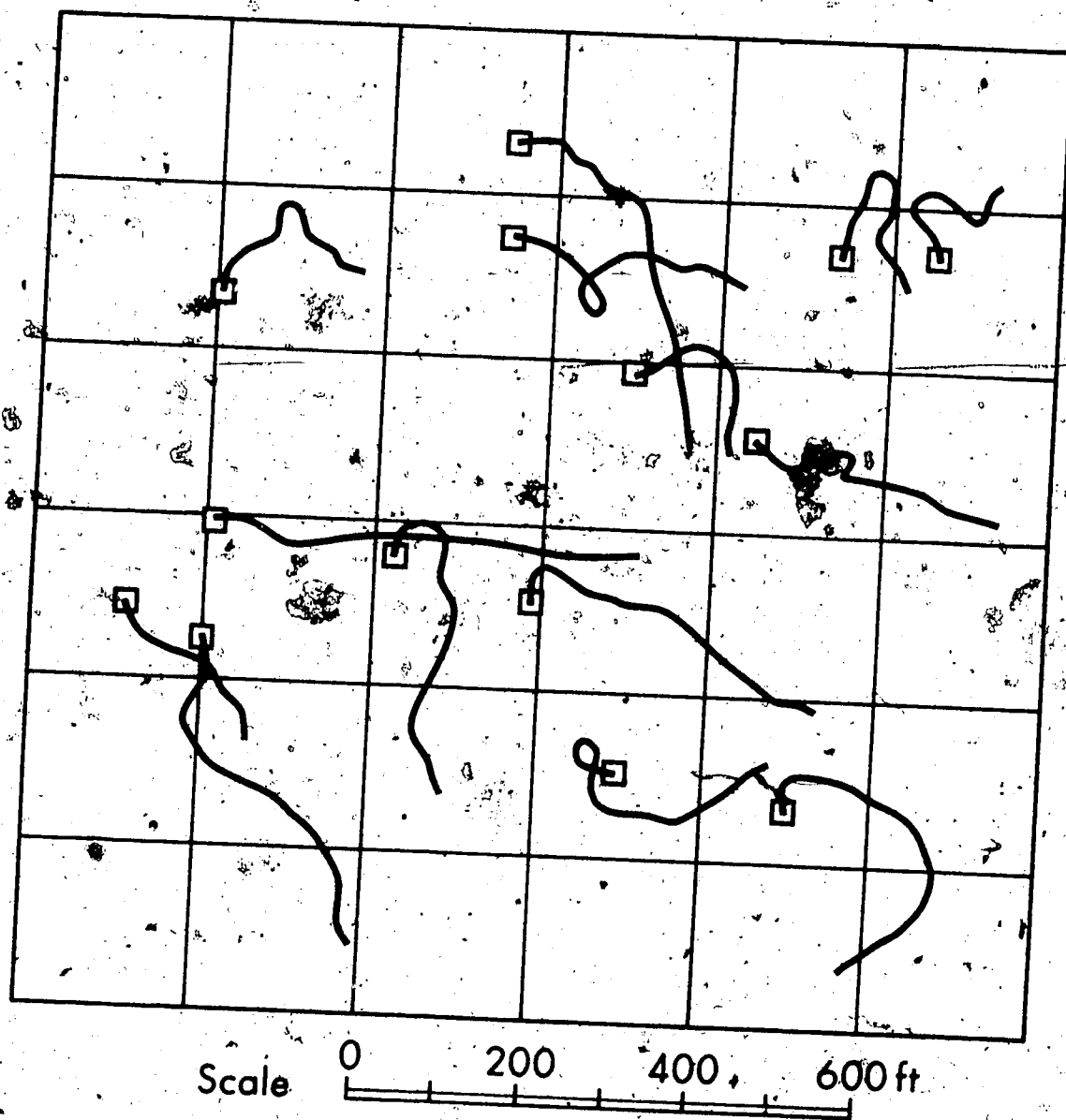


FIGURE 5. PLAN VIEW OF 14 VERTICAL WELLS
DRILLED TO 6000 FEET (166)

estimate the flowing bottom-hole pressure.

6. Variation of the Overall Heat-Transfer Coefficient:

Models U_2 and U_3 (due to necessity) were employed as a means for estimating the overall heat-transfer coefficient for complex geometries such as those encountered in gas wells. However, these approximations may not be reasonable from a physical standpoint. Consequently the heat-transfer between the flowing fluid and the surroundings may be in error which may affect the calculated flowing bottom-hole pressure.

Main models AA and BB employ models U_2 and U_3 respectively. It has been found that these main models yield calculated pressure and temperature profiles which are essentially the same; however, the cost of computation may be reduced by about a factor of 10 when using main model BB instead of AA as may be seen from Table 15. For example in six test-points the range of cost for the calculations when the computer programs were stored on disks, was:

Main model AA . . . \$3.18 - \$10.83

Main model BB . . . \$0.44 - \$ 0.75

7. Uncertainty of Absolute Roughness of Tubing:

This is probably one of the most important sources of error. The absolute roughness of tubing was taken as 0.0006 inch. This value is that shown by Cullender and Binckley (169) for new tubing. This value would in all likelihood be too low for old pipe as indicated by Colebrook (170). He reported

that an increase in roughness may cause 20 to 30 per cent reduction in the carrying capacity of pipes while the reduction in the area perpendicular to flow due to such increase in roughness may cause a reduction of only 2 to 3 per cent of the pipe carrying capacity. This implies that the same increase in roughness may cause an increase of 25 to 43 per cent of pressure loss due to friction if the pipe carrying capacity were kept at its initial value. Furthermore, Ippen (171) has reported a practical example where the absolute roughness was doubled within three years and the associated friction factor was increased by 20 per cent.

This suggests that flowing bottom-hole pressure calculated using the value of roughness of 0.0006 inch will yield values which are usually low. The 12 test-points presented here tend to support this.

CONCLUSIONS

As a result of this investigation, the following conclusions can be made:

1. The assumption of adiabatic flow in gas wells is unrealistic and yields the lower bound estimate for flowing bottom-hole pressure.
2. The apparent failure of main models AA and BB may be attributed to unreliable data and errors in different terms such as viscosity and absolute roughness, etc....
3. The complete solution of the Force-Momentum-Balance differential equation, as suggested in this study, yields the closest estimates to the measured flowing bottom-hole pressures.
4. The assumption of absolute roughness of 0.0006 inch for old and new tubing may be in error. In all likelihood this value is too small.

RECOMMENDATIONS

The following recommendations are made:

1. The Lee, et al. (172) viscosity correlation should be replaced by a more accurate one capable of representing the viscosity of sour as well as sweet natural gases. In this regard it is recommended that the Carr, et al. (173) Dean and Stiel (174) viscosity correlations be examined.
2. The parameters in the equations giving the pseudocriticals of natural gases as a function of gas gravity should be re-estimated by using data for real natural gases rather than fitting the Brown, et al. (175) correlation.
3. The absolute roughness of tubing used in gas wells should be re-estimated. This can be done by employing accurate surface and bottom-hole measurements for a large number of dry gas wells and applying main model A or B. The values of absolute roughness estimated may then be correlated to age of tubing, the degree of sourness of flowing gas, etc.

NOMENCLATURE

A, B	= functions of specified arguments or constants depending upon usage.
B^*	= a function of specified argument.
C_p	= isobaric heat capacity, Btu/lbm.°R.
C^*	= a function of specified argument.
C_p°	= molal ideal isobaric heat capacity, Btu/lb mole.°R.
$C_{p,m}^\circ$	= molal ideal isobaric heat capacity of gas mixtures, Btu/lb mole.°R.
$C_{p,m}^\circ \text{H.C.}$	= molal ideal isobaric heat capacity of pure hydrocarbon mixtures, Btu/lb mole.°R.
C_v°	= molal specific heat at infinite volume, Btu/lb mole.°R.
C_p	= molal isobaric heat capacity, Btu/lb mole.°R.
$C_{p,m}$	= molal isobaric heat capacity of gas mixtures, Btu/lb mole.°R.
C_v	= molal constant-volume specific heat, Btu/lb mole.°R.
d	= inner pipe diameter, inches.
D	= inner pipe diameter, feet.
D^*	= a function of specified argument.
D_o	= inner pipe diameter, feet.
exp	= exponential.
E^*	= a function of specified argument.
E_p	= estimated truncation error in calculated pressure, lb/in ² .
E_T	= estimated truncation error in calculated temperature, °F.

- f = Moody's friction factor, Dimensionless.
- f_1, f_2, f_{en} = functions
- F^*, F_{CO_2} = functions, specified arguments.
- F_{H_2S}, F_{N_2}
- g = local gravitational acceleration, 32.174 ft/sec².
- g_c = conversion factor, 32.174 lbm.ft/lb_f.sec².
- G = molar flow rate per unit area, lb mole/ft².sec.
- G = gas gravity, (air = 1).
- G_f = gravity of flowing stream, (air = 1).
- h = enthalpy, Btu/lbm.
- h_o = gas-film heat-transfer coefficient, Btu/ft².sec.°R.
- h_k = trial step size at kth interval, feet.
- J = conversion factor, 778.16 lb_f.ft/Btu.
- k = upper limit integer.
- k = thermal conductivity, cal/cm.sec.K.
- k^o = thermal conductivity at atmospheric pressure, cal/cm.sec.K.
- k_m^o = thermal conductivity of gas mixtures, at atmospheric pressure, cal/cm.sec.K.
- k_m = thermal conductivity of gas mixtures, cal/cm.sec.K.

$k_{P_1}, k_{P_2}, k_{P_3},$

$k_{P_4}, k_{P_5}, k_{T_1},$

$k_{T_2}, k_{T_3}, k_{T_4},$

= functions of specified arguments.

k_{T_5}

K

= a function of specified argument.

$\ln$

= natural logarithm.

$\log$

= logarithm, base 10.

L

= length variable, feet.

L^*

= length of tubing, feet.

M_w

= molecular weight or average molecular weight depending upon usage, lbm/lb mole.

n

= upper limit integer.

P

= pressure variable, lb_f/ft^2 absolute.

P_c

= critical pressure or pseudocritical pressure depending upon context, lb_f/in^2 absolute.

P_c

= pseudocritical pressure, lb_f/in^2 absolute.

$P_{cH.C.}$

= pseudocritical pressure of pure hydrocarbon mixtures, lb_f/in^2 absolute.

P_r

= reduced or pseudocritical pressure depending upon context, Dimensionless.

$\bar{P}_r$

= Prandtl Number; Dimensionless.

q

= heat flux, Btu/sec.

Q_o

= flow rate at standard conditions, MMSCF/Day.

R

= universal gas constant, $1545 lb_f \cdot ft/lb \text{ mole } ^\circ R$.

R	= Gas/Condensate Ratio, SCF/bbl.
R^*	= average thermal resistance of the surrounding as defined by Equation 13, $\text{ft}^2 \cdot \text{sec} \cdot ^\circ\text{F}/\text{Btu}$.
R^{**}	= average overall heat-transfer resistance, $\text{ft}^2 \cdot \text{sec} \cdot ^\circ\text{F}/\text{Btu}$.
Re	= Reynolds Number, Dimensionless.
Re_b	= Reynolds Number at the boundary between the transition and rough pipe regions, Dimensionless.
S	= Sutherland Constant, K.
$\sin \theta$	= $\sin \theta$, dimensionless.
T	= temperature variable, $^\circ\text{R}$.
T_b	= normal boiling point temperature, K.
T_c	= critical or pseudocritical temperature depending upon context, $^\circ\text{R}$.
p^T_c	= pseudocritical temperature, $^\circ\text{R}$.
$p^T_{c,H.C.}$	= pseudocritical temperature of pure hydrocarbon mixtures, $^\circ\text{R}$.
T_r	= reduced or pseudoreduced temperature, Dimensionless.
T_s	= temperature of surroundings, $^\circ\text{R}$.
$\frac{T_s}{T_c}$	= $\frac{T_s}{T_c}$, Dimensionless.
u	= velocity of flow, ft/sec .
$\langle u \rangle$	= average velocity of flow, ft/sec .
U	= overall heat-transfer coefficient, $\text{Btu}/\text{ft}^2 \cdot \text{sec} \cdot ^\circ\text{F}$.
V	= molal volume variable, $\text{ft}^3/\text{lb mole}$.
W	= shaft work done by gas, $\text{lb}_f \cdot \text{ft}$.

X, Y = functions of specified arguments.

Y = mole fraction, Dimensionless.

Y_{CO_2} = mole fraction of carbon dioxide.

Y_{H_2S} = mole fraction of hydrogen sulphide.

Y_{N_2} = mole fraction of nitrogen.

z = height above datum, feet.

Z = compressibility factor, Dimensionless.

Z_c = critical compressibility factor, Dimensionless.

Greek

α = inverse of the momentum correction factor.

γ = a parameter defined by Equation 58.

δ = a change in the value of a variable.

δ_1, δ_2 = tolerances used as conversion factors.

ΔP = pressure drop, lb_f/in^2 .

ϵ = absolute roughness, inches.

ϵ^* = $\sqrt{\frac{f}{8}} \cdot R_e \cdot \frac{\epsilon}{d}$ = Shear Reynolds number x relative roughness.

ϵ_0 = maximum energy of attraction for the Lennard-Jones potential, erg.

ϵ_3 = critical temperature correction as defined by Equation 24, °F.

ϵ_s = equivalent sand roughness, inches.

θ = dip angle of pipe, degrees.

μ = viscosity, c.p.

- μ^o = viscosity at one atmosphere, c.p.
- μ_m = viscosity of gas mixtures, c.p.
- ρ = density variable, lbm/ft³.
- ρ_c = critical density, lbm/ft³.
- ρ_{cond} = density of condensate, gm/cm³.
- ρ_r = reduced density, Dimensionless.
- σ = collision diameter for the Lennard-Jones Potential, Å.
- Σ = summation operator.
- χ = a function of specified argument as defined by Equation 49.
- Ω_v = collision integral.

REFERENCES

- (1) Ferguson, J.W., "Calculations of Back-Pressure Tests on Natural Gas Wells", O. & G. J., Vol. 37, No. 36, pp. 47-53, (Jan. 19, 1939).
- (2) Smith, R.V., "Determining Friction Factors for Measuring Productivity of Gas Wells", Trans. AIME, Vol. 189, pp. 73-82, (1950).
- (3) Poettmann, F.H., "The Calculation of Pressure Drop in the Flow of Natural Gas Through Pipe", Trans. AIME, pp. 317-326, (1951).
- (4) Sukkar, Y.K., and Cornell, D., "Direct Calculation of Bottom-Hole Pressure in Natural Gas Wells", Trans. AIME, Vol. 204, pp. 43-48, (1955).
- (5) Cullender, M.H., and Smith, R.V., "Practical Solution of Gas Flow Equations for Wells and Pipelines with Large Temperature Gradients", Trans. AIME, Vol. 207, pp. 281-287, (1956).
- (6) Young, K.L., "Effect of Assumptions Used to Calculate Bottom-Hole Pressures in Gas Wells", Trans. AIME, Vol. 240, pp. 547-550, (1967).
- (7) Dranchuk, P.M., Pet. E. 644 Lecture Notes, Department of Chemical and Petroleum Engineering, University of Alberta, Edmonton, Alberta, (1972).
- (8) Dranchuk, P.M., and McFarland, J.D., "The Effect of the Time Rate of Change of Momentum on Bottom-Hole Pressure in Flowing Gas Wells", J. of Can. Pet. Tech., Vol. 13, No. 2, pp. 34-38, (1974).
- (9) Dranchuk, P.M., and Quon, D., "A General Solution of the Equations Describing Steady State Turbulent Compressible Flow in Circular Conduits", J. of Can. Pet. Tech., Vol. 3, No. 2, pp. 60-65, (1964).
- (10) Ferguson, J.W., op. cit.
- (11) Smith, R.V., op. cit.
- (12) Poettmann, F.H., op. cit.

- (13) Sukkar, Y.K., and Cornell, D., op. cit.
- (14) Cullender, M.H., and Smith, R.V., op. cit.
- (15) Young, K.L., op. cit.
- (16) Dranchuk, P.M., op. cit.
- (17) Dranchuk, P.M., and McFarland, J.D., op. cit.
- (18) Aziz, K., "Ways to Calculate Gas Flow and Static Head",
Pet. Eng., Parts 1 through 10, (November, 1962,
through September, 1963).
- (19) Dranchuk, P.M., and Quon, D., op. cit.
- (20) Young, K.L., op. cit., p. 548.
- (21) Dranchuk, P.M., and McFarland, J.D., op. cit.
- (22) Cullender, M.H., and Smith, R.V., op. cit., pp. 285-286.
- (23) Standing, M.B., and Katz, D.L., "Density of Natural
Gases", Trans. AIME, Vol. 146, p. 144, (1942).
- (24) Elfrink, E.B., Sandberg, C.R., and Pollard, T.A., "A
New Compressibility Correlation for Natural
Gases and Its Application to Estimates of Gas-
in-Place", Trans. AIME, Vol. 186, pp. 219-223,
(1949).
- (25) Sarem, A.M., "Z-Factor Equation Developed for Use in
Digital Computers", O. & G. J., Vol. 59, No. 38,
p. 118, (September 18, 1961).
- (26) Dranchuk, P.M., and Quon, D., op. cit., p. 62.
- (27) Gray, E.H., and Sims, H.L., "New Factor Permits Z
Factor Determination in a Digital Computer",
O. & G. J., Vol. 57, No. 30, pp. 80-81, (July 20,
1959).
- (28) Benedict, M., Webb, G.B., and Rubin, L.C., "An Empiri-
cal Equation for Thermodynamic Properties of
Light Hydrocarbons and Their Mixtures", Chem.
Eng. Prog., Vol. 47, pp. 419-422, 449-454, and
609-620, (1951); J. Chem. Phys., Vol. 8, pp. 334-
345, (1940); Vol. 10, pp. 747-758, (1942).
- (29) Redlich, O., and Kwong, J.N.S., "On the Thermodynamics
of Solutions. V An Equation of State. Fuga-
cities of Gaseous Solutions", Chem. Rev., Vol.
44, pp. 233-244, (1949).

- (30) Carnahan, N.F., and Starling, K.E., "Equation of State for Nonattracting Rigid Spheres", J. Chem. Phys., Vol. 51, No. 2, p. 635 and 636, (1969).
- (31) Starling, K.E., "Thermo Data Refined for LPG", Hydrocarbon Processing, Parts 1 through 15, (March, 1971, through June, 1972).
- (32) Wichert, E., and Aziz, K., "Compressibility Factor of Sour Natural Gases", University of Calgary, Report No. CER/3, (July, 1970).
- (33) Wichert, E., and Aziz, K., "Compressibility Factor of Sour Natural Gases", J. Can. Chem. Eng., Vol. 49, pp. 267-273, (1971).
- (34) Redlich, O., and Kwong, J.N.S., op. cit.
- (35) Benedict, M., Webb, G.B., and Rubin, L.C., op. cit.
- (36) Barner, H.E., and Adler, S.B., "Low Temperature B-W-R Applications", Hydrocarbon Processing, Vol. 47, No. 10, pp. 150-156, (1968).
- (37) Barner, H.E., and Adler, S.B., "A Brief Review and an Extension of the B-W-R Equation of State", Paper Presented at Annual NGPA Meeting, Denver, Colorado, (March, 1970).
- (38) Bishnoi, P.R., and Robinson, D.B., "New Mixing Rules for the BWR Parameters to Predict Mixture Properties", J. Can. Chem. Eng., Vol. 50, pp. 101-107, (February, 1972).
- (39) Dranchuk, P.M., Purvis, R.A., and Robinson, D.B., "Computer Calculation of Natural Gas Compressibility Factors Using the Standing and Katz Correlation", Paper Presented at the 24th Annual Technical Meeting of the Petroleum Society of CIM, Edmonton, Alberta, (May 8-12, 1973).
- (40) Hall, K.R., and Yarborough, L., "A New Equation of State for Z-Factor Calculations", O. & G. J., Vol. 71, No. 25, pp. 82-92, (June 18, 1973).
- (41) Abou-Kassem, J.H., and Dranchuk, P.M., "Prediction of Compressibility Factors of Natural Gases by the Use of the Generalized Starling Equation of State", Unpublished Work, Department of Mineral Engineering, University of Alberta, Edmonton, Alberta, (1974).

- (42) Katz, D.L., Cornell, D., Kobayashi, R., Poettmann, F., Vary, J.A., Elenbaas, J.R., and Weinaug, C.F., "Handbook of Natural Gas Engineering", McGraw-Hill Book Co., Inc., N.Y., pp. 720-750, (1959).
- (43) Robinson, D.B., Macrygeorgos, C.A., and Govier, G.W., "The Volumetric Behavior of Natural Gases Containing Hydrogen Sulfide and Carbon Dioxide", Trans. AIME, Vol. 219, pp. 54-60, (1960).
- (44) Wichert, E., and Aziz, K., "Compressibility Factor of Sour Natural Gases", Report No. CEF/3, op. cit.
- (45) Wichert, E., and Aziz, K., "Compressibility Factor of Sour Natural Gases", op. cit.
- (46) Wichert, E., and Aziz, K., "Calculate Z's for Sour Gases", Hydrocarbon Processing, Vol. 51, pp. 119-122, (May, 1972).
- (47) Kobe, K.A., et al., "Thermochemistry for Petrochemical Industry", Pet. Refiner, (January, 1949, through November, 1954).
- (48) Tinh, T.P., Duran, J.L., Ramalho, R.S., and Kaliaguine, S., "Equations Improve C_p Predictions", Hydrocarbon Processing, Vol. 50, pp. 98-104, (January, 1971).
- (49) Kothari, M.S., and Doraiswamy, L.K., "Alternative Method for Vapor Heat Capacity", Hydrocarbon Processing and Petroleum Refiner, Vol. 43, pp. 133-137, (1964).
- (50) Reid, R.C., and Sherwood, T.K., "The Properties of Gases and Liquids", McGraw-Hill Book Co., New York, 2nd Edition, p. 324, (1966).
- (51) Dodge, B.F., "Physico-Chemical Factors in High Pressure Design", Ind. Eng. Chem., Vol. 24, No. 12, pp. 1353-1363, (1932).
- (52) Watson, K.M., and Smith, R.L., "Generalized High Pressure Properties of Gases", Nat. Pet. News, Vol. 28, No. 27, p. 29, 30, 34, 35, and 36, (July 1, 1936).
- (53) Lydersen, A.L., Greenkorn, R.A., and Hougen, O.A., "Generalized Thermodynamic Properties of Pure Fluids", Coll. Eng., University of Wisconsin, Eng. Expt. Station, Report No. 4, Madison, Wisconsin, (October, 1955).

- (54) Edmister, W.C., "Thermodynamic Properties of Hydrocarbons", Ind. Eng. Chem., Vol. 30, pp. 352-358, (March, 1948).
- (55) Weiss, A.H., and Joffe, J., "Generalized Correlation for Effect of Pressure on Isobaric Heat Capacity of Gases", Ind. Eng. Chem., Vol. 49, No. 1, pp. 120-124, (January, 1957).
- (56) Sherwood, T.K., "Velocity of Sound in Compressed Gases", J. Chem. and Eng. Data, Vol. 7, No. 1, pp. 47-50, (January, 1962).
- (57) Comings, E.W., and Nathan, M.F., "Thermal Conductivity of Gases at High Pressures", Ind. Eng. Chem., Vol. 39, No. 8, p. 969, (1947).
- (58) Reid, R.C., and Sherwood, T.K., op. cit., p. 270.
- (59) Sledjeski, E., "Isobaric Heat Capacity of Methane", Ind. Eng. Chem., Vol. 43, No. 12, pp. 2913-2915, (1951).
- (60) Seifarth, J.H., and Joffe, J., "Isobaric Heat Capacity of Propane", Ind. Eng. Chem., Vol. 44, No. 12, pp. 2894-2897, (1952).
- (61) Colebrook, C.F., "Turbulent Flow in Pipes, with Particular Reference to the Transition Region Between the Smooth and Rough Pipe Laws", Inst. of Civil Eng. J., Vol. 11, p. 133 and 137, (1939).
- (62) Moody, L.F., "Friction Factors for Pipe Flow", Trans. ASME, Vol. 66, p. 671, (1944).
- (63) Ibid., p. 673.
- (64) Dipprey, D.F., and Sabersky, R.H., "Heat and Momentum Transfer in Smooth and Rough Tubes at Various Prandtl Numbers", Int. J. Heat Mass Transfer, Vol. 6, p. 337, (1963).
- (65) Nikuradse, J., "Laws of Fluid Flow in Rough Pipes", Pet. Eng., Vol. 11, Part 3, p. 124, (June, 1940).
- (66) Colebrook, C.F., op. cit., pp. 134-135.
- (67) Ibid., p. 137.
- (68) Moody, L.F., op. cit., p. 672.
- (69) Cullender, M.H., and Smith, R.V., loc. cit.

- (70) Colebrook, C.F., "The Reduction of Carrying Capacity of Pipes with Age", Inst. of Civil Eng. J., Vol. 7, pp. 99-118, (1938).
- (71) Cullender, M.H., and Binckley, C.W., "Adaption of the Relative Roughness Correlation of the Coefficient of Friction to the Flow of Natural Gas in Gas Well Casings", Report Presented to Railroad Commission of Texas, Amarillo, (November 9, 1950).
- (72) Smith, R.V., Williams, R.H., and Dewees, E.J., "Measurement of Resistance to Flow of Fluids in Natural Gas Wells", Trans. AIME, Vol. 201, p. 285, (1954).
- (73) Smith, R.V., Miller, J.S., and Ferguson, J.W., "Flow of Natural Gas Through Experimental Pipe Lines and Transmission Lines", U.S. Bur. Mines Monograph 9, (1956).
- (74) Bird, R.B., Stewart, W.E., and Lightfoot, E.N., "Transport Phenomena", John Wiley and Sons, Inc., New York, p. 393, (1960).
- (75) Holman, J.P., "Heat Transfer", McGraw-Hill Book Co., New York, 3rd Edition, pp. 322-323, (1972).
- (76) Holst, P.H., M.Sc. Thesis, Department of Chemical and Petroleum Engineering, University of Alberta, Edmonton, Alberta, (1965).
- (77) Dranchuk, P.M., and Quon, D., loc. cit.
- (78) Colburn, A.P., Drew, T.B., and Worthington, H., "Heat Transfer in a 3-1 Hydrogen-Nitrogen Mixture at High Pressure", Ind. Eng. Chem., Vol. 39, No. 8, p. 961, (1947).
- (79) Martinelli, R.C., "Heat Transfer to Molten Metals", Trans. ASME, Vol. 69, p. 952, (1947).
- (80) Dipprey, D.F., and Sabersky, R.H., op. cit., p. 348.
- (81) Ibid., p. 351.
- (82) Gowen, R.A., and Smith, J.W., "Turbulent Heat Transfer from Smooth and Rough Surfaces", Int. J. Heat Mass Transfer, Vol. 11, p. 1665 and 1670, (1968).
- (83) Sheriff, N., and Gumley, P., "Heat Transfer and Friction Properties of Surfaces with Discrete Roughnesses", Int. J. Heat Mass Transfer, Vol. 9, p. 1313, (1966).

- (84) Reid, R.C., and Sherwood, T.K., op. cit., pp. 456-493.
- (85) Misic, D., and Thodos, G., "The Thermal Conductivity of Hydrocarbon Gases at Normal Pressures", AICHE J., Vol. 7, No. 2, pp. 264-266, (June, 1961).
- (86) Misic, D., and Thodos, G., "Atmospheric Thermal Conductivities for Gases of Simple Molecular Structure", J. Chem. and Eng. Data, Vol. 8, pp. 540-543, (1963).
- (87) Lindsay, A.L., and Bromley, L.A., "Thermal Conductivity of Gas Mixtures", Ind. Eng. Chem., Vol. 42, No. 8, pp. 1508-1511, (1950).
- (88) Stiel, L.I., and Thodos, G., "The Thermal Conductivity of Non-Polar Substances in the Dense Gaseous and Liquid Regions", AICHE J., Vol. 10, pp. 26-30, (1964).
- (89) Reid, R.C., and Sherwood, T.K., op. cit., pp. 396-428.
- (90) Ibid., p. 423.
- (91) Bicher, L.B., and Katz, D.L., "Viscosity of Natural Gases", Trans. AIME, Vol. 155, pp. 246-252, (1944).
- (92) Carr, N.L., Kobayashi, R., and Burrows, D.B., "Viscosity of Hydrocarbon Gases Under Pressure", Trans. AIME, Vol. 201, pp. 264-272, (1954).
- (93) Lee, A.L., "Viscosity of Light Hydrocarbons", API, New York, (1950).
- (94) Lee, A.L., Gonzalez, M.H., and Eakin, B.E., "The Viscosity of Natural Gases", Trans. AIME, Vol. 237, pp. 997-1000, (1966).
- (95) McCabe, W.L., and Smith, J.C., "Unit Operation of Chemical Engineering", McGraw-Hill Book Co., Inc., New York, 2nd Edition, p. 102, (1967).
- (96) Dranchuk, P.M., and Quon, D., loc. cit.
- (97) Cullender, M.H., and Smith, R.V., loc. cit.
- (98) Ibid.
- (99) Dranchuk, P.M., and McFarland, J.D., op. cit.

- (100) Purvis, R.A., Ph.D. Thesis, Department of Chemical and Petroleum Engineering, University of Alberta, Edmonton, Alberta, p. 102, (1972).
- (101) Dranchuk, P.M., Purvis, R.A., and Robinson, D.B., op. cit.
- (102) Wichert, E., and Aziz, K., "Calculate Z's for Sour Gases", op. cit., p. 121.
- (103) Moody, L.F., loc cit.
- (104) Colebrook, C.F., "Turbulent Flow in Pipes, with Particular Reference to the Transition Region Between the Smooth and Rough Pipe Laws", loc. cit.
- (105) Moody, L.F., op. cit., p. 684.
- (106) Cullender, M.H., and Binckley, C.W., op. cit.
- (107) Cullender, M.H., and Smith, R.V., loc. cit.
- (108) Bird, R.B., Stewart, W.E., and Lightfoot, E.N., loc. cit.
- (109) Dranchuk, P.M., and Quon, D., loc. cit.
- (110) Holman, J.P., op. cit., p. 150.
- (111) Gowen, R.A., and Smith, J.W., loc. cit.
- (112) Lindsay, A.L., and Bromley, L.A., op. cit., p. 1509.
- (113) Misic, D., and Thodos, G., "The Thermal Conductivity of Hydrocarbon Gases at Normal Pressures", op. cit., p. 265.
- (114) Misic, D., and Thodos, G., "Atmospheric Thermal Conductivities for Gases of Simple Molecular Structure", op. cit., p. 542.
- (115) Stiel, L.I., and Thodos, G., op. cit., p. 28.
- (116) Lee, A.L., Gonzalez, M.H., and Eakin, B.E., op. cit., p. 999.
- (117) Reid, R.C., and Sherwood, T.K., op. cit., p. 398.
- (118) Luft, N.W., and Kharabanda, O.P., "Simple Formula for Viscosity and Thermal Conductivity of Gases", Research Correspondence, Vol. 7, p. S36, (1954).

- (119) Reid, R.C., and Sherwood, T.K., op. cit., pp. 632-633.
- (120) Bird, R.B., Hirschfelder, J.O., and Curtiss, C.F., "Theoretical Calculation of the Equation of State and Transport Properties of Gases and Liquids", Trans. ASME, Vol. 76, p. 1015 and 1016, (1954).
- (121) Stiel, L.I., and Thodos, G., "Lennard-Jones Force Constants Predicted from Critical Properties", J. Chem. and Eng. Data, Vol. 7, No. 2, p. 235, (April, 1962).
- (122) Reid, R.C., and Sherwood, T.K., op. cit., p. 410.
- (123) Kay, W.B., "Density of Hydrocarbon Gases and Vapors", Ind. Eng. Chem., Vol. 28, No. 9, pp. 1014-1019, (September, 1936).
- (124) Brown, G.G., Katz, D.L., Oberfell, G.G., and Alden, R.C., "Natural Gasoline and the Volatile Hydrocarbons", Natural Gasoline Association of America, Tulsa, Oklahoma, p. 44, (1948).
- (125) Flores, J., Private Communication, Mobil Oil Company, Calgary, Alberta, (1974).
- (126) Carr, N.L., Kobayashi, R., and Burrows, D.B., op. cit., p. 270.
- (127) Hall, K.R., and Yarborough, L., op. cit., pp. 85-92.
- (128) Brown, G.G., Katz, D.L., Oberfell, G.G., and Alden, R.C., op. cit., pp. 24-39.
- (129) Wichert, E., and Aziz, K., "Compressibility Factor of Sour Natural Gases", Report No. CER/3, op. cit., pp. 129-152.
- (130) Ibid., p. 145 and 147.
- (131) Hall, K.R., and Yarborough, L., op. cit., p. 85 and 86.
- (132) Weiss, A.H., and Joffe, J., op. cit., p. 120.
- (133) Workman, E.J., "The Variation of the Specific Heats (C_p) of Oxygen, Nitrogen, and Hydrogen with Pressure", Phys. Rev., Vol. 37, pp. 1345-1355, (1931).

- (134) Krase, N.W., and Mackey, B.H., "The Specific Heats of Gases at High Pressures. I. Method and Apparatus at Room Temperature", J. Am. Chem. Soc., Vol. 52, pp. 108-115, (1930).
- (135) Sage, B.H., Botkin, D.F., and Lacey, W.N., "Joule-Thomson Coefficients for Two Natural Gases", Trans. AIME, Vol. 151, pp. 216-220, (1943).
- (136) Moody, L.F., op. cit., p. 676.
- (137) Cullender, M.H., and Smith, R.V., loc. cit.
- (138) Lee, A.L., Gonzalez, M.H., and Eakin, B.E., loc. cit.
- (139) Mumby, M.E., Private Communication, Gas Department, Alberta Energy Resources Conservation Board, Calgary, Alberta, (1974).
- (140) Robinson, D.B., Private Communication, Department of Chemical Engineering, University of Alberta, Edmonton, Alberta, (1974).
- (141) Thinh, T.P., Duran, J.L., Ramalho, R.S., and Kaliaguine, S., op. cit., pp. 99-103.
- (142) Weiss, A.H., and Joffe, J., loc. cit.
- (143) Ibid., p. 123.
- (144) Sherwood, T.K., op. cit., p. 48.
- (145) Weiss, A.H., and Joffe, J., op. cit., pp. 122-124.
- (146) Workman, E.J., op. cit.
- (147) Krase, N.W., and Mackey, B.H., op. cit.
- (148) Reid, R.C., and Sherwood, T.K., loc. cit.
- (149) Colebrook, C.F., "Turbulent Flow in Pipes, with Particular Reference to the Transition Region Between the Smooth and Rough Pipe Laws", op. cit., pp. 141-143.
- (150) Cullender, M.H., and Smith, R.V., loc. cit.
- (151) Bird, R.B., Stewart, W.E., and Lightfoot, E.N., op. cit., p. 286.
- (152) Gowen, R.A., and Smith, J.W., op. cit., p. 1671 and 1672.
- (153) Ibid.

- (154) Reid, R.C., and Sherwood, T.K., op. cit., p. 511.
- (155) Misic, D., and Thodos, G., "The Thermal Conductivity of Hydrocarbon Gases at Normal Pressures", op. cit., p. 266.
- (156) Misic, D., and Thodos, G., "Atmospheric Thermal Conductivities for Gases of Simple Molecular Structure", op. cit., p. 543.
- (157) Lindsay, A.L., and Bromley, L.A., op. cit., p. 1510.
- (158) Reid, R.C., and Sherwood, T.K., op. cit., p. 491.
- (159) Ibid., p. 492.
- (160) Ibid., p. 398.
- (161) Ibid., pp. 399-401.
- (162) Ibid., p.
- (163) Lee, A.L., Gonzalez, M.H., and Eakin, B.E., loc. cit.
- (164) Carr, N.L., Kobayashi, R., and Burrows, D.B., op. cit., p. 268 and 269.
- (165) Railroad Commission of Texas, "Back Pressure Test for Natural Gas Wells", State of Texas.
- (166) Gatlin, C., "Petroleum Engineering Drilling and Well Completion", Prentice-Hall, Inc., Englewood Cliffs, N.J., p. 143, (1965).
- (167) Carr, N.L., Kobayashi, R., and Burrows, D.B., loc. cit.
- (168) Lee, A.L., Gonzalez, M.H., and Eakin, B.E., loc. cit.
- (169) Cullender, M.H., and Binckley, C.W., op. cit.
- (170) Colebrook, C.F., "The Reduction of Carrying Capacity of Pipes with Age", op. cit., p. 105.
- (171) Ippen, A.T., Discussion of Moody's Paper, Trans. ASME, Vol. 66, p. 679, (1944).
- (172) Lee, A.L., Gonzalez, M.H., and Eakin, B.E., loc. cit.
- (173) Carr, N.L., Kobayashi, R., and Burrows, D.B., op. cit., pp. 267-268.

- (174) Dean, D.E., and Stiel, L.I., "The Viscosity of Non-Polar Gas Mixtures at Moderate and High Pressures", AICHE J., Vol. 11, No. 3, P. 528 and 529, (1965).
- (175) Brown, G.G., Katz, D.L., Oberfell, G.G., and Alden, R.C., loc. cit.
- (176) Sledjeski, E.W., op. cit., p. 2913.
- (177) Colebrook, C.F., loc. cit.
- (178) Niesse, D.H., "Technical Comment", Simulation, Vol. 14, pp. 93-94, (1970).
- (179) Chai, A.S., and Burgin, G.H., "Comment on the Runge-Kutta-Merson Algorithm", Simulation, Vol. 15, pp. 89-91, (1970).

APPENDIX A

DEVELOPMENT OF A MATHEMATICAL MODEL FOR STEADY STATE COMPRESSIBLE FLOW OF GAS IN PIPES

For the steady state flow of gas through a differential length of pipe of known dimensions and position in space, the Total Energy Balance Equation and the Force Momentum Balance Equation, based on the fluid element contained in a differential length (dL) may be written.

The Total Energy Balance Equation:

$$\delta q - \frac{dW}{J} = \frac{\pi D^2 G}{4 M_w} \left[dh + \frac{\langle u \rangle d\langle u \rangle}{g_c J} + \frac{g}{g_c} \frac{dz}{J} \right] \quad (A-1)$$

The Force Momentum Balance Equation:

$$\left(\frac{V}{M_w} \right) dP + \frac{\langle u \rangle d\langle u \rangle}{g_c} + \frac{g}{g_c} dz + \frac{f \langle u \rangle^2 dL}{2 g_c D} = 0 \quad (A-2)$$

where δq = the transferred heat from the surroundings to the fluid element, Btu.

W = shaft work, $lb_f \cdot ft$.

G = molal flow rate per unit area, $lb \text{ mole/sec. ft.}^2$

h = enthalpy, Btu/lbm.

$\langle u \rangle$ = mean velocity of fluid, ft/sec.

P = pressure, lb_f/ft^2 absolute.

V = molal volume, $\text{ft}^3/\text{lb mole}$.

f = Moody's hydraulic friction factor.

z = elevation above datum, ft.

D = I.D. of pipe, ft.

M_w = the molecular weight of gas, $\text{lbm}/\text{lb mole}$.

α = the inverse of the momentum correction factor

$$= \frac{\langle u \rangle^2}{u^2} \approx 1.0$$

J = conversion factor.

$$= 778.16 \text{ lb}_f \cdot \text{ft}/\text{Btu}$$

g = local gravitational acceleration.

$$= 32.174 \text{ ft}/\text{sec}^2 \text{ (assumed).}$$

g_c = conversion factor.

$$= 32.174 \frac{\text{lbm} \cdot \text{ft.}}{\text{lb}_f \cdot \text{sec}^2}$$

Equations (A-1) and (A-2) were rearranged into forms appropriate to numerical solution. This task was accomplished as follows:

By definition:

$$G = \frac{\langle u \rangle}{V}$$

Thus

$$\langle u \rangle = GV \quad (\text{A-3})$$

and

$$d\langle u \rangle = G dV \quad (\text{A-4})$$

The dip angle of the pipe is θ , where

$$\sin \theta = S = \frac{dz}{dL} \quad (A-5)$$

The equation of state is:

$$PV = ZRT \quad (A-6)$$

From which

$$dV = \frac{ZR}{P} dT + \frac{RT}{P} dZ - \frac{ZRT}{P^2} dP \quad (A-7)$$

and

$$\left(\frac{\partial V}{\partial T}\right)_P = \frac{ZR}{P} + \frac{RT}{P} \left(\frac{\partial Z}{\partial T}\right)_P \quad (A-8)$$

where $Z = \text{fcn}(T_r, P_r)$. Thus

$$dZ = \left(\frac{\partial Z}{\partial T_r}\right)_{P_r} dT_r + \left(\frac{\partial Z}{\partial P_r}\right)_{T_r} dP_r$$

or

$$dZ = \left(\frac{\partial Z}{\partial T}\right)_P dT + \left(\frac{\partial Z}{\partial P}\right)_T dP \quad (A-9)$$

The gas does no shaft work; that is,

$$dW = 0 \quad (A-10)$$

The transferred heat to the fluid element from the surroundings may be given by

$$\delta q = \pi D d L U (T_s - T) \quad (A-11)$$

where T_s = the temperature of the surroundings, °F.

T = the temperature of the flowing fluid element, °F.

U = the overall heat-transfer coefficient, Btu/ft². °F.

It is known that $h = f(T, P)$. Thus

$$dh = \left(\frac{\partial h}{\partial T}\right)_P \cdot dT + \left(\frac{\partial h}{\partial P}\right)_T \cdot dP$$

$$dh = c_p \cdot dT + \left(\frac{\partial h}{\partial P}\right)_T \cdot dP \quad (A-12)$$

But $\left(\frac{\partial h}{\partial P}\right)_T = \frac{V - T\left(\frac{\partial V}{\partial T}\right)_P}{J M_w}$ since h is in Btu/lbm units. (A-13)

Equations (A-6), (A-7), (A-12), and (A-13) give

$$dh = c_p \cdot dT - \frac{RT^2}{J P M_w} \left(\frac{\partial Z}{\partial T}\right)_P \cdot dP \quad (A-14)$$

where c_p is in Btu/lbm. °R units.

$$\frac{\langle u \rangle d \langle u \rangle}{\alpha g_c} = \frac{G^2 V dV}{\alpha g_c} \quad \text{by virtue of Equations (A-3) and (A-4).}$$

$$= \frac{G^2 Z^2 R^2 T^2}{\alpha g_c P^2} \left(\frac{dT}{T} + \frac{dZ}{Z} - \frac{dP}{P} \right) \quad \text{with the aid of Equation (A-7).} \quad (\text{A-15})$$

$$\frac{f \langle u \rangle^2 dL}{2g_c D} = \frac{f G^2 Z^2 R^2 T^2}{2g_c D P^2} \cdot dL \quad \text{with the aid of Equation (A-3)} \quad (\text{A-16})$$

$$\frac{V}{M_w} dP = \frac{ZRT}{PM_w} dP \quad \text{with the aid of Equation (A-6)} \quad (\text{A-17})$$

$$\frac{g}{g_c} dz = \frac{g}{g_c} s dL \quad \text{by virtue of Equation (A-5)} \quad (\text{A-18})$$

Combining Equations (A-1), (A-9), (A-10), (A-11), (A-14), (A-15), and (A-18) and rearranging yield

$$\frac{dP}{dL} \left[\frac{G^2 Z R^2 T^2}{\alpha g_c J P^2} \left(\frac{\partial Z}{\partial P} \right)_T - \frac{RT^2}{J P M_w} \left(\frac{\partial Z}{\partial T} \right)_P - \frac{G^2 Z^2 R^2 T^2}{\alpha g_c J P^3} \right] +$$

$$\frac{dT}{dL} \left[c_p + \frac{G^2 Z^2 R^2 T}{\alpha g_c J P^2} + \frac{G^2 Z^2 R^2 T^2}{\alpha g_c J P^2} \left(\frac{\partial Z}{\partial T} \right)_P \right] =$$

$$\frac{4U(T_s - T)/M_w - \frac{g}{g_c} DGS}{DGJ} \quad (\text{A-19})$$

Combining Equations (A-2), (A-9), (A-15), (A-16), (A-17) and (A-18) and rearranging yield

$$\frac{dP}{dL} \left[\frac{g_c P^2}{M_w} + \frac{G^2 RTP}{\alpha} \left(\frac{\partial Z}{\partial P} \right)_T - \frac{G^2 ZRT}{\alpha} \right] +$$

$$\frac{dT}{dL} \left[\frac{G^2 ZRP}{\alpha} + \frac{G^2 RTP}{\alpha} \left(\frac{\partial Z}{\partial T} \right)_P \right] = - \frac{g_c P^2}{ZT} - \frac{f G^2 Z RTP}{2D} \quad (A-20)$$

By employing the definitions of P_r and T_r and multiplying both sides of Equation (A-19) by -1 , Equations (A-19) and (A-20) may be written as:

$$\frac{dP_r}{dL} \left[- \frac{G^2 Z R^2 T^2}{\alpha J g_c P^2} \left(\frac{\partial Z}{\partial P_r} \right)_{T_r} + \frac{RTT_r}{JP_r M_w} \left(\frac{\partial Z}{\partial T_r} \right)_{P_r} + \frac{G^2 Z^2 R^2 T^2}{\alpha g_c JP^2 P_r} \right] +$$

$$\frac{dT_r}{dL} \left[- \frac{T_r}{T_r} c_p - \frac{G^2 Z^2 R^2 T^2}{\alpha g_c JP^2 T_r} - \frac{G^2 Z R^2 T^2}{\alpha J g_c P^2} \left(\frac{\partial Z}{\partial T_r} \right)_{P_r} \right] =$$

$$\frac{\frac{g}{g_c} DGS - 4UJ(T_{sr} - T_r)/M_w}{DGJ} \quad (A-21)$$

and

$$\frac{dP_r}{dL} \left[\frac{g_c P^3}{P_r M_w} - \frac{G^2 RTZP}{\alpha P_r} + \frac{G^2 RTP}{\alpha} \left(\frac{\partial Z}{\partial P_r} \right)_{T_r} \right] +$$

$$\frac{dT_r}{dL} \left[\frac{G^2 ZRPT}{\alpha T_r} + \frac{G^2 RTP}{\alpha} \left(\frac{\partial Z}{\partial T_r} \right)_{P_r} \right] =$$

$$- \frac{g_s P^3}{ZRT} - \frac{f G^2 Z RTP}{2D} \quad (A-22)$$

Equations (A-21) and (A-22) may be written as:

$$D^* \frac{dP_r}{dL} + E^* \frac{dT_r}{dL} = F^* \quad (A-23)$$

and $A^* \frac{dP_r}{dL} + B^* \frac{dT_r}{dL} = C^*$ respectively (A-24)

where

$$A^* = \left[\frac{g_c P^3}{P_r M_w} - \frac{G^2 R T Z P}{\alpha P_r} + \frac{G^2 R T P}{\alpha} \left(\frac{\partial Z}{\partial P_r} \right)_{T_r} \right]$$

$$= \text{fcn}(T_r, P_r, T_c, P_c, M_w, Q_0, D) \quad (A-25)$$

$$B^* = \left[\frac{G^2 Z R P T}{\alpha T_r} + \frac{G^2 R T P}{\alpha} \left(\frac{\partial Z}{\partial T_r} \right)_{P_r} \right]$$

$$= \text{fcn}(T_r, P_r, T_c, P_c, M_w, Q_0, D) \quad (A-26)$$

$$C^* = \left[-\frac{g_s P^3}{Z R T} - \frac{f G^2 Z R T P}{2 D} \right]$$

$$= \text{fcn}(T_r, P_r, T_c, P_c, M_w, Q_0, D, \left(\frac{E}{d} \right), R_e, S)$$

But $R_e = \frac{\rho \langle u \rangle D}{\mu} = \text{fcn}(T_r, P_r, Q_0, D)$

$$C^* = \text{fcn}(T_r, P_r, T_c, P_c, M_w, Q_o, D, (\frac{\epsilon}{d}), S) \quad (A-27)$$

$$D^* = \left[-\frac{G^2 Z R^2 T^2}{\alpha g_c J P^2} \left(\frac{\partial Z}{\partial P_r} \right)_{T_r} + \frac{R T T_r}{J P_r M_w} \left(\frac{\partial Z}{\partial T_r} \right)_{P_r} + \frac{G^2 Z^2 R^2 T^2}{\alpha g_c J P^2 P_r} \right]$$

$$= \text{fcn}(T_r, P_r, T_c, P_c, M_w, Q_o, D) \quad (A-28)$$

$$E^* = \left[-\frac{T_r}{T_r} c_p - \frac{G^2 Z^2 R^2 T^2}{\alpha g_c J P^2 T_r} - \frac{G^2 Z R^2 T^2}{\alpha g_c J P^2} \left(\frac{\partial Z}{\partial T_r} \right)_{P_r} \right]$$

$$= \text{fcn}(T_r, P_r, T_c, P_c, M_w, Q_o, D) \quad (A-29)$$

and

$$F^* = \frac{\frac{g}{g_c} D G S - 4 U J (T_{sr} - T_r) / M_w}{D G J}$$

$$= \text{fcn}(T_r, P_r, T_c, P_c, M_w, Q_o, S, U) \quad (A-30)$$

Solving Equations (A-23) and (A-24) simultaneously for $\frac{dP_r}{dL}$ and $\frac{dT_r}{dL}$ gives

$$\frac{dP_r}{dL} = \frac{C^* \cdot E^* - B^* \cdot F^*}{A^* \cdot E^* - B^* \cdot D^*} \quad (A-31)$$

and

$$\frac{dT_r}{dL} = \frac{A^* \cdot F^* - C^* \cdot D^*}{A^* \cdot E^* - B^* \cdot D^*} \quad (A-32)$$

In order to obtain a solution, $\frac{dP_r}{dL}$ and $\frac{dT_r}{dL}$ must be finite.

Equations (A-31) and (A-32) with the aid of Equations (A-25) through (A-30) may be rewritten as

$$\frac{dP_r}{dL} = \text{fon}(T_r, P_r, T_c, P_c, M_w, Q_0, D, \frac{\epsilon}{d}, U, S) \quad (A-33)$$

and

$$\frac{dT_r}{dL} = \text{fon}(T_r, P_r, T_c, P_c, M_w, Q_0, D, \frac{\epsilon}{d}, U, S) \quad (A-34)$$

For a specific flow case, that is, when the pipe dimensions and position in space (L^*, D, S), flow rate (Q_0) and gas properties (T_c, P_c, M_w), are given, Equations (A-33) and (A-34) reduce to

$$\frac{dP_r}{dL} = \text{fon}(T_r, P_r, (\frac{\epsilon}{d}), U) \quad (A-35)$$

and

$$\frac{dT_r}{dL} = \text{fon}(T_r, P_r, (\frac{\epsilon}{d}), U) \quad (A-36)$$

Equations (A-35) and (A-36) constitute the mathematical model for a steady-state compressible flow of gas in pipes.

APPENDIX B

DEVELOPMENT OF MATHEMATICAL EXPRESSIONS FOR THE PARTIAL DERIVATIVES OF Z WITH RESPECT TO T_r AND P_r

The reduced Benedict-Webb-Rubin equation of state has the form

$$Z = 1 + \left(A_1 + \frac{A_2}{T_r} + \frac{A_3}{T_r^3}\right)\rho_r + \left(A_4 + \frac{A_5}{T_r}\right)\rho_r^2 + A_5 A_6 \frac{\rho_r^5}{T_r} + A_7 \frac{\rho_r^2}{T_r^3} (1 + A_8 \rho_r^2) \exp(-A_8 \rho_r^2) \quad (B-1)$$

where

$$Z = \frac{PV}{RT} = \frac{PM_w}{\rho RT} \quad (B-2)$$

At the critical point

$$Z_c = \frac{P_c M_w}{\rho_c RT_c} \quad (B-3)$$

Dividing Equation (B-2) by Equation (B-3) and rearranging give

$$Z = Z_c \frac{P_r}{\rho_r T_r} \quad (B-4)$$

Combining Equations (B-1) and (B-4) produces another form of reduced Benedict-Webb-Rubin equation of state

$$P_r = \frac{1}{Z_c} \left[\rho_r T_r + (A_1 T_r + A_2 + \frac{A_3}{T_r^2}) \rho_r^2 + (A_4 T_r + A_5) \rho_r^3 + A_5 A_6 \rho_r^6 + A_7 \frac{\rho_r^3}{T_r^2} (1 + A_8 \rho_r^2) \exp(-A_8 \rho_r^2) \right] \quad (B-5)$$

Equation (B-4) was employed to derive $(\frac{\partial Z}{\partial T_r})_{P_r}$ and $(\frac{\partial Z}{\partial P_r})_{T_r}$ by partial differentiation as follows:

$$(\frac{\partial Z}{\partial T_r})_{P_r} = \frac{\partial}{\partial T_r} \left(\frac{Z_c P_r}{T_r \rho_r} \right)_{P_r}$$

$$= Z_c P_r \left[-\frac{1}{\rho_r T_r^2} - \frac{1}{T_r \rho_r^2} \left(\frac{\partial \rho_r}{\partial T_r} \right)_{P_r} \right]$$

$$= \frac{Z_c P_r}{\rho_r T_r} \left[-\frac{1}{T_r} - \frac{1}{\rho_r} \left(\frac{\partial \rho_r}{\partial T_r} \right)_{P_r} \right]$$

$$(\frac{\partial Z}{\partial T_r})_{P_r} = Z \left[-\frac{1}{T_r} - \frac{1}{\rho_r} \left(\frac{\partial \rho_r}{\partial T_r} \right)_{P_r} \right] \quad (B-6)$$

And

$$\begin{aligned}
 \left(\frac{\partial Z}{\partial P_r}\right)_{T_r} &= \frac{\partial}{\partial P_r} \left(\frac{Z_c P_r}{T_r \rho_r} \right)_{T_r} \\
 &= \frac{Z_c}{T_r} \left[\frac{1}{\rho_r} - \frac{P_r}{\rho_r^2} \left(\frac{\partial \rho_r}{\partial P_r}\right)_{T_r} \right] \\
 &= \frac{Z_c P_r}{\rho_r T_r} \left[\frac{1}{P_r} - \frac{1}{\rho_r} \left(\frac{\partial \rho_r}{\partial P_r}\right)_{T_r} \right] \\
 \left(\frac{\partial Z}{\partial P_r}\right)_{T_r} &= Z \left[\frac{1}{P_r} - \frac{1}{\rho_r} \left(\frac{\partial \rho_r}{\partial P_r}\right)_{T_r} \right] \quad (B-7)
 \end{aligned}$$

Equations (B-6) and (B-7) imply that the partial derivatives of Z with respect to T_r and P_r can be evaluated once the partial derivatives of ρ_r with respect to T_r and P_r are known.

$$\text{Now } \left(\frac{\partial \rho_r}{\partial P_r}\right)_{T_r} = \frac{1}{\left(\frac{\partial P_r}{\partial \rho_r}\right)_{T_r}} ; \text{ provided } \left(\frac{\partial P_r}{\partial \rho_r}\right)_{T_r} \neq 0 \quad (B-8)$$

$$\left(\frac{\partial \rho_r}{\partial T_r}\right)_{P_r} = - \left(\frac{\partial P_r}{\partial T_r}\right)_{\rho_r} / \left(\frac{\partial P_r}{\partial \rho_r}\right)_{T_r} \quad (B-9)$$

$\left(\frac{\partial P_r}{\partial T_r}\right)_{\rho_r}$ and $\left(\frac{\partial P_r}{\partial \rho_r}\right)_{T_r}$ were obtained by performing partial differentiation on Equation (B-5)

$$\left(\frac{\partial P_r}{\partial T_r}\right)_{\rho_r} = \frac{1}{Z_c} \left[\rho_r + \left(A_1 - \frac{2A_3}{T_r}\right) \rho_r^2 + A_4 \rho_r^3 - 2A_7 \frac{\rho_r^3}{T_r^3} (1 + A_8 \rho_r^2) \exp(-A_8 \rho_r^2) \right] \quad (B-10)$$

$$\left(\frac{\partial P_r}{\partial \rho_r}\right)_{T_r} = \frac{1}{Z_c} \left[T_r + 2\left(A_1 T_r + A_2 + \frac{A_3}{T_r}\right) \rho_r + 3(A_4 T_r + A_5) \rho_r^2 + 6A_5 A_6 \rho_r^5 + A_7 \frac{\rho_r^2}{T_r^2} \exp(-A_8 \rho_r^2) (3 + 3A_8 \rho_r^2 - 2A_8^2 \rho_r^4) \right] \quad (B-11)$$

Substituting Equation (B-11) into Equation (B-8) gives

$$\left(\frac{\partial \rho_r}{\partial P_r}\right)_{T_r} = \frac{Z_c}{\left[T_r + 2\left(A_1 T_r + A_2 + \frac{A_3}{T_r^2}\right)\rho_r + 3(A_4 T_r + A_5)\rho_r^2 + 6A_5 A_6 \rho_r^5 + A_7 \frac{\rho_r^2}{T_r^2} \exp(-A_8 \rho_r^2)(3 + 3A_8 \rho_r^2 - 2A_8^2 \rho_r^4) \right]} \quad (B-12)$$

Substituting Equations (B-10) and (B-11) into Equation (B-9) gives

$$\left(\frac{\partial \rho_r}{\partial T_r}\right)_{P_r} = - \frac{\left[\rho_r + \left(A_1 - \frac{2A_3}{T_r^3}\right)\rho_r^2 + A_4 \rho_r^3 - 2A_7 \frac{\rho_r^3}{T_r^3} (1 + A_8 \rho_r^2) \exp(-A_8 \rho_r^2) \right]}{\left[T_r + 2\left(A_1 T_r + A_2 + \frac{A_3}{T_r^2}\right)\rho_r + 3(A_4 T_r + A_5)\rho_r^2 + 6A_5 A_6 \rho_r^5 + A_7 \frac{\rho_r^2}{T_r^2} \exp(-A_8 \rho_r^2)(3 + 3A_8 \rho_r^2 - 2A_8^2 \rho_r^4) \right]} \quad (B-13)$$

Equations (B-6) and (B-7) with the aid of Equations (B-12) and (B-13) made it possible to estimate $(\frac{\partial Z}{\partial T})_{P_r}$ and $(\frac{\partial Z}{\partial P_r})_{T_r}$ at T_r and P_r analytically.

APPENDIX C

DEVELOPMENT OF A MATHEMATICAL EXPRESSION FOR HEAT CAPACITY DEPARTURE

The reduced form of the Benedict-Webb-Rubin equation of state is

$$P_r = \frac{1}{Z_c} \left[\rho_r T_r + (A_1 T_r + A_2 + \frac{A_3}{T_r^2}) \rho_r^2 + (A_4 T_r + A_5) \rho_r^3 + A_5 A_6 \rho_r^6 + A_7 \frac{\rho_r^3}{T_r^2} (1 + A_8 \rho_r^2) \exp(-A_8 \rho_r^2) \right] \quad (C-1)$$

The heat capacity departure from the ideal state is defined as $(C_p - C_p^\circ)$.

For a gas in the ideal state

$$C_p^\circ = C_v^\circ + R \quad (C-2)$$

Thus

$$(C_p - C_p^\circ) = C_p - (C_v^\circ + R) \quad (C-3)$$

By adding and subtracting C_v from the left hand side of Equation (C-3) and rearranging, one may get

$$(C_p - C_p^\circ) = (C_p - C_v) + (C_v - C_v^\circ) - R \quad (C-4)$$

where C_p = isobaric heat capacity at pressure P and temperature T , Btu/lb mole. $^\circ$ R.

C_p° = isobaric heat capacity at zero pressure and temperature T , Btu/lb mole. $^\circ$ R.

C_v = constant-volume heat capacity at pressure P and temperature T , Btu/lb mole. $^\circ$ R.

C_v° = constant-volume specific heat capacity at zero pressure or at infinite volume and temperature T , Btu/lb mole. $^\circ$ R.

R = universal gas constant, 1.98654 Btu/lb mole. $^\circ$ R.

Expressions for $(C_p - C_v)$ and $(C_v - C_v^\circ)$ may be obtained from Equation (C-1) by applying the following thermodynamic relationships (176):

$$\left(\frac{\partial C_v}{\partial V}\right)_T = T \left(\frac{\partial^2 P}{\partial T^2}\right)_V \quad (C-5)$$

$$(C_p - C_v) = - T \left(\frac{\partial P}{\partial T}\right)_V^2 / \left(\frac{\partial P}{\partial V}\right)_T \quad (C-6)$$

Equations (C-5) and (C-6) may be rewritten as:

$$\left(\frac{\partial C_v}{\partial \rho_r}\right)_{T_r} = - R Z_c \frac{T_r}{\rho_r^2} \left(\frac{\partial^2 P_r}{\partial T_r^2}\right)_{\rho_r} \quad (C-7)$$

$$(C_p - C_v) = R Z_c \frac{T_r}{\rho_r^2} \left(\frac{\partial P_r}{\partial T_r}\right)_{\rho_r}^2 / \left(\frac{\partial P_r}{\partial \rho_r}\right)_{T_r} \quad (C-8)$$

By performing partial differentiation on Equation (C-1), the following expressions were obtained:

$$\left(\frac{\partial P_r}{\partial T_r}\right) \rho_r = \frac{1}{Z_c} \left[\rho_r + \left(A_1 - \frac{2A_3}{T_r}\right) \rho_r^2 + A_4 \rho_r^3 - 2A_7 \frac{\rho_r^3}{T_r^3} (1 + A_8 \rho_r^2) \exp(-A_8 \rho_r^2) \right] \quad (C-9)$$

$$\left(\frac{\partial^2 P_r}{\partial T_r^2}\right) \rho_r = \frac{1}{Z_c} \left[6A_3 \frac{\rho_r^2}{T_r^4} + 6A_7 \frac{\rho_r^3}{T_r^4} (1 + A_8 \rho_r^2) \exp(-A_8 \rho_r^2) \right] \quad (C-10)$$

$$\left(\frac{\partial P_r}{\partial \rho_r}\right) T_r = \frac{1}{Z_c} \left[T_r + 2\left(A_1 T_r + A_2 + \frac{A_3}{T_r^2}\right) \rho_r + 3\left(A_4 T_r + A_5\right) \rho_r^2 + 6A_5 A_6 \rho_r^5 + A_7 \frac{\rho_r^2}{T_r^2} \exp(-A_8 \rho_r^2) (3 + 3A_8 \rho_r^2 - 2A_8^2 \rho_r^4) \right] \quad (C-11)$$

Substituting Equations (C-9), and (C-11) into Equation (C-7) gives

$$(c_p - c_v) = \frac{RT_r \left[1 + \left(A_1 - \frac{2A_3}{T_r} \right) \rho_r + A_4 \rho_r^2 - 2A_7 \frac{\rho_r^2}{T_r} (1 + A_8 \rho_r^2) \exp(-A_8 \rho_r^2) \right]^2}{\left[T_r + 2 \left(A_1 T_r + A_2 + \frac{A_3}{T_r} \right) \rho_r + 3(A_4 T_r + A_5) \rho_r^2 + 6A_5 A_6 \rho_r^5 + A_7 \frac{\rho_r^2}{T_r} \exp(-A_8 \rho_r^2) (3 + 3A_8 \rho_r^2 - 2A_8^2 \rho_r^4) \right]} \quad (C-12)$$

Substituting Equation (C-10) into Equation (C-7) gives

$$\left(\frac{\partial c_v}{\partial \rho_r} \right)_{T_r} = -R \left[\frac{6A_7}{T_r} + 6A_7 \frac{\rho_r}{T_r} (1 + A_8 \rho_r^2) \exp(-A_8 \rho_r^2) \right] \quad (C-13)$$

Boundary Conditions:

$$\text{At } \rho_r = 0, \quad c_v = c_v^0$$

$$\rho_r = \rho_r, \quad c_v = c_v$$

By separating variables in Equation (C-13), integrating, and applying the above Boundary Conditions, one may obtain:

$$\int_{C_v^0}^{C_v} (\partial C_v)_{T_r} = -R \left[\int_0^{\rho_r} \left(\frac{6A_7}{T_r^3} \right)_{T_r} d\rho_r + \int_0^{\rho_r} \left(6A_7 \frac{\rho_r}{T_r^3} \exp(-A_8 \rho_r^2) \right)_{T_r} d\rho_r \right. \\ \left. + \int_0^{\rho_r} \left(6A_7 A_8 \frac{\rho_r^3}{T_r^3} \exp(-A_8 \rho_r^2) \right)_{T_r} d\rho_r \right]$$

$$(C_v - C_v^0) = -R[X + Y + Q] \quad (C-14)$$

where

$$X = 6A_3 \frac{\rho_r}{T_r^3} \quad (C-15)$$

$$Y = -\frac{3A_7 \exp(-A_8 \rho_r^2)}{A_8 T_r^3} + \frac{3A_7}{A_8} \frac{1}{T_r^3} \quad (C-16)$$

$$Q = -3A_7 \frac{\rho_r^2}{T_r^3} \exp(-A_8 \rho_r^2) - \frac{3A_7}{A_8} \frac{\exp(-A_8 \rho_r^2)}{T_r^3} + \frac{3A_7}{T_r^3} \quad (C-17)$$

or

$$(C_v - C_v^0) = R \left[-6A_3 \frac{\rho_r}{T_r^3} - \frac{6A_7}{A_8 T_r^3} + \left(\frac{6A_7}{A_8 T_r^3} + 3A_7 \frac{\rho_r^2}{T_r^3} \right) \exp(-A_8 \rho_r^2) \right] \quad (C-18)$$

Substituting Equations (C-12) and (C-18) into Equation

(C-4) results in the final form of heat capacity departure equation, which may be written as:

$$(C_p - C_p^o) = R + \left[\begin{aligned} & -1 - 6A_3 \frac{\rho_r}{T_r^3} - \frac{6A_7}{A_8 T_r^3} \\ & + \left(\frac{6A_7}{A_8 T_r^3} + 3A_7 \frac{\rho_r^2}{T_r^3} \right) \exp(-A_8 \rho_r^2) \\ & T_r \left[1 + \left(A_1 - \frac{2A_3}{T_r^3} \right) \rho_r + A_4 \rho_r^2 \right. \\ & \quad \left. - 2A_7 \frac{\rho_r^2}{T_r^3} (1 + A_8 \rho_r^2) \exp(-A_8 \rho_r^2) \right]^2 \\ & T_r + 2(A_1 T_r + A_2 + \frac{A_3}{T_r^2}) \rho_r \\ & + 3(A_4 T_r + A_5) \rho_r^2 + 6A_5 A_6 \rho_r^5 \\ & + A_7 \frac{\rho_r^2}{T_r^2} \exp(-A_8 \rho_r^2) (3 + 3A_8 \rho_r^2 - 2A_8^2 \rho_r^4) \end{aligned} \right]$$

(C-19)

APPENDIX D

THE HYDRAULIC FRICTION FACTOR IN BOTH THE TRANSITION AND THE ROUGH PIPE REGIONS

Transition Region

This region is characterized by $4000 \leq R_e \leq \frac{3200}{(\epsilon/d)}$.
The hydraulic friction factor is given by the Colebrook transition law (177).

$$f = \left[-\frac{1}{2 \log\left(\frac{\epsilon/d}{3.7} + \frac{2.51}{R_e \sqrt{f}}\right)} \right]^2 \quad (D-1)$$

where f = Moody's friction factor.

ϵ = absolute roughness, inch.

d = I.D. of tubing, inch.

R_e = Reynolds number.

By taking the square root of both sides of Equation (D-1) and by letting $x = \sqrt{f}$, one may obtain

$$x = -\frac{1}{2 \log\left(\frac{\epsilon/d}{3.7} + \frac{2.51}{R_e x}\right)} \quad (D-2)$$

In order to solve for x in Equation (D-2), the Newton-

Raphson iterative method was used. The iterative equation has the following form:

$$x_{i+1} = x_i - \frac{g(x_i)}{g'(x_i)} ; \text{ providing } g'(x_i) \neq 0 \quad (D-3)$$

where $g(x_i)$ is obtained from Equation (D-2) as:

$$g(x_i) = 1 + 2x_i \log\left(\frac{(\epsilon/d)}{3.7} + \frac{2.51}{Re^{x_i}}\right) \quad (D-4)$$

Thus,

$$g'(x_i) = 2 \log\left(\frac{(\epsilon/d)}{3.7} + \frac{2.51}{Re^{x_i}}\right) - \frac{5.02 \log e}{Re^{x_i} \left(\frac{(\epsilon/d)}{3.7} + \frac{2.51}{Re^{x_i}}\right)} \quad (D-5)$$

Combining Equations (D-3), (D-4), and (D-5) provides the final form of the iterative equation for x ,

$$x_{i+1} = x_i - \frac{\left[1 + 2x_i \log\left(\frac{(\epsilon/d)}{3.7} + \frac{2.51}{Re^{x_i}}\right) \right]}{\left[2 \log\left(\frac{(\epsilon/d)}{3.7} + \frac{2.51}{Re^{x_i}}\right) - \frac{5.02 \log e}{Re^{x_i} \left(\frac{(\epsilon/d)}{3.7} + \frac{2.51}{Re^{x_i}}\right)} \right]} \quad (D-6)$$

Once Equation (D-6) is satisfied after n iterations

$$f = x_n^2$$

The Fully Rough Pipe Region

Flow of fluid falls in this region when $R_e > \frac{3200}{(\epsilon/d)}$.
The Van-Kármán equation was employed to estimate the hydraulic friction factor in this region.

$$f = \left[\frac{1}{2 \log\left(\frac{\epsilon/d}{3.7}\right)} \right]^2 \quad (D-7)$$

APPENDIX E

AN ALGORITHM FOR SOLVING THE SYSTEM OF DIFFERENTIAL EQUATIONS IN MAIN MODELS AA AND BB

The system of differential equations that describes the flow problem is as follows

$$\frac{dP_r}{dL} = f_1(T_r, P_r, R^*, \epsilon/d) \quad (E-1)$$

$$\frac{dT_r}{dL} = f_2(T_r, P_r, R^*, \epsilon/d) \quad (E-2)$$

$$\frac{dR^*}{dL} = 0 \quad (E-3)$$

The Boundary Conditions are:

$$\text{At } L = 0 : T = T_{(L=0)} = T_1$$

$$L = L^* : T = T_{(L=L^*)} = T_2 \text{ and } P = P_{(L=L^*)} = P_2$$

The modified Runge-Kutta-Merson algorithm as suggested by Niesse (178), Chai and Burgin (179) was utilized to solve this system of differential equations as follows:

(1) Let $(\epsilon/d) = 0.0006/d$ (for new tubing)

$$R^* = 7000 ; 1000 < R^* \leq 7000$$

(2) For a trial step size (h_k), compute k_{p_i} and k_{T_i} , $i=1,5$:

$$k_{p_1} = \frac{h_k}{3} f_1(T_{r_n}, P_{r_n}, R^*, \epsilon/d) ;$$

$$k_{T_1} = \frac{h_k}{3} f_2(T_{r_n}, P_{r_n}, R^*, \epsilon/d)$$

$$k_{p_2} = \frac{h_k}{3} f_1(T_{r_n} + k_{T_1}, P_{r_n} + k_{p_1}, R^*, \epsilon/d) ;$$

$$k_{T_2} = \frac{h_k}{3} f_2(T_{r_n} + k_{T_1}, P_{r_n} + k_{p_1}, R^*, \epsilon/d)$$

$$k_{p_3} = \frac{h_k}{3} f_1(T_{r_n} + \frac{k_{T_1}}{2} + \frac{k_{T_2}}{2}, P_{r_n} + \frac{k_{p_1}}{2} + \frac{k_{p_2}}{2}, R^*, \epsilon/d)$$

$$k_{T_3} = \frac{h_k}{3} f_2(T_{r_n} + \frac{k_{T_1}}{2} + \frac{k_{T_2}}{2}, P_{r_n} + \frac{k_{p_1}}{2} + \frac{k_{p_2}}{2}, R^*, \epsilon/d)$$

$$k_{p_4} = \frac{h_k}{3} f_1(T_{r_n} + \frac{3}{8} k_{T_1} + \frac{9}{8} k_{T_3}, P_{r_n} + \frac{3}{8} k_{p_1} + \frac{9}{8} k_{p_3}, R^*, \epsilon/d)$$

$$k_{T_4} = \frac{h_k}{3} f_2(T_{r_n} + \frac{3}{8} k_{T_1} + \frac{9}{8} k_{T_3}, P_{r_n} + \frac{3}{8} k_{p_1} + \frac{9}{8} k_{p_3}, R^*, \epsilon/d)$$

$$k_{p5} = \frac{h_k}{3} f_1(T_{r_n} + \frac{3}{2} k_{T1} - \frac{9}{2} k_{T3} + 6k_{T4}, P_{r_n} + \frac{3}{2} k_{p1} - \frac{9}{2} k_{p3} + 6k_{p4}, R^*, \epsilon/d)$$

$$k_{T5} = \frac{h_k}{3} f_2(T_{r_n} + \frac{3}{2} k_{T1} - \frac{9}{2} k_{T3} + 6k_{T4}, P_{r_n} + \frac{3}{2} k_{p1} - \frac{9}{2} k_{p3} + 6k_{p4}, R^*, \epsilon/d)$$

and calculate

$$P_{r_{n+1}} = P_{r_n} + \frac{3}{2} k_{p1} - \frac{9}{2} k_{p3} + 6k_{p4}$$

$$T_{r_{n+1}} = T_{r_n} + \frac{3}{2} k_{T1} - \frac{9}{2} k_{T3} + 6k_{T4}$$

(3) Compute the estimated truncation error

$$E_p = P_c \left[k_{p1} - \frac{9}{2} k_{p3} + 4k_{p4} - \frac{k_{p5}}{2} \right]$$

$$E_T = T_c \left[k_{T1} - \frac{9}{2} k_{T3} + 4k_{T4} - \frac{k_{T5}}{2} \right]$$

(4) Let δ_1 = tolerance in predicted pressure P .

δ_2 = tolerance in predicted temperature T .

If either $E_p > \delta_1$ or $E_T > \delta_2$, halve the trial step size

and start over from step (2) until $E_p \leq \delta_1$ and $E_T \leq \delta_2$.

- (5) Compute the corrected value of $P_{r_{n+1}}$ and $T_{r_{n+1}}$:

$$P_{r_{n+1}}(\text{corrected}) = P_{r_{n+1}} - \frac{E_p}{P_c}$$

$$T_{r_{n+1}}(\text{corrected}) = T_{r_{n+1}} - \frac{E_T}{T_c}$$

This correction raises the accuracy to the fourth order.

- (6) If $E_p < \frac{\delta_1}{16}$ and $E_T < \frac{\delta_2}{16}$, double the trial step size for the next integration step. If not, leave the trial step size as is for the next integration interval, h_{k+1} .

- (7) If $\left[L^* - \sum_{i=1}^k h_i \right] = 0$, go to step (8).

If $\left[L^* - \sum_{i=1}^k h_i \right] \geq h_{k+1}$, go to step (2).

If $\left[L^* - \sum_{i=1}^k h_i \right] < h_{k+1}$, let $h_{k+1} = \left[L^* - \sum_{i=1}^k h_i \right]$

and go to step (2).

- (8) At this stage, $T_{(L=0)}$ and $P_{(L=0)}$ are estimated.

If $T_{(L=0)} = T_1$, then $P_1 = P_{(L=0)}$. STOP.

If $T_{(L=0)} \neq T_1$, modify the value of R^* .

(9) Let $k = 1$ and start over from step (2).

For main model BB, R^{**} replaces R^* .

APPENDIX F

SOURCE DATA

WELL A

FIELD: Dunvegan

WELL NAME: Anderson et al.
Dunvegan 6-29

POOL: Debolt

LOCATION: 6-29-80-3 W6M

PERFORATIONS: 4748'-4785' K.B. DISTANCE TO M.P.P.: 4753' K.B.

TUBING SIZE: 2.441" I.D.

RESERVOIR TEMPERATURE: 115°F

GAS PROPERTIES:

$G = 0.6402$

$p_c = 669.9 \text{ PSIA}$

$T_g = 367.2 \text{ }^\circ\text{R}$

MOL.% $\text{CO}_2 = 0.56$

MOL.% $\text{N}_2 = 1.05$

MOL.% $\text{H}_2\text{S} = 0.00$

MOL.% $\text{C}_1 = 90.33$

MOL.% $\text{C}_2 = 4.04$

MOL.% $\text{C}_3 = 1.95$

MOL.% $i\text{-C}_4 = 0.31$

MOL.% $n\text{-C}_4 = 0.70$

MOL.% $i\text{-C}_5 = 0.20$

MOL.% $n\text{-C}_5 = 0.25$

MOL.% $\text{C}_6^{\text{S}} = 0.22$

MOL.% $\text{C}_7^{\text{S}} = 0.25$

MOL.% $\text{C}_8^+ = 0.07$

MOL.% $\text{H}_e = 0.07$

Density of Condensate = 0.777 gm/c.c.

WELL OPERATOR: Anderson Exploration Limited

TYPE OF TEST: Modified Isochronal

DATE OF TEST: April 5, 1971

GAS PRODUCED THROUGH: Tubing

WELL A (CONTINUED)

PRODUCTION AND PRESSURE DATA

Total Flow Rate (MMSCF/Day)	Tubing Head		Press. at M.P.P. (PSIA)	Condensate Prod. Rate (bbl/Day)	Water/Gas Ratio (bbl/MMSCF)
	Press. (PSIA)	Temp. (°F)			
0.833	1660.0	47	1928.7	3.48	0.0
1.606	1566.0	53	1808.4	5.22	0.0
2.186	1473.0	57	1701.8	10.44	0.0
2.910	1360.0	58	1567.5	5.22	0.0
4.291	895.0	68	1103.7	1.74	0.0
4.043	862.0	69	1048.0	12.97	0.0

WELL B

FIELD: BigstoneWELL NAME: Pan Am. HB,G-2POOL: D-3ALOCATION: 2-25-60-22 W5MPERFORATIONS: 11006'-11052' K.B. DISTANCE TO M.P.P.: 11029' K.B.TUBING SIZE: 2.992" I.D.RESERVOIR TEMPERATURE: 246°FGAS PROPERTIES:

G = 0.6997

a.

 $p_c = 801.2 \text{ PSIA}$ $p_{T_c} = 410.9 \text{ }^\circ\text{R}$ MOL.% CO₂ = 2.88MOL.% N₂ = 4.56MOL.% H₂S = 18.89

(Complete gas analysis at the time of test is not available.)

WELL OPERATOR: Amoco Canada Petroleum Company LimitedTYPE OF TEST: Multi-PointDATE OF TEST: September 20, 1970GAS PRODUCED THROUGH: Tubing

WELL B (CONTINUED)

PRODUCTION AND PRESSURE DATA

Total Flow Rate (MMSCF/Day)	Tubing Head		Temp. (°F)	Press. at M.P.P. (PSIA)	Condensate Prod. Rate (bbl/Day)	Water/Gas Ratio (bbl/MMSCF)
	Press. (PSIA)					
6.527	3249.4		132	4249.6	0.0	1.5
9.598	3168.4		140	4205.8	0.0	1.5
12.048	3078.4		153	4163.2	0.0	1.5
14.277	2990.4		158	4123.2	0.0	1.5

WELL C

FIELD: BigstoneWELL NAME: Pan Am. HB, C-1POOL: D-3ALOCATION: 13-35-60-22 W5MPERFORATIONS: 10957'-10974' K.B. DISTANCE TO M.P.P.: 10965' K.B.TUBING SIZE: 2.992" I.D.RESERVOIR TEMPERATURE: 243°FGAS PROPERTIES:

G = 0.702

 $p_c = 798.3 \text{ PSIA}$ $p_c^T = 409.5 \text{ }^\circ\text{R}$ MOL.% CO₂ = 2.78MOL.% N₂ = 4.56MOL.% H₂S = 18.50

(Complete gas analysis at the time of test is not available.)

WELL OPERATOR: Amoco Canada Petroleum Company LimitedTYPE OF TEST: Single-PointDATE OF TEST: October 10, 1972GAS PRODUCED THROUGH: Tubing

WELL C (CONTINUED)

PRODUCTION AND PRESSURE DATA

Total Flow Rate (MMSCF/Day)	Tubing Head		Temp. (°F)	Press. at M.P.P. (PSIA)	Condensate Prod. Rate (bbl/Day)	Water/Gas Ratio (bbl/MMSCF)
	Press. (PSIA)					
15.606	2314.5*	155		3295.5	0.0	5.21
* Dead Weight Test Data						

* Dead Weight Test Data.

WELL D

FIELD: BigstoneWELL NAME: Pan Am. HB,G-2POOL: D-3ALOCATION: 2-25-60-22 W5MPERFORATIONS: 11006'-11052' K.B. DISTANCE TO M.P.P.: 11029' K.B.TUBING SIZE: 2.992" I.D.RESERVOIR TEMPERATURE: 243°FGAS PROPERTIES: $G_v = 0.702$ $p_c = 798.3 \text{ PSIA}$ $p_{T_c} = 409.5 \text{ }^\circ\text{R}$ $\text{MOL.\% CO}_2 = 2.78$ $\text{MOL.\% N}_2 = 3.52$ $\text{MOL.\% H}_2\text{S} = 18.50$

(Complete gas analysis at the time of test is not available.)

WELL OPERATOR: Amoco Canada Petroleum Company LimitedTYPE OF TEST: Single-PointDATE OF TEST: October 10, 1972GAS PRODUCED THROUGH: Tubing

WELL D (CONTINUED)

PRODUCTION AND PRESSURE DATA

Total Flow Rate (MMSCF/Day)	Tubing Head		Temp. (°F)	Press. at M.P.P. (PSIA)	Condensate Prod. Rate (bbl/Day)	Water/Gas Ratio (bbl/MMSCF)
	Press. (PSIA)					
17.359	2.599.5*	170	3690.4	0.0	1.3	

* Dead Weight Test Data.

APPENDIX G

COPIES OF THE FORTRAN COMPUTER PROGRAMS FOR MAIN MODELS AA AND BB

C*****
C***** MAIN MODEL AA *****
C*****

C** THIS MAIN MODEL, 'AA', REQUIRES THE FOLLOWING
C** FUNCTION AND SUBROUTINE PROGRAMS :
C** F, SEARCH, AMEIN, KASEM, BWR, LEE, MOODY,
C** QUON & MISIC . THESE PROGRAMS ARE LISTED
C** IN PAGES TO FOLLOW .

C** HOW TO USE THIS COMPUTER PROGRAM

C ITYPE IS A FLAG INTEGER
C ITYPE=0 : INJECTION PROBLEM
C ITYPE=1 : PRODUCTION PROBLEM
C M IS A FLAG INTEGER
C M=1 : SUPPLY GAS ANALYSIS
C *** MMM IS A FLAG INTEGER
C MMM=0 : ONE-PHASE FLOW
C MMM=1 : TWO-PHASE FLOW

C COMBINATIONS OF THE FLAG INTEGERS ARE AS FOLLOWS:
C ITYPE M MMM DESCRIPTION
C 1 1 1 TWO-PHASE FLOW PRODUCTION
C PROBLEM. SUPPLY GAS COMPOSITION
C AND SG OF FLOWING STREAM
C 1 1 0 ONE-PHASE FLOW PRODUCTION
C PROBLEM. SUPPLY GAS COMPOSITION

C GASUC IS GAS UNIVERSAL CONSTANT
C GC IS CONVERSION FACTOR
C AJ IS CONVERSION FACTOR
C SLOPE IS SLOPE OF PIPE SEGMENTS

C** THE FOLLOWING TERMS BELONG TO PURE GASES
C YMOLF IS MOLE FRACTION
C YMW IS MOLECULAR WEIGHT
C YMW1 IS MOLECULAR WEIGHT
C TEMC IS CRITICAL TEMPERATURE IN DEG. RANKIN
C TEMC1 IS CRITICAL TEMPERATURE IN DEG. RANKIN
C PREC IS CRITICAL PRESSURE IN PSIA
C PREC1 IS CRITICAL PRESSURE IN PSIA
C VOLC IS CRITICAL VOLUME IN CU. FT/LB
C VOLC1 IS CRITICAL VOLUME IN CU. FT/LB
C TNBO1 IS NORMAL BOILING POINT IN DEG. R
C TNBO1 IS NORMAL BOILING POINT IN DEG. R
C SEGM1 IS COLLISION DIAMETER IN ANGSTROMS
C SEGMA IS COLLISION DIAMETER IN ANGSTROMS

C EOK IS A PARAMETER IN LENNARD-JONES POTENTIAL,
 C IT IS (EO /K) IN THE TEXT
 C EOK1 IS THE SAME AS EOK
 C CPOC IS THE FOUR COEFFICIENTS OF THE IDEAL
 C HEAT CAPACITY EQ. OF PURE GASES
 C (KOBE ET AL.)
 C CPOC1 IS THE SAME AS CPOC

 C SG IS GAS GRAVITY
 C PTEMC IS PSEUDOCRITICAL TEMP. IN DEG. RANKIN
 C PPREC IS PSEUDOCRITICAL PRESS. IN PSIA
 C DIAM IS TUBING DIAMETER IN INCHES
 C DEPTH IS LENGTH OF TUBING IN FEET
 C TSOD IS GROUND SURFACE TEMP. IN DEG. F
 C TSD IS TUBING HEAD TEMP. IN DEG. F
 C TDD IS BOTTOM-HOLE TEMP. IN DEG. F
 C PSD IS TUBING HEAD PRESSURE IN PSIA
 C Q IS TOTAL FLOW RATE IN MSCF/DAY

IMPLICIT REAL*8 (A-H,O-Z)

EXTERNAL F

DIMENSION VOLC1(15),TNBO1(15),SEGM1(15),EOK1(15)

DIMENSION YMW1(15),TEMC1(15),PREC1(15)

DIMENSION CPOC1(15,4)

COMMON /BL1/GASUC,GC,AJ,SLOPE/BL2/ITYPE

COMMON /BL10/TNBO1(15)/BL4/SEGMA(15),EOK(15)

COMMON /BL5/YMW(15),TC(15),ZC(15),GAMA(15)

COMMON /BL8/CPOC(15,4)

COMMON /BL11/TEMC(15),PREC(15),VOLC(15)

COMMON /BL12/PTEMC,ZTEMC,PPREC,ZPREC/BL18/M

COMMON /BL13/DIAM/BL14/SG/BL9/YMOLF(15)

COMMON /BL15/Q,TSD,PSD,TDD,DEPTH,TSOD,DL

COMMON /BL30/VMASS

COMMON /BL50/MMH

DATA YMW1/0.44013D 02,0.28013D 02,0.34076D 02,

1 0.16043D 02,0.30070D 02,0.44097D 02,

2 0.58124D 02,0.58124D 02,0.72151D 02,

3 0.72151D 02,0.86178D 02,0.86178D 02,

4 0.10021D 03,0.10021D 03,0.11423D 03/

DATA TEMC1/0.54757D 03,0.22727D 03,0.67237D 03,

1 0.34304D 03,0.54976D 03,0.66568D 03,

2 0.73465D 03,0.76532D 03,0.82877D 03,

3 0.84537D 03,0.89550D 03,0.91337D 03,

4 0.95467D 03,0.97247D 03,0.10239D 04/

DATA PREC1/0.1071D 04,0.4930D 03,0.1306D 04,

1 0.6678D 03,0.7078D 03,0.6163D 03,

2 0.5291D 03,0.5507D 03,0.4904D 03,

3 0.4886D 03,0.4366D 03,0.4369D 03,

4 0.3965D 03,0.3968D 03,0.3606D 03/

DATA VOLC1/0.342D-01,0.514D-01,0.459D-01,

1 0.991D-01,0.788D-01,0.737D-01,

2 0.724D-01,0.702D-01,0.679D-01,

```

3      0.675D-01,0.681D-01,0.688D-01,
4      0.673D-01,0.691D-01,0.690D-01/
DATA TNBO1/0.35037D 03,0.13927D 03,0.38307D 03,
1      0.20098D 03,0.33219D 03,0.41600D 03,
2      0.47057D 03,0.49077D 03,0.54179D 03,
3      0.55659D 03,0.60014D 03,0.61539D 03,
4      0.65376D 03,0.66884D 03,0.71789D 03/
DATA SEGM1/0.3996D 01,0.3798D 01,0.3623D 01,
1      0.3758D 01,0.4443D 01,0.5118D 01,
2      0.5278D 01,0.4687D 01,0.6051D 01,
3      0.5784D 01,0.6256D 01,0.5949D 01,
4      0.7015D 01,0.7004D 01,0.7451D 01/
DATA EOK1/0.1900D 03,0.7140D 02,0.3011D 03,
1      0.1486D 03,0.2157D 03,0.2371D 03,
2      0.3301D 03,0.5314D 03,0.2694D 03,
3      0.3411D 03,0.2472D 03,0.3993D 03,
4      0.2742D 03,0.2882D 03,0.3200D 03/
DATA CPOC1/0.5316D 01, 0.6903D 01, 0.7070D 01,
A      0.4750D 01, 0.1648D 01,-0.9660D 00,
B      -0.1890D 01, 0.9450D 00,-0.2273D 01,
C      0.1618D 01,-0.9081D 00, 0.1657D 01,
D      -0.26069D 01, 0.31899D 01, 0.35844D 01,
E      0.7936111D-02,-0.2085000D-03, 0.1737777D-02,
F      0.6666666D-02, 0.2291111D-01, 0.4043888D-01,
G      0.5519999D-01, 0.4929444D-01, 0.6907777D-01,
H      0.6027777D-01, 0.7843494D-01, 0.7327777D-01,
I      0.9611055D-01, 0.8315605D-01, 0.9502444D-01,
J      -0.2580864D-05, 0.5956790D-06, 0.4209876D-06,
K      0.9351851D-06,-0.4722222D-05,-0.1158950D-04,
L      -0.1695987D-04,-0.1351851D-04,-0.2190432D-04,
M      -0.1655864D-04,-0.2396910D-04,-0.2112345D-04,
O      -0.3089506D-04,-0.2373148D-04,-0.2742592D-04,
P      0.3058984D-09,-0.1176440D-09,-0.1348936D-09,
Q      -0.4509602D-09, 0.2983539D-09, 0.1299725D-08,
R      0.2043895D-08, 0.1433470D-08, 0.2719478D-08,
S      0.1731824D-08, 0.2848422D-08, 0.2362825D-08,
T      0.3834362D-08, 0.2614881D-08, 0.3060012D-08/
GC=0.32174D 02
GASUC=0.1545D 04
AJ=0.77816D 03
SLOPE=0.1D 01
DO 5 I=1,15
YMW(I)=YMW1(I)
TEMC(I)=TEMC1(I)
PREC(I)=PREC1(I)
VOLC(I)=VOLC1(I)
TNBOI(I)=TNBO1(I)
SEGMA(I)=SEGM1(I)
EOK(I)=EOK1(I)
DO 5 J=1,4
5 CPOC(I,J)=CPOC1(I,J)
WRITE(6,208)

```

```

WRITE(6,202) (YMW(I),TEMC(I),PREC(I),VOLC(I),
*TNBOI(I),SEGMA(I),EOK(I),I=1,15)
WRITE(6,213)
WRITE(6,205) ((CPOC(I,J),J=1,4),I=1,15)
WRITE(6,214)

C READ 3 CARDS. GAS COMPOSITION. EACH CARD HAS
C THE MOLE FRACTIONS OF 5 COMPONENTS. THE
C ORDER OF COMPONENTS IS AS FOLLOWS :
C CO2,N2,H2S,C1,C2,C3,I-C4,N-C4,I-C5,N-C5,
C I-C6,N-C6,I-C7,N-C7,N-C8 .
C IF GAS COMPOSITION IS NOT AVAILABLE,
C SUPPLY 3 BLANK CARDS .

READ(5,206) (YMOLF(I),I=1,15)
WRITE(6,215)
WRITE(6,207) (YMOLF(I),I=1,15)

C READ 1 CARD ( COMBINATION OF FLAG INTEGERS
C AND SG )

READ(5,209) ITYPE,M,MMM,SG

C READ 2 CARDS ( WELL PRODUCTION DATA )

READ(5,206) DIAM,DEPTH,TSOD
READ(5,206) Q,TSO,PSD,TDD
IF(ITYPE) 1,1,2
1 WRITE(6,218)
GO TO 3
2 WRITE(6,219)
WRITE(6,208)
3 WRITE(6,216) DIAM,DEPTH
WRITE(6,217) Q,TSO,TSO,TDD,PSD
CALL AMEIN
WRITE(6,227) SG,PTMC,PPREC
WRITE(6,225)
WRITE(6,226)

C CALCULATING MOLAL FLOW RATE
TR=(0.6D 02+0.45967D 03)/ZTEMC
PR=0.1465D 02/ZPREC
CALL BWR (TR,PR,ITER,ZST,DR,CPDM,DZDTR,DZDPR)
VMAS=0.55744134D-02*Q/(ZST*DIAM**2)

C READ 1 CARD ( RANGE OF R* )
C XMIN=1000.0
C XOUT=1010.0
C XMAX=7000.0
C IF CONVERGENCE IS NOT ACHIEVED,THE RANGE
C MAY BE PUT BETWEEN 200.0 & 1000.0 OR
C 7000.0 & 200000.0

```

```

C** REMARK ** FOR SAVING COMPUTER TIME IN CASE R*
C    IS OUTSIDE THE SPECIFIED RANGE,
C    ALWAYS PUT XOUT=XMIN+10.0

```

```

READ(5,201) XMIN, T,XMAX
TEMESP=0.1D 00

```

```

C MINIMIZING FUNCTION F

```

```

CALL SEARCH (F,T,TEMESP,X,XMIN,XMAX,XOUT)
201 FORMAT(7F10.4)
202 FORMAT(10X,7F16)
204 FORMAT(4D15.7)
205 FORMAT(10X,4D20)
206 FORMAT(5F15.8)
207 FORMAT(18X,F20.5)
208 FORMAT('1')
209 FORMAT(3I5,5X,F10.4)
211 FORMAT(10X,5F15.7)
213 FORMAT(////////,29X,'A',19X,'B',19X,'C',19X,'D',/)
214 FORMAT(//,47X,'CPO(I)=A+B*T+C*T**2+D*T**3')
215 FORMAT('1',////////,27X,'MOLE FRACTION',/)
216 FORMAT(////////,10X,'TUBING I. D.',27X,'=',F10.3,
A ' INCHES',/,
B10X,'THE LENGTH OF FLOW STRING',/,
C F10.3,' FEET')
217 FORMAT(10X,'FLOW RATE',30X,'=',F10.3,
A ' MSCF/DAY',/,
B10X,'THE AVE. TEMP. OF THE GROUND SURFACE',/,
C F10.3,' F',/,
D10X,'THE TEMP. OF FLOWING STREAM AT SURFACE',/,
E F10.3,' F',/,
F10X,'THE TEMP. OF THE PRODUCING FORMATION',/,
G F10.3,' F',/,
H10X,'THE WELL HEAD TUBING PRESSURE',/,
I F10.3,' PSIA',/)
218 FORMAT(10X,'INJECTION PROBLEM')
219 FORMAT(10X,'PRODUCTION PROBLEM')
225 FORMAT(10X,'THE CALCULATED BOTTOM-HOLE PRESSURE',
1 ' IS THAT WHEN THE ABSOLUTE DIFFERENCE BETWEEN',
2 //,10X,'EARTH AND CALCULATED FLUID TEMPERATURES',
3 ' AT THE BOTTOM OF WELL IS WITHIN 0.1 DEG. F. ')
226 FORMAT(//,10X,'FIRST COLUMN IS DEPTH IN FEET',/,
1 10X,'SECOND COL. IS EARTH TEMP. IN DEG. F.',/,
2 10X,'THIRD COL. IS CALC. FLUID TEMP. IN DEG. F.',
3 //,10X,'FOURTH COL. IS CALC. PRESS. IN PSIA',/,
4 //,10X,'FIFTH COL. IS R*',/)
227 FORMAT(10X,'GAS GRAVITY, PSEUDO REDUCED TEMP.',
1 'AND PSEUDO REDUCED PRESS. ARE :',//,10X,3F10.4)
STOP
END

```

FUNCTION F(X)

C**** THIS FUNCTION SUBROUTINE IS TO BE
C**** USED WITH MAIN MODEL AA ONLY

C F IS THE FUNCTION TO BE MINIMIZED
C E IS ABSOLUTE ROUGHNESS OF TUBING IN INCHES
C DIAM IS TUBING DIAMETER IN INCHES
C DL IS MONITORED INTEGRATION STEP SIZE IN FEET
C TGEOG IS GEOTHERMAL GRADIENT , DEG. F/ FOOT
C PRE IS PRESSURE IN PSIA
C TEM IS TEMPERATURE IN DEG. RANKIN
C DENS IS DENSITY OF NAT. GAS IN GM/C.C.
C REN IS REYNOLDS NUMBER
C PRAND IS PRANDTL NUMBER
C STANT IS STANTON NUMBER
C GFTCO IS GAS-FILM HEAT-TRANSFER COEFFICIENT
C OHTCO IS THE OVERALL HEAT-TRANSFER COEFFICIENT

IMPLICIT REAL*8 (A-H,O-Z)
DIMENSION PRY(5),TRY(5),PRK(5),TRK(5)
COMMON /BL3/AVEMW,BL13/DIAM,BL14/SG
COMMON /BL12/PTEMC,ZTEMC,PPREC,ZPREC
COMMON /BL15/Q,TSD,PSD,TDD,DEPTH,TSOD,DL
COMMON /BL6/CONST,BL18/M
COMMON /BL30/VMASS
COMMON /BL50/MMM
D1=0.1D-01
D2=0.1D-01
DL=C.1D 02
E=C.6D-03
RSTAR=X
RR=E/DIAM
TGEOG=(TDD-TSOD)/DEPTH

C THE RUNGE-KUTTA-MERSON INTEGRATION ALGORITHM

TEM=TSOD+0.45967D 03
PRE=PSD
TS=TSOD
TR=TEM/PTEMC
PR=PRE/PPREC
TEN=TR
PRN=PR
TL=C.0D C0
WRITE(6,211) TL,TS,TSD,PRE,RSTAR
1 DO 14 J1=1,5
IF(J1-1) 2,2,3
2 PRY(1)=PRN
TRY(1)=TRN
TS=TSOD+TGEOG*TL
3 IF(J1-2) 11,4,5

```

4  PRY (2) = PRN + PRK (1)
   TRY (2) = TRN + TRK (1)
   TS = TSOD + TGEOG * (TL + DL / 0.3D 01)
5  IF (J1-3) 11, 6, 7
6  PRY (3) = PRN + 0.5D 00 * PRK (1) + 0.5D 00 * PRK (2)
   TRY (3) = TRN + 0.5D 00 * TRK (1) + 0.5D 00 * TRK (2)
   TS = TSOD + TGEOG * (TL + DL / 0.3D 01)
7  IF (J1-4) 11, 8, 9
8  PRY (4) = PRN + 0.375D 00 * PRK (1) + 0.1125D 01 * PRK (3)
   TRY (4) = TRN + 0.375D 00 * TRK (1) + 0.1125D 01 * TRK (3)
   TS = TSOD + TGEOG * (TL + DL / 0.2D 01)
9  IF (J1-5) 11, 10, 14
10 ABC1 = PRN + 0.15D 01 * PRK (1)
   ABC2 = -0.45D 01 * PRK (3) + 0.6D 01 * PRK (4)
   PRY (5) = ABC1 + ABC2
   ABC1 = TRN + 0.15D 01 * TRK (1)
   ABC2 = -0.45D 01 * TRK (3) + 0.6D 01 * TRK (4)
   TRY (5) = ABC1 + ABC2
   TS = TSOD + TGEOG * (TL + DL)
   PRJAM = PRY (5)
   TRJAM = TRY (5)
11 PR = PRY (J1)
   TR = TRY (J1)
   TRS = (TS + 0.45967D 03) / PTEM
   PRE = PR * PPREC
   TEM = TR * PTEM

```

C CALCULATING IDEAL HEAT CAPACITY OF NAT. GAS

```

CALL KASEM (TR, TEM, SG, M, CPOM)
TRC = TEM / ZTEM
PRC = PRE / ZPREC

```

C CALCULATING COMPRESSIBILITY FACTOR, ITS PARTIAL C DERIVATIVES W.R.T. PR & TR, REDUCED C DENSITY, AND HEAT CAPACITY DEPARTURE C GF NAT. GAS —

```

CALL BWH (TRC, PRC, ITER, Z, DR, CPDM, DEDTR, DZDPR)
IF (ITER) 22, 22, 12
12 DENS = 0.149256D - 02 * AVEMW * PRE / (Z * TEM)
   CPM = CPDM + CPOM

```

C CALCULATING VISCOSITY OF NAT. GAS IN C.P.

```

CALL LEE (TEM, AVEMW, DENS, VISM)
REN = 0.20014855D 02 * Q * SG / (VISM * DIAM)

```

C CALCULATING THE HYDRAULIC FRICTION FACTOR

```

CALL MOODY (RR, REN, FRI)

```

C CALCULATING THERMAL COND. OF NAT. GAS

CALL MISIC (TEM,DR,CONST,THKM)

C CALCULATING THE PRANDTL NUMBER

PRAND=CPM*(VISM/0.1D 03)/(AVEMW*THKM)

C CALCULATING THE STANTON NUMBER AND THE
C GAS-FILM HEAT-TRANSFER COEFFICIENT

B1=FF/0.8D 01

B2=DSQRT(B1)

ABC3=DSQRT(PRAND)

ABC4=(B1*REN*REN)**0.27D 00

PSI=ABC3*(0.155D 00*ABC4+C.1D 01/B2)

STANT=B2/(PSI-0.45D 01)

GFTCO=VMASS*CPM*STANT

U1=0.1D 01/GFTCO+RSTAR

OHTCO=0.1D 01/U1

C SOLVING FOR THE RATE OF CHANGE OF REDUCED
C PRESSURE AND TEMPERATURE WITH LENGTH

CALL QUON(TEM,PRE,TR,PR,TRS,Z,DZDPR,DZDTR,
1FF,CPM,OHTCO,VMASS,DPRDL,DTRDL,NASER)

IF(NASER-1) 22,13,13

13 PRK(J1)=DL*DPRDL/0.3D 01

TRK(J1)=DL*DTRDL/0.3D 01

14 CONTINUE

C MONITORING THE INTEGRATION STEP SIZE

ABC5=PRK(1)-0.45D 01*PRK(3)

ABC6=0.4D 01*PRK(4)-0.5D 00*PRK(5)

EPR=ABC5+ABC6

ABC5=TRK(1)-0.45D 01*TRK(3)

ABC6=0.4D 01*TRK(4)-0.5D 00*TRK(5)

ETR=ABC5+ABC6

ETEM=DABS(ETR*PTMC)

EPRE=DABS(EPR*PPREC)

IF(EPRE-D1) 15,15,16

15 IF(ETEM-D2) 17,17,16

16 DL=DL/0.2D 01

GO TO 1

17 PR=PRJAM-EPR

TR=TRJAM-ETR

TI=TL+DL

IF(EPRE-D1/0.16D 02) 18,20,20

18 IF(ETEM-D2/0.16D 02) 19,20,20

19 DL=0.2D 01*DL

20 PRN=PR

```
TRN=TR
PRE=PRN*PPREC
TEM=TRN*PTMC
TED=TEM-C.45967D 03
WRITE(6,211) TL,TS,TED,PRE,RSTAR
DEL=DEPTH-TL
IF(DABS(DEL)-0.1D-04) 23,23,25
25 IF(DEL-DL) 21,1,1
21 DL=DEL
GO TO 1
22 WRITE(6,227)
23 TED=TEM-0.45967D 03
TEMTL=TED
F=TDD-TEMTL
211 FCRMAT(10X,5F10.2)
227 FORMAT(//,10X,'*** THE COURSE OF ITERATIONS IS'
*, ' INTERRUPTED ***',//,7X,'*** THE PROGRAM IS'
*, ' GOING TO ASSUME ANOTHER "RSTAR"***')
RETURN
END
```

SUBROUTINE AMEIN

C**** THIS SUBROUTINE CAN BE USED WITH
C**** MAIN MODEL AA ONLY

C THIS SUBROUTINE CALCULATES THE MOLAL AVERAGE
C PROPERTIES, THE WICHERT & AZIZ CORRECTED
C PSEUDOCRITICAL PROPERTIES OF NAT. GAS ,
C AND OTHER PROPERTIES

C THE FOLLOWING TERMS BELONG TO PURE GASES

C S IS THE SUTHERLAND CONSTANT IN DEG. K
C VC IS CRITICAL MOLAL VOLUME, C.C./GM MOLE
C VOLC IS CRITICAL VOLUME , CU. FT/ LB.
C ZC IS CRITICAL COMPRESSIBILITY FACTOR
C GAMA. IS CONSTANT DEFINED BY EQ. 58 IN TEXT

C THE FOLLOWING TERMS BELONG TO NAT. GAS

C AVMW IS AVERAGE MOLECULAR WEIGHT
C SG IS GAS GRAVITY
C PPREC IS PSEUDOCRITICAL PRESSURE IN PSIA
C PPC IS PSEUDOCRITICAL PRESSURE IN ATMOSPHERES
C PTEMC IS PSEUDOCRITICAL TEMPERATURE, DEG. R
C PTC IS PSEUDOCRITICAL TEMPERATURE, DEG. K
C GAM IS CONSTANT DEFINED BY EQ. 58 IN TEXT
C ZTEMC IS CORRECTED PSEUDOCRITICAL TEMPERATURE
C ZPREC IS CORRECTED PSEUDOCRITICAL PRESSURE

IMPLICIT REAL*8 (A-H, O-Z)

DIMENSION PC (15)

COMMON /BL3/AVMW/BL4/SEGMA (15) , EOK (15)

COMMON /BL5/YMW (15) , TC (15) , ZC (15) , GAMA (15)

COMMON /BL9/YMOLF (15) /BL10/TNBOI (15) /BL20/S (15)

COMMON /BL11/TEMC (15) , PREC (15) , VOLC (15) /BL6/CONST

COMMON /BL12/PTMC, ZTEMC, PPREC, ZPREC

COMMON /BL14/SG/BL18/M

COMMON /BL50/MMH

COST1=0.15D 01/0.18D 01

CCST2=0.1D 01/C.6D 01

COST3=0.2D 01/0.3D 01

Y11=YMOLF (1)

Y22=YMOLF (2)

Y33=YMOLF (3)

IF (M-1) 3, 1, 1

1 AVMW=0.0D 00

PTMC=0.0D 00

PPREC=0.0D 00

DO 2 I=1, 15

S (I)=COST1*TNBOI (I)

PC (I)=PREC (I)/0.146964757D 02

```

TC(I)=TEMC(I)/0.18D 01
VC=0.623826896D 02*YMW(I)*VOLC(I)
ZC(I)=PC(I)*VC/(0.8206D 02*TC(I))
Y1=TC(I)**COST2
Y2=PC(I)**COST3
GAMA(I)=0.1D 07*Y1*DSQRT(YMW(I))/Y2
AVEMW=AVEMW+YMOLF(I)*YMW(I)
PTMC=PTMC+YMOLF(I)*TEMC(I)
2 PPREC=PPREC+YMOLF(I)*PREC(I)
IF(MMM-1) 10,3,3
10 SG=AVEMW/0.2897D 02
GO TO 7
3 AVEMW=SG*0.2897D 02

```

C THE BROWN ET AL AND CARR ET AL CORRELATIONS
C FOR THE PSEUDOCRITICAL PROPERTIES OF
C SWEET AND SOUR NAT. GASES

```

PTMC=0.171137D 03+0.313725D 03*SG
TCCOR=-0.8D 02*Y11-0.25D 03*Y22+0.13D 03*Y33
PTMC=PTMC+TCCOR
IF(SG-0.85D 00) 4,4,5
4 PPREC=C.6951D 03-C.4D 02*SG
GO TO 6
5 PPREC=C.704396D 03-0.51724D 02*SG
6 PCCOR=0.44D 03*Y11-0.17D 03*Y22+0.6D 03*Y33
PPREC=PPREC+PCCOR

```

C THE WICHERT & AZIZ CORRECTED PSEUDOCRITICALS
C DUE TO PRESENCE OF CO2 AND H2S

```

7 S1=YMOLF(1)+YMOLF(3)
S2=YMOLF(3)
E31=0.12D 03*(S1**0.9D 00-S1**0.16D 01)
E32=0.15D 02*(DSQRT(S2)-S2**4)
E3=E31+E32
S4=PTMC+S2*(0.1D 01-S2)*E3
ZTEMC=PTMC-E3
ZPREC=PPREC*ZTEMC/S4
PTC=PTMC/0.18D 01
PPC=PPREC/0.14696476D 02
GAM=PTC**COST2*DSQRT(AVEMW)/PPC**COST3
CCNST=GAM*0.14348907D 06
RETURN
END

```

SUBROUTINE MISIC (TEM,DR,CONST,THKM)

C**** THIS SUBROUTINE CAN BE USED WITH
C**** MAIN MODEL AA ONLY

C THIS SUBROUTINE CALCULATES THE THERMAL
C CONDUCTIVITY OF GAS MIXTURES

C** THE MISIC AND THODOS METHODS ARE USED TO
C CALCULATE THKO,THKMO,THKMD,AND THKM

C TEM IS TEMP. IN DEGREE RANKIN
C DR IS REDUCED DENSITY OF GAS MIXTURE
C THKMO IS THERMAL COND. OF GAS MIX. AT 0 PRESS.
C THKMD IS THERMAL COND. DEPARTURE
C THKM IS THERMAL CONDUCTIVITY OF GAS MIXTURE
C IN CAL/CM SEC K
C CONST IS $10^{**8} \cdot 0.27^{**5} \cdot \text{GAM}$

C** THE FOLLOWING TERMS BELONG TO PURE GASES
C YMW IS MOLECULAR WEIGHT
C TC IS CRITICAL TEMPERATURE IN DEG. KELVIN
C ZC IS CRITICAL COMPRESSIBILITY FACTOR
C S IS SUTHERLAND'S CONSTANT IN DEG. KELVIN
C CPO IS HEAT CAPACITY AT ZERO PRESSURE
C IN CAL/GM MOLE K
C THKO IS THERMAL COND. AT 0 PRESSURE
C YMOLF IS MOLE FRACTION

IMPLICIT REAL*8 (A-H,O-Z)
DIMENSION A(15,15),SS(15,15),VISO(15),
* THKO(15),F(2),TRED(15)
COMMON /BL4/SEGMA(15),EOK(15)
COMMON /BL5/YMW(15),TC(15),ZC(15),GAMA(15)
COMMON /BL7/CPO(15)/BL9/YMOLF(15)/BL20/S(15)
UNITY=0.1D 01
FRAC=2.0D 00/3.0D 00
TK=TEM/0.18D 01
TKK=TK
DO 8 I=1,15
TRED(I)=TK/TC(I)
DO 5 K=1,2
IFLAG=K
T=TK/EOK(I)
IF(T-0.2D 01) 1,2,2
1 F(K)=0.634D 00*(UNITY+0.475D 00*DLOG(T))
GO TO 3
2 F(K)=0.697D 00*(UNITY+0.323D 00*DLOG(T))
3 IF(TRED(I)-0.2D 01) 4,6,6
4 TK=0.2D 01*TC(I)
5 CONTINUE
Y=DSQRT(YMW(I)*TK)

```

6 VISO(I)=0.26693D-02*Y*F(IFLAG)/SEGMA(I)**2
  IF(IFLAG-1) 8,8,7
7 TK=TKK
  VISO(I)=VISO(I)*0.5D 00*TRED(I)*F(1)/F(2)
8 CONTINUE
  DO 11 I=1,15
    IF(I-3) 9,9,10
9 A1=0.195D 03*ZC(I)-0.3194D 02
  A2=0.1683D 02-0.825D 02*ZC(I)
  A3=0.1524D 01-0.28D 01*ZC(I)
  A4=(A1*TRED(I)+A2)**A3
  THKO(I)=CPO(I)**0.75D 00*A4/GAMA(I)
  GO TO 11
10 B1=0.1452D 02*TRED(I)-0.514D 01
  THKO(I)=CPO(I)*B1**FRAC/GAMA(I)
11 CONTINUE
  DO 12 I=1,15
    DO 12 J=1,15
      SS(I,J)=DSQRT(S(I)*S(J))
      C1=(UNITY+SS(I,J)/TK)/(UNITY+S(I)/TK)
      C2=(VISO(I)/VISO(J))*(YMW(J)/YMW(I))**0.75D 00
      C3=(UNITY+S(I)/TK)/(UNITY+S(J)/TK)
      A(I,J)=0.25D 00*C1*(UNITY+DSQRT(C2*C3))**2
12 CONTINUE
  THKMO=0.0D 00
  DO 15 I=1,15
    IF(YMOLF(I)-0.1D-06) 15,15,25
25 DUMY=0.1D 01
  DO 14 J=1,15
    IF(YMOLF(J)-0.1D-06) 14,14,26
26 IF(I-J) 13,14,13
13 DUMY=DUMY+A(I,J)*YMOLF(J)/YMOLF(I)
14 CCNTINUE
  THKMO=THKMO+THKO(I)/DUMY
15 CONTINUE
  IF(DR-0.5D 00) 16,16,17
16 X1=DEXP(0.535D 00*DR)
  THKMD=0.14D 02*(X1-UNITY)/CONST
  GO TO 20
17 IF(DR-0.2D 01) 18,18,19
18 X1=DEXP(0.67D 00*DR)
  THKMD=0.131D 02*(X1-0.1069D 01)/CONST
  GO TO 20
19 X1=DEXP(0.1155D 01*DR)
  THKMD=C.2976D 01*(X1+0.2016D 01)/CONST
20 THKM=THKMO+THKMD
  RETURN
  END

```

SUBROUTINE SEARCH(F,EPS,X,YMIN,YMAX,RSTOUT)

C**** THIS SUBROUTINE CAN BE USED WITH
 C**** EITHER MAIN MODEL AA OR BB

C THIS SUBROUTINE SEARCHES FOR THE VALUE OF
 C R* OR R** WHICH MINIMIZES THE FUNCTION 'F'

```

    IMPLICIT REAL*8 (A-H,O-Z)
    EXTERNAL F
    DIMENSION RST(3),TEST(2),DELTA(3)
    RST(1)=YMIN
    RST(3)=YMAX
    ML=0
30  IF(ML-17) 1,1,21
21  WRITE(6,225)
    GO TO 13
    1 RST(2)=(RST(1)+RST(3))/0.2D 01
      IF(RST(2)-RSTOUT) 13,13,2
    2 DO 6 K3=1,3
      IF(ML) 4,4,3
    3  IF(K3-2) 6,4,6
    4  RSTAR=RST(K3)
      X=RSTAR
      DIFF=F(X)
      IF(DABS(DIFF)-EPS) 13,13,5
    5  DELTA(K3)=DIFF
    6  CONTINUE
      TEST(1)=DELTA(1)*DELTA(2)
      TEST(2)=DELTA(3)*DELTA(2)
      DESCN=DELTA(1)*DELTA(3)
      IF(ML) 23,23,24
23  IF(DESCN) 24,24,25
25  WRITE(6,226)
    GO TO 13
24  ML=ML+1
    DO 8 K5=1,2
      IF(TEST(K5)) 7,13,8
    7  ITEST=K5
    GO TO 9
    8  CONTINUE
    9  IF(ITEST-1) 10,10,11
10  RST(3)=RST(2)
    DELTA(3)=DELTA(2)
    GO TO 30
11  IF(TEST(2)) 12,13,12
12  RST(1)=RST(2)
    DELTA(1)=DELTA(2)
    GO TO 30
225 FORMAT(/,20X,'*** NO SOLUTION IS ',
  *'REACHED WITH 20 ITERATIONS ***')
226 FORMAT(/,20X,'*** NO SOLUTION IS ',

```

*'PROVIDED TO THIS PROBLEM ***')
13 RETURN
END

SUBROUTINE KASEM (TR, TEM, SG, M, CPOM)

C**** THIS SUBROUTINE CAN BE USED WITH
 C**** EITHER MAIN MODEL AA OR BB

C TR IS REDUCED TEMP.
 C TEM IS TEMPERATURE IN DEGREE RANKIN
 C SG IS GAS GRAVITY
 C M IS A FLAG INTEGER
 C M=0 : SUPPLY SG, H₂S, CO₂, AND N₂
 C M=1 : SUPPLY GAS ANALYSIS
 C CPOM IS HEAT CAPACITY AT ZERO PRESS. OF NAT. GAS
 C CPC IS HEAT CAPACITY AT ZERO PRESS. OF PURE GAS
 C YMOLF IS MOLE FRACTION

IMPLICIT REAL*8 (A-H, O-Z)

DIMENSION A(14)

CCOMON /BL7/CPO(15)/BL8/CPOC(15,4)/BL9/YMOLF(15)

DATA A/0.5596695D 01, -0.2233480D 01,
 1 0.807265D 00, -0.100390D 01, 0.314160D 01,
 2 0.5758700D 01, 0.425860D 00, 0.124323D -02,
 3 -0.40560D -01, 0.100889D -02, 0.3623622D 00,
 4 -0.4661581D 00, 0.9757D -03, 0.27082167D -02/

IF (M-1) 4, 1, 1

1 CPOM=0.0D 00

DO 3 I=1, 15

CPO(I)=0.0D 00

DO 2 J=1, 4

2 CPO(I)=CPO(I)+CPOC(I,J)*TEM** (J-1)

3 CPOM=CPOM+CPO(I)*YMOLF(I)

GO TO 5

4 B=A(1)+A(2)*SG+A(3)*SG**2

C=A(4)+A(5)*SG+A(6)*SG**2

CPCM=B+C*TR

FCO2=0.1D 01+YMOLF(1)*(A(7)+A(8)*TEM)

FH2S=0.1D 01+YMOLF(3)*(A(9)+A(10)*TEM)

Y1=A(11)*YMOLF(2)+A(12)*YMOLF(2)**2

Y2=A(13)*YMOLF(2)+A(14)*YMOLF(2)**2

FN2=0.1D 01+Y1+Y2*TEM

CPOM=CPOM/(FCO2*FN2*FH2S)

5 RETURN

END

SUBROUTINE BWR (TR, PR, ITER, Z, DR, CPDM, DZDTR, DZDPR)

C**** THIS SUBROUTINE CAN BE USED WITH
C**** EITHER MAIN MODEL AA OR BB

C TR IS THE PSEUDO-REDUCED TEMPERATURE
C PR IS THE PSEUDO-REDUCED PRESSURE
C Z IS THE GAS COMPRESSIBILITY FACTOR
C DR IS THE REDUCED DENSITY OF GAS
C DZDTR IS THE PARTIAL DERIVATIVE OF Z W.R.T. TR
C DZDPR IS THE PARTIAL DERIVATIVE OF Z W.R.T. PR
C CPDM IS THE HEAT CAPACITY DEPARTURE
C ITER IS AN ITERATION FLAG
C ITER=0 EITHER OR BOTH TR AND PR ARE
C OUTSIDE RANGE OF CORRELATION
C ITER>0 NUMBER OF ITERATIONS

IMPLICIT REAL*8 (A-H, O-Z)

DIMENSION A(8)

DATA A/C.31506237D C0, 1.0467099D 00,

A 0.57832729D 00, 0.53530771D 00, 0.61232032D C0,

B 0.10488813D C0, 0.68157C01D 00, 0.68446549D 00/

ITER=0

DR=C.1D C1

IF (TR-0.1C5D 01) 10, 1, 1

1 IF (TR-0.3D 01) 2, 2, 1C

2 IF (PR-0.3C0D 02) 3, 3, 10

3 DO 8 ITER=1, 30

DR2=DR*DR

T1=(A(1)*TR-A(2)-A(3)/TR**2)*DR

T2=(A(4)*TR-A(5))*DR2

T3=A(5)*A(6)*DR**5

T4=A(7)*DR2/TR**2

T5=A(8)*DR2

T6=DEXP(-T5)

P=(TR+T1+T2+T3)*DR+T4*DR*(0.1D 01+T5)*T6

DP1=TR+0.2D 01*T1+0.3D C1*T2+0.6D 01*T3

DP2=T4*T6*(0.3D 01+0.3D 01*T5-0.2D 01*T5*T5)

DP=DP1+DP2

DR1=DR-(P-0.270D 00*PR)/DP

IF (DR1) 4, 4, 5

4 DR1=0.5D C0*DR

5 IF (DR1-0.22D 01) 7, 7, 6

6 DR1=DR+0.9D C0*(0.22D 01-DR)

7 IF (DABS(DR-DR1)-0.1D-04) 9, 8, 8

8 DR=DR1

9 Z=C.270D C0*PR/(DR1*TR)

T7=(A(1)*TR+0.2D 01*A(3)/TR**2)*DR

T8=A(4)*TR*DR**2

T9=A(3)*DR/TR**2

T10=A(7)/A(8)/TR**2

R1=DP/TR

R2=(TR+T7+T8-0.2D 01*T4*T6*(0.1D 01+T5))/TR

P2=R2*DR

P3=0.270D 00/DP

Y1=0.3D 01*T6*(0.2D 01*T10+T4)

R3=(-TR+0.6D 01*(T9-T10)+Y1)/TR

P4=P2/DP

DZDTR=Z*(P4/DR-0.1D 01/TR)

DZDPR=Z*(0.1D 01/PR-P3/DR)

CPDM=0.198654284D 01*(R3+R2**2/R1)

10 RETURN

END

SUBROUTINE LEE(TEM,AVEMW,DENS,VISM)

C**** THIS SUBROUTINE CAN BE USED WITH
 C**** EITHER MAIN MODEL AA OR BB

C THIS SUBROUTINE IS THE LEE,ET AL. VISCOCITY EQ.
 C TEM IS TEMPERATURE IN DEGREE RANKIN
 C AVEMW IS AVERAGE MOLECULAR WEIGHT
 C DENS IS THE GAS DENSITY IN GRAM/C.C.
 C VISM IS THE GAS VISCOCITY IN C.P.

IMPLICIT REAL*8 (A-H,O-Z)

DIMENSION V(11)

DATA V/0.94D 01,0.2D-01,0.15D 01,0.209D 03,

1 0.19D 02,0.35D 01,0.986D 03,0.1D-01,

2 0.24D 01,0.2D 00,0.1D 05/

VK1=(V(1)+V(2)*AVEMW)*TEM**V(3)

VK2=V(4)+V(5)*AVEMW+TEM

VK=VK1/VK2

VX=V(6)+V(7)/TEM+V(8)*AVEMW

VY=V(9)-V(10)*VX

VY1=VX*DENS**VY

VISM=VK*DEXP(VY1)/V(11)

RETURN

END

SUBROUTINE MOODY (RR,REN,FF,I)

C**** THIS SUBROUTINE CAN BE USED WITH
 C**** EITHER MAIN MODEL AA OR BB

C RR IS THE RELATIVE ROUGHNESS
 C REN IS REYNOLDS NUMBER
 C RENCER IS THE BOUNDARY BETWEEN THE
 C TRANSITION AND ROUGH PIPE REGIONS
 C FF IS THE HYDRAULIC FRICTION FACTOR
 C I IS NO. OF ITERATIONS REQUIRED
 C I=0 ROUGH PIPE REGION
 C I>0 TRANSITION REGION .

IMPLICIT REAL*8 (A-H,O-Z)

DIMENSION G(15)

I=0

C=DLOG(0.1D 02)

IF(RR) 2,2,1

1 RENCER=C.32D 04/RR

IF(REN-RENCER) 2,6,6

C THE COLEBROOK TRANSITION LAW

2 X=C.8D-01

DO 4 I=1,15

X1=RR/0.37D 01+0.251D 01/(REN*X)

X2=DLOG10(X1)

FX=C.1D 01+0.2D 01*X*X2

FXD=0.2D 01*X2-0.502D 01/(REN*X*X1*C)

G(I)=X-FX/FXD

X=G(I)

IF(I-1) 4,4,3

3 IF(DABS(G(I)-G(I-1))-0.1D-05) 5,5,4

4 CCNTINUE

5 FF=X*X

GO TO 7

C VON-KARMAN'S ROUGH PIPE FORMULA

6 FF=C.1D 01/(0.2D 01*DLOG10(0.37D 01/RR))**2

7 CCNTINUE

RETURN

END

SUBROUTINE QUON(TEM,PRE,TR,PR,TRS,Z,DZDPR,
1DZDTR,FF,CPM,OHTCO,VMASS,DPRDL,DTRDL,NASER)

C**** THIS SUBROUTINE CAN BE USED WITH
C**** EITHER MAIN MODEL AA OR BB

C THIS SUBROUTINE SOLVES FOR THE RATE OF CHANGE
C OF REDUCED PRESSURE & TEMP. WITH LENGTH
C (SEE EQUATIONS 6 & 7 IN TEXT)

IMPLICIT REAL*8 (A-H,O-Z)
COMMON /BL1/GASUC,GC,AJ,SLOPE/BL2/ITYPE
COMMON /BL3/AVEMW/BL13/DIAM
NASER=1
O=SLOPE
BETA=0.1D 01
DIAN=DIAM/0.12D 02
A1=BETA*VMASS**2*TEM*GASUC
PRF=PR*0.144D 03
A2=GC*AJ*PRF**2
A3=A1/A2
A4=A3*TEM*GASUC
A5=Z/PR-DZDPR
A6=Z/TR+DZDTR
ASTAR=-A1*A5*PRF+GC*PRF**3/(PR*AVEMW)
BSTAR=A1*A6*PRF
CSTA1=-A2*PRF*O/(Z*TEM*AJ*GASUC)
CSTA2=-A1*Z*PRF*FF/(0.2D 01*DIAN*BETA)
CSTAR=CSTA1+CSTA2
DSTAR=A4*A5*Z+TEM*TR*GASUC*DZDTR/(AJ*PR*AVEMW)
ESTAR=-A4*A6*Z-TEM*CPM/(TR*AVEMW)
CHRIS=VMASS*DIAN
FSTA1=CHRIS*O
FSTA2=0.4D 01*OHTCO*TEM*AJ*(TRS-TR)/(TR*AVEMW)
FSTAR=(FSTA1-FSTA2)/(CHRIS*AJ)
A7=ASTAR*ESTAR-BSTAR*DSTAR
IF(A7) 1,3,1
1 DPRDL=(CSTAR*ESTAR-BSTAR*FSTAR)/A7
DTRDL=(ASTAR*FSTAR-CSTAR*DSTAR)/A7
IF(ITYPE-1) 4,2,2
2 DPRDL=-DPRDL
DTRDL=-DTRDL
GO TO 4
3 NASER=0
4 RETURN
END

```

C*****
C***** MAIN MODEL BB *****
C*****
C**  THIS MAIN MODEL, 'BB', REQUIRES THE FOLLOWING
C**  FUNCTION AND SUBROUTINE PROGRAMS:
C**  F & AMEIN (LISTED IN PAGES TO FOLLOW),
C**  SEARCH, KASEM, BWR, LEE, MOODY & QUINN (LISTED
C**  IN PREVIOUS PAGES) .

```

```

C**  HOW TO USE THIS COMPUTER PROGRAM

```

```

C  ITYPE IS A FLAG INTEGER
C    ITYPE=0 : INJECTION PROBLEM
C    ITYPE=1 : PRODUCTION PROBLEM
C  M  IS A FLAG INTEGER
C    M=0 : SUPPLY SG, H2S, CO2, AND N2
C    M=1 : SUPPLY GAS ANALYSIS
C  MM IS A FLAG INTEGER
C    MM=0 : DO NOT SUPPLY PSEUDOCRITICALS
C    MM=1 : SUPPLY PSEUDOCRITICALS
C  MMM IS A FLAG INTEGER
C    MMM=0 : ONE-PHASE FLOW
C    MMM=1 : TWO-PHASE FLOW

```

```

C  COMBINATIONS OF THE FLAG INTEGERS ARE AS FOLLOWS:
C  ITYPE  M  MM  MMM  DESCRIPTION

```

ITYPE	M	MM	MMM	DESCRIPTION
1	1	1	1	TWO-PHASE FLOW PRODUCTION PROBLEM. SUPPLY GAS COMPOSITION AND SG OF FLOWING STREAM
1	1	1	0	ONE-PHASE FLOW PRODUCTION PROBLEM. SUPPLY GAS COMPOSITION
1	0	0	1	TWO-PHASE FLOW PRODUCTION PROBLEM. SUPPLY SG OF FLOWING STREAM, CO2, N2 & H2S
1	0	0	0	ONE-PHASE FLOW PRODUCTION PROBLEM. SUPPLY SG, CO2, N2 & H2S
1	0	1	0	ONE-PHASE FLOW PRODUCTION PROBLEM. SUPPLY SG, CO2, N2, H2S AND THE REPORTED PTMC & PPREC

```

C  GASUC IS GAS UNIVERSAL CONSTANT
C  GC    IS CONVERSION FACTOR
C  AJ    IS CONVERSION FACTOR
C  SLOPE IS SLOPE OF PIPE SEGMENTS

```

```

C**  THE FOLLOWING TERMS BELONG TO PURE GASES
C  YMOLF IS MOLE FRACTION
C  YMW   IS MOLECULAR WEIGHT
C  YMW1  IS MOLECULAR WEIGHT
C  TEMC  IS CRITICAL TEMPERATURE IN DEG. RANKIN

```

```

C TEMC1 IS CRITICAL TEMPERATURE IN DEG. RANKIN
C PREC IS CRITICAL PRESSURE IN PSIA
C PREC1 IS CRITICAL PRESSURE IN PSIA
C CPOC IS THE FOUR COEFFICIENTS OF THE IDEAL
C HEAT CAPACITY EQ. OF PURE GASES
C (KOBÉ ET AL.)
C CPOC1 IS THE SAME AS CPOC
C
C SG IS GAS GRAVITY
C PTEMC IS PSEUDOCRITICAL TEMP. IN DEG. RANKIN
C PPREC IS PSEUDOCRITICAL PRESS. IN PSIA
C DIAM IS TUBING DIAMETER IN INCHES
C DEPTH IS LENGTH OF TUBING IN FEET
C TSOD IS GROUND SURFACE TEMP. IN DEG. F
C TSD IS TUBING HEAD TEMP. IN DEG. F
C TDD IS BOTTOM-HOLE TEMP. IN DEG. F
C PSD IS TUBING HEAD PRESSURE IN PSIA
C Q IS TOTAL FLOW RATE IN MSCF/DAY

```

```

IMPLICIT REAL*8 (A-H,O-Z)
EXTERNAL F
DIMENSION YMW(15),TEMC1(15),PREC1(15)
DIMENSION CPOC1(15,4)
COMMON /BL1/GASUC,GC,AJ,SLOPE/BL2/ITYPE
COMMON /BL8/CPOC(15,4)/BL5/YMW(15)
COMMON /BL11/TEMC(15),PREC(15)
COMMON /BL12/PTEMC,ZTEMC,PPREC,ZPREC/BL18/M
COMMON /BL13/DIAM/BL14/SG/BL9/YMOLF(15)
COMMON /BL15/Q,TSD,PSD,TDD,DEPTH,TSOD,DL
COMMON /BL30/VMASS/BL31/MM
COMMON /BL50/MMM
DATA YMW1/0.44013D 02,0.28013D 02,0.34076D 02,
1 0.16043D 02,0.39070D 02,0.44097D 02,
2 0.58124D 02,0.58124D 02,0.72151D 02,
3 0.72151D 02,0.86178D 02,0.86178D 02,
4 0.10021D 03,0.10021D 03,0.11423D 03/
DATA TEMC1/0.54757D 03,0.22727D 03,0.67237D 03,
1 0.34304D 03,0.54976D 03,0.66568D 03,
2 0.73465D 03,0.76532D 03,0.82877D 03,
3 0.84537D 03,0.89550D 03,0.91337D 03,
4 0.95467D 03,0.97247D 03,0.10239D 04/
DATA PREC1/0.1071D 04,0.4930D 03,0.1306D 04,
1 0.6678D 03,0.7078D 03,0.6163D 03,
2 0.5291D 03,0.5507D 03,0.4904D 03,
3 0.4886D 03,0.4366D 03,0.4369D 03,
4 0.3965D 03,0.3968D 03,0.3606D 03/
DATA CPOC1/0.5316D 01, 0.6903D 01, 0.7070D 01,
A 0.4750D 01, 0.1648D 01, -0.9660D 00,
B -0.1890D 01, 0.9450D 00, -0.2273D 01,
C 0.1618D 01, -0.9081D 00, 0.1657D 01,
D -0.26069D 01, 0.31899D 01, 0.35844D 01,
E 0.791611D-C2, -0.2095000D-03, 0.1737777D-02,

```

```

F      0.6666666D-02, 0.2291111D-01, 0.4043888D-01,
G      0.5519999D-01, 0.4929444D-01, 0.6907777D-01,
H      0.6027777D-01, 0.7843494D-01, 0.7327777D-01,
I      0.9614444D-01, 0.8315605D-01, 0.9502444D-01,
J      -0.2580864D-05, 0.5956790D-06, 0.4209876D-06,
K      0.9351851D-06, -0.4722222D-05, -0.1158950D-04,
L      -0.1695987D-04, -0.1351851D-04, -0.2190432D-04,
M      -0.1655864D-04, -0.2396950D-04, -0.2112345D-04,
O      -0.3089506D-04, -0.2373148D-04, -0.2742592D-04,
P      0.3058984D-09, -0.1174403D-09, -0.1348936D-09,
Q      -0.4509602D-09, 0.2981539D-09, 0.1299725D-08,
R      0.2043895D-08, 0.1433470D-08, 0.2719478D-08,
S      0.1731824D-08, 0.2848422D-08, 0.2362825D-08,
T      0.3834362D-08, 0.2614881D-08, 0.3060012D-08/

```

```

* GC=0.32174D 02
  GASUC=0.1545D 04
  AJ=0.77816D 03
  SLOPE=0.1D 01
  DO 5 I=1, 15
    YMW(I)=YMW1(I)
    TEMC(I)=TEMC1(I)
    PREC(I)=PREC1(I)
  DO 5 J=1, 4
5 CPOC(I,J)=CPOC1(I,J)
  WRITE(6,208)
  WRITE(6,202) (YMW(I),TEMC(I),PREC(I),I=1,15)
  WRITE(6,213)
  WRITE(6,205) ((CPOC(I,J),J=1,4),I=1,15)
  WRITE(6,214)

```

```

C READ 3 CARDS. GAS COMPOSITION. EACH CARD HAS
C THE MOLE FRACTIONS OF 5 COMPONENTS. THE
C ORDER OF COMPONENTS IS AS FOLLOWS :
C CO2,N2,H2S,C1,C2,C3,I-C4,N-C4,I-C5,N-C5,
C I-C6,N-C6,I-C7,N-C7,N-C8 .
C IF GAS COMPOSITION IS NOT AVAILABLE,
C SUPPLY 3 BLANK CARDS .

```

```

READ(5,206) (YMOLF(I),I=1,15)
WRITE(6,215)
WRITE(6,207) (YMOLF(I),I=1,15)

```

```

C READ 1 CARD ( COMBINATION OF FLAG INTEGERS
C AND SG,PTEMC & PPREC )

```

```

READ(5,209) ITYPE,M,MM,MMM,SG,PTEMC,PPREC

```

```

C READ 2 CARDS ( WELL PRODUCTION DATA )

```

```

READ(5,206) DIAM,DEPTH,TSOD
READ(5,206) Q,TSO,PSD,TDD
IF (ITYPE) 1,1,2

```

```

      C F10.3, ' FEET')
217 FORMAT(10X, 'FLOW RATE', 30X, '=', F10.3,
      A ' MSCF/DAY', //,
      B10X, 'THE AVE. TEMP. OF THE GROUND SURFACE  =',
      C F10.3, ' F', //,
      D10X, 'THE TEMP. OF FLOWING STREAM AT SURFACE =',
      E F10.3, ' F', //,
      F10X, 'THE TEMP. OF THE PRODUCING FORMATION  =',
      G F10.3, ' F', //,
      H10X, 'THE WELL HEAD TUBING PRESSURE          =',
      I F10.3, ' PSIA', //)
218 FORMAT(10X, 'INJECTION PROBLEM')
219 FORMAT(10X, 'PRODUCTION PROBLEM')
225 FORMAT(10X, 'THE CALCULATED BOTTOM-HOLE PRESSURE',
      1 ' IS THAT WHEN THE ABSOLUTE DIFFERENCE BETWEEN',
      2 //, 10X, 'EARTH AND CALCULATED FLUID TEMPERATURES',
      3 ' AT THE BOTTOM OF WELL IS WITHIN 0.1 DEG. F .')
226 FORMAT(//, 10X, 'FIRST COLUMN IS DEPTH IN FEET .', //,
      1 10X, 'SECOND COL. IS EARTH TEMP. IN DEG. F .', //,
      2 10X, 'THIRD COL. IS CALC. FLUID TEMP. IN DEG. F .',
      3 //, 10X, 'FOURTH COL. IS CALC. PRESS. IN PSIA .',
      4 //, 10X, 'FIFTH COL. IS R** .', //)
227 FORMAT(10X, 'GAS GRAVITY, PSEUDO REDUCED TEMP.',
      1 ' AND PSEUDO REDUCED PRESS. ARE :', //, 10X, 3F10.4)
STOP
END

```

FUNCTION F(X)

C**** THIS FUNCTION SUBROUTINE IS TO BE
C**** USED WITH MAIN MODEL BB ONLY

C F IS THE FUNCTION TO BE MINIMIZED
C E IS ABSOLUTE ROUGHNESS OF TUBING IN INCHES
C DIAM IS TUBING DIAMETER IN INCHES
C DL IS MONITORED INTEGRATION STEP SIZE IN FEET
C TGEOG IS GEOTHERMAL GRADIENT , DEG. F/ FOOT
C PRE IS PRESSURE IN PSIA
C TEM IS TEMPERATURE IN DEG. RANKIN
C DENS IS DENSITY OF NAT. GAS IN GM/C.C.
C REN IS REYNOLDS NUMBER
C OHTCO IS THE OVERALL HEAT-TRANSFER COEFFICIENT

IMPLICIT REAL*8 (A-H,O-Z)
DIMENSION PRY(5),TRY(5),PRK(5),TRK(5)
COMMON /BL3/AVEMW,BL13/DIAM,BL14/SG,BL18/M
COMMON /BL12/PTEMC,ZTEMC,PPREC,ZPREC
COMMON /BL15/Q,TSD,PSD,TDD,DEPTH,TSOD,DL
COMMON /BL30/VMASS
COMMON /BL50/MMM
D1=0.1D-01
D2=0.1D-01
DL=0.1D 02
E=0.6D-03
RSTAR=X
RR=E/DIAM
TGEOG=(TDD-TSOD)/DEPTH

C THE RUNGE-KUTTA-MERSON INTEGRATION ALGORITHM

TEM=TSOD+0.45967D 03
PRE=PSD
TS=TSOD
TRN=TEM/PTEMC
PRN=PRE/PPREC
TL=0.0D 00
WRITE(6,211) TL,TS,TSOD,PRE,RSTAR
1 DO 14 J1=1,5
IF(J1-1) 2,2,3
2 PRY(1)=PRN
TRY(1)=TRN
TS=TSOD+TGEOG*TL
3 IF(J1-2) 11,4,5
4 PRY(2)=PRN+PRK(1)
TRY(2)=TRN+TRK(1)
TS=TSOD+TGEOG*(TL+DL/0.3D 01)
5 IF(J1-3) 11,6,7
6 PRY(3)=PRN+0.5D 00*PRK(1)+0.5D 00*PRK(2)
TRY(3)=TRN+0.5D 00*TRK(1)+0.5D 00*TRK(2)

$TS = TSOD + TGEOG * (TL + DL / 0.3D \ 01)$
 7 IF (J1-4) 11, 8, 9
 8 $PRY(4) = PRN + 0.375D \ 00 * PRK(1) + 0.1125D \ 01 * PRK(3)$
 $TRY(4) = TRN + 0.375D \ 00 * TRK(1) + 0.1125D \ 01 * TRK(3)$
 $TS = TSOD + TGEOG * (TL + DL / 0.2D \ 01)$
 9 IF (J1-5) 11, 10, 14
 10 $ABC1 = PRN + C.15D \ 01 * PRK(1)$
 $ABC2 = -0.45D \ 01 * PRK(3) + 0.6D \ 01 * PRK(4)$
 $PRY(5) = ABC1 + ABC2$
 $ABC1 = TRN + C.15D \ 01 * TRK(1)$
 $ABC2 = -0.45D \ 01 * TRK(3) + 0.6D \ 01 * TRK(4)$
 $TRY(5) = ABC1 + ABC2$
 $TS = TSOD + TGEOG * (TL + DL)$
 $PRJAM = PRY(5)$
 $TRJAM = TRY(5)$
 11 $PR = PRY(J1)$
 $TR = TRY(J1)$
 $TRS = (TS + 0.45967D \ 03) / PTEM C$
 $PRE = PR * PPREC$
 $TEM = TR * PTEM C$

C CALCULATING IDEAL HEAT CAPACITY OF NAT. GAS

CALL KASEM (TR, TEM, SG, M, CPOM)
 $TRC = TEM / ZTEM C$
 $PRC = PRE / ZPREC$

C
 C
 C
 C CALCULATING COMPRESSIBILITY FACTOR, ITS PARTIAL
 DERIVATIVES W.R.T. PR & TR, REDUCED
 DENSITY, AND HEAT CAPACITY DEPARTURE
 OF NAT. GAS

CALL BWR (TRC, PRC, ITER, Z, DR, CPDM, DZDTR, DZDPR)
 IF (ITER) 22, 22, 12
 12 $DENS = 0.14925826D - 02 * AVE MW * PRE / (Z * TEM)$
 $CPM = CPDM + CPOM$

C CALCULATING VISCOSITY OF NAT. GAS IN C.P.

CALL LEE (TEM, AVE MW, DENS, VISM)
 $REN = 0.20014855D \ 02 * Q * SG / (VISM * DIAM)$

C CALCULATING THE HYDRAULIC FRICTION FACTOR

CALL MOODY (RR, REN, FF, I)
 $U1 = RSTAR$
 $OHTCO = 0.1D \ 01 / U1$

C
 C SOLVING FOR THE RATE OF CHANGE OF REDUCED
 PRESSURE AND TEMPERATURE WITH LENGTH
 CALL QUON (TEM, PRE, TR, PR, TRS, Z, DZDPR, DZDTR,
 1FF, CPM, OHTCO, VMASS, DPRDL, DTRDL, NASER)

```

      IF (NASER-1) 22,13,13
13  PRK(J1)=DL*DPDRDL/0.3D 01
      TRK(J1)=DL*DTRDL/0.3D 01
14  CONTINUE

```

C MONITORING THE INTEGRATION STEP SIZE

```

      ABC5=PRK(1)-0.45D 01*PRK(3)
      ABC6=0.4D 01*PRK(4)-0.5D 00*PRK(5)
      EPR=ABC5+ABC6
      ABC5=TRK(1)-0.45D 01*TRK(3)
      ABC6=0.4D 01*TRK(4)-0.5D 00*TRK(5)
      ETR=ABC5+ABC6
      ETEM=DABS(ETR*PTMC)
      EPRE=DABS(EPR*PPREC)
      IF (EPRE-D1) 15,15,16
15  IF (ETEM-D2) 17,17,16
16  DL=DL/0.2D 01
      GO TO 1
17  PR=PRJAM-EPR
      TR=TRJAM-ETR
      TL=TL+DL
      IF (EPRE-D1/0.16D 02) 18,20,20
18  IF (ETEM-D2/0.16D 02) 19,20,20
19  DL=0.2D 01*DL
20  PRN=PR
      TRN=TR
      PRE=PRN*PPREC
      TEM=TRN*PTMC
      TED=TEM-0.45967D 03
      WRITE(6,211) TL,TS,TED,PRE,RSTAR
      DEL=DEPTH-TL
      IF (DABS(DEL)-0.1D-04) 23,23,25
25  IF (DEL-DL) 21,1,1
21  DL=DEL
      GO TO 1
22  WRITE(6,227)
23  TED=TEM-C.45967D 03
      TEMTL=TED
      F=TED-TEMTL
211 FORMAT(10X,5F10.2)
227 FORMAT(//,10X,'*** THE COURSE OF ITERATIONS IS'
*, ' INTERRUPTED ***',//,7X,'*** THE PROGRAM IS'
*, ' GOING TO ASSUME ANOTHER "RSTAR"***')
      RETURN
      END

```

SUBROUTINE AMEIN

C**** THIS SUBROUTINE CAN BE USED WITH
C**** MAIN MODEL BB ONLY

C THIS SUBROUTINE CALCULATES THE MOLAL AVERAGE
C PROPERTIES AND THE WICHERT & AZIZ CORRECTED
C PSEUDOCRITICAL PROPERTIES OF NAT. GAS

C AVMW IS AVERAGE MOLECULAR WEIGHT
C SG IS GAS GRAVITY
C PPREC IS PSEUDOCRITICAL PRESSURE IN PSIA
C PTEMC IS PSEUDOCRITICAL TEMPERATURE, DEG. R
C ZTEMC IS CORRECTED PSEUDOCRITICAL TEMPERATURE
C ZPREC IS CORRECTED PSEUDOCRITICAL PRESSURE

IMPLICIT REAL*8 (A-H,O-Z)
COMMON /BL3/AVMW/BL5/YMW(15)/BL9/YMOLF(15)
COMMON /BL11/TEMC(15),PREC(15)
COMMON /BL12/PTEMC,ZTEMC,PPREC,ZPREC
COMMON /BL14/SG/BL18/M/BL31/MM
COMMON /BL50/MMM
Y11=YMOLF(1)
Y22=YMOLF(2)
Y33=YMOLF(3)
IF(MMM-1) 10,3,3
10 IF(M-1) 3,1,1
1 AVMW=0.0D 00
PTEMC=0.0D 00
PPREC=0.0D 00
DO 2 I=1,15
AVMW=AVMW+YMOLF(I)*YMW(I)
PTEMC=PTEMC+YMOLF(I)*TEMC(I)
2 PPREC=PPREC+YMOLF(I)*PREC(I)
SG=AVMW/0.2897D 02
GO TO 7
3 AVMW=SG*0.2897D 02
IF(MM-1) 11,7,7

C THE BROWN ET AL AND CARR ET AL CORRELATIONS
C FOR THE PSEUDOCRITICAL PROPERTIES OF
C SWEET AND SOUR NAT. GASES

11 PTEMC=0.171137D 03+0.313725D 03*SG
TCCOR=-0.8D 02*Y11-0.25D 03*Y22+0.13D 03*Y33
PTEMC=PTEMC+TCCOR
IF(SG-0.85D 00) 4,4,5
4 PPREC=0.6951D 03-0.4D 02*SG
GO TO 6
5 PPREC=0.704396D 03-0.51724D 02*SG
6 PCCOR=0.44D 03*Y11-0.17D 03*Y22+0.6D 03*Y33
PPREC=PPREC+PCCOR

C THE WICHELT & AZIZ CORRECTED PSEUDOCRITICALS
C DUE TO PRESENCE OF CO2 AND H2S

```
7 S1=YMOLF(1)+YMOLF(3)
  S2=YMOLF(3)
  E31=0.12D 03*(S1**0.9D 00-S1**0.16D 01)
  E32=0.15D 02*(DSQRT(S2)-S2**4)
  E3=E31+E32
  S4=PTMC+S2*(0.1D 01-S2)*E3
  ZTMC=PTMC-E3
  ZPREC=PPREC*ZTMC/S4
  RETURN
  END
```