

487 ± 98 μm particle size: before compression

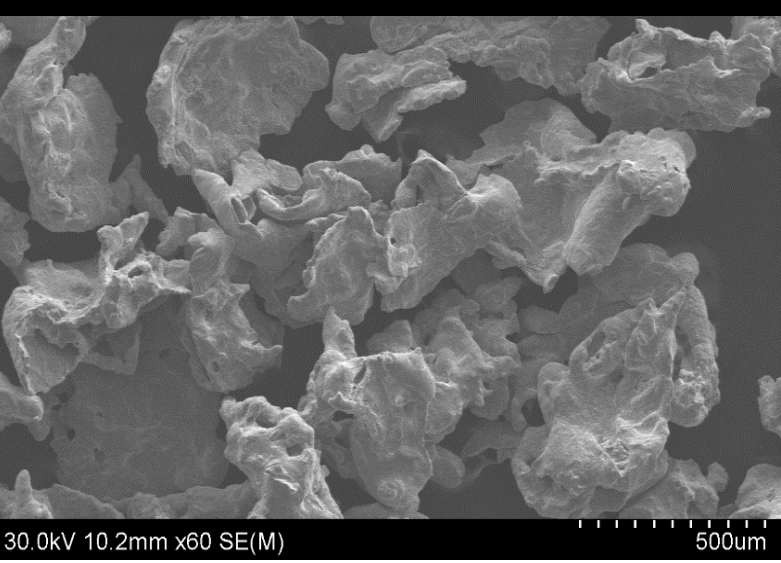
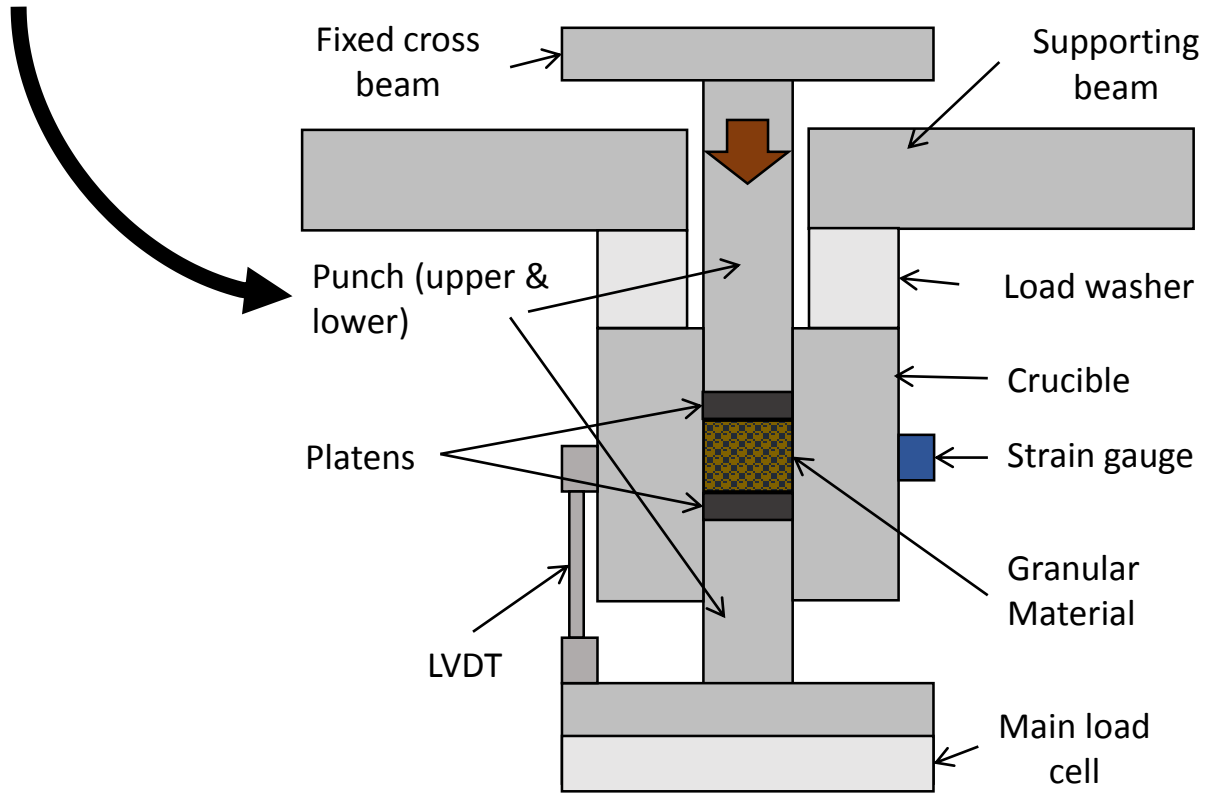
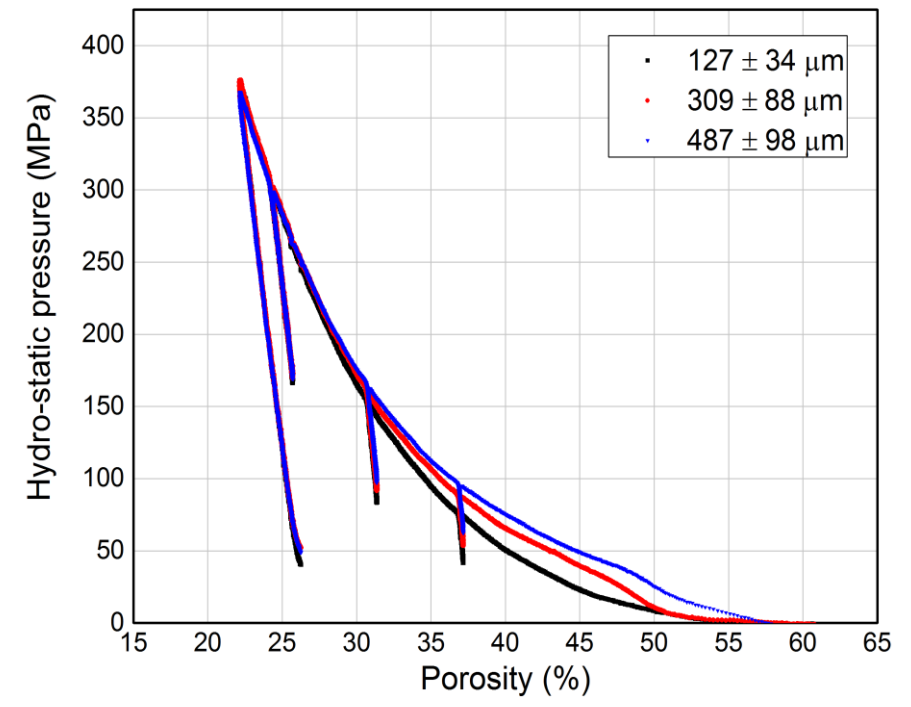


Figure 1 - In the figure above SEM images were taken of the SS316 powder at an average size of 487 μm before the material was compresses axially.



487 ± 98 μm particle size: after compression

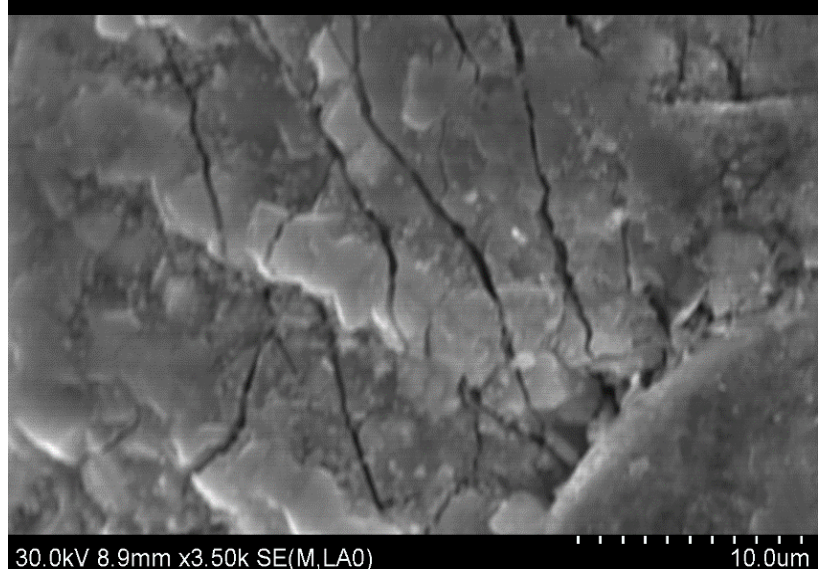


Figure 2 - Post-mortem SEM images were taken to visualize the mechanism responsible for material failure under compaction. Micro fractures were observed to be dominant.

*Highlights (for review)

- Particle size dependents exists; for the porosity vs hydro-static pressure relation
- Bulk modulus: sensitive to the particle size and is dictated by different failure
- 127 ± 34 μm particles exhibited higher stresses when compared to 487 ± 98 μm particles

Confined Uniaxial Compression of Granular Stainless Steel 316

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Abstract

The quasi-static compaction response of granular stainless steel 316 was studied under triaxial loading conditions using a confined crucible experiment apparatus, specifically looking at effects of the particle size ($127 \pm 34 \mu\text{m}$, $309 \pm 88 \mu\text{m}$, $487 \pm 98 \mu\text{m}$) on mechanical behavior. The material response was captured using load washer and strain gauges to relate: the porosity effects as a function of hydro-static pressure, particle size dependency on wall friction effects, and particle size-dependent failure mechanism. Our observations revealed that the path of crushing out porosity varied based on the particle size and the frictional effects. Scanning Electron Microscope (SEM) images were taken to observe the surface features of the compacted material and comment on failure mechanisms. Using these techniques it was observed that the smaller particles exhibited significant plastic deformation and flow, while the larger appeared to show micro cracking which lead to inelastic deformation and particle fracture. Altogether, these results are important because granular behavior is critical in powder flow applications such as additive manufacturing.

Keywords: Stainless steel, Granular, Triaxial compaction, Static, Failure of granular materials

1. Introduction

Granular materials are an agglomeration of discrete solid macroscopic particles that behave differently than continuous solids, liquids and gases. The understanding of granular behavior is critical in manufacturing [1], pharmaceutical [2], and geo-technical [3] applications. Researchers choose to pursue two main approaches when describing the granular behavior: discrete analysis and

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6 continuum mechanics. In the discrete approach, research is focused on individual particle behavior
7 [4, 5, 6]. The discrete models are based on kinematics of the particles and allow for simulating
8 mechanical characteristics [7]. More recently, discrete models exhibit deformability and allow for
9 the simulation of the yielding behavior of materials [8]. Conversely, in continuum-scale studies
10 researchers focus on using constitutive laws and conservation equations to model the ensemble of
11 granular behavior [9, 10]. Specifically, researchers have derived a list of yield and friction models
12 to express various characteristics such as variation in friction as a function of localized stress [7].
13 However, challenges exist with both approaches. Discrete modeling is computationally expensive
14 and struggles with complex non-linear behavior [7, 11]. On the other hand, continuum models
15 struggle with dry grains, for example, that do not support tension. This means the behavior changes
16 from plastic media to a gas-like disconnected state which is difficult to represent in a unified model
17 [9]. Researchers have been interested in the micro-scale of solid-like behavior of grains [12] while
18 capturing the macro-scale fluid-like characteristics [13] due to the large application possibilities.
19 Challenges exist when working with granular materials and so, to address this issue, the current
20 study of granular media has been undertaken by looking at granular failure in cold press compaction.

21 The compaction of cold-pressed powder is of particular interest to pharmaceutical researchers
22 [14], when looking at particle distribution, Jansen coefficient for particle friction, and porosity
23 as a function of applied stress [2, 15]. For example, Michrafy et al. [15] studied the frictional
24 effects of pill compaction lubricant effects on the Columb wall friction coefficient. They showed
25 that wall friction and particle friction primarily account for non-uniform stress distribution and
26 produce a density gradient of the material in the compact; the lack of consistent density results in
27 skewed mechanical properties [16]. Similar work was conducted by Kadiri et al. [14]. They looked at
28 predicting the axial density distribution of microcrystalline material while determining the material
29 properties during quasi-static triaxial compression. The axial density of cellulose decreases from
30 the top to the bottom of compaction and Kadiri et al. [14], concluded that the particle size and
31 shape significantly influence the density distribution throughout pharmaceutical tablet compaction.

32 Similar work on cold compaction of granular materials has been conducted in additive manu-
33 facturing (AM), where it is important to understand the relationship between the powder charac-

34 teristics and the mechanics of the consolidated part [17, 18]. For example, mechanical properties,
35 surface finish, and integrity of the final structure are highly influenced by the characteristics of
36 the powder material that is used in the manufacturing process. Further, flowability of the powder
37 highly influences the finish, grade, and strength of the final product. The correlation between fac-
38 tors such as the shape distribution, size distribution, density, and packing density are a complex
39 and an ongoing research topic [19]. In one example, Spierings et al. [20] determined that larger
40 stainless steel (SS316) particles in the raw powder consequently resulted in bigger pores in the
41 final steel part. Thicker layers created inhomogeneous regions in the structure that promoted frac-
42 ture [20] and increased void density creates nucleation sites for brittle failure, and promotes crack
43 propagation. Comparable research was done by Bai et al. [21] when analyzing the binder jetting
44 AM technique with copper particles. Similar to laser sintering, binder jetting combines the metal
45 powder with a binder and creates a semi consolidated green structure surrounded by the remaining
46 powder, which is then sintered. By using a bimodal powder distribution, Bai et al. [21] was able to
47 improve the powder density by 8.2% and flowability by 10.5%. Powder flow properties (flowability)
48 has been shown to play an important role in the cohesive strength, friction, compressibility, and
49 transportation of the powder [19, 22]. The tailoring of particle sizes is required to improve the final
50 density of the structure.

51 In this study, we are primarily concerned with the behavior of granular stainless steel particles.
52 Due to its high hardness and elasticity, stainless steel 316 (SS316) powder has been used in various
53 industry applications. These applications include making artificial joints in medical research [23],
54 creating tooth implants in the dentistry industry [24], and as a AM material in creating complex
55 geometries in the manufacturing industry [17]. To optimize the usage of SS316 powder in the various
56 industries, the shortcomings of powder flowability and particle variability relating to strength must
57 be addressed and understood [17]. Further research is required to obtain a better understanding
58 between particle characteristics, strength, and product performance. Previous research shows a
59 gap in understanding the behavior of particles in granular material. Addressing this gap would
60 improve the repeatability of manufacturing and the final performance of the products by limiting
61 the large variation we see in granular SS316. The behavior of granular SS316 is studied in this

62 paper with an emphasis on the effect of particle sizes on flow behavior and particle variability in
63 strength. A uniaxial confined experimental technique is used to evaluate the triaxial flow stress
64 SS316 exhibits for increasing particle sizes. This allows for progress in evaluating accurate failure
65 mechanisms in the material and linking the overall effects of particle size on material behavior.
66 In this paper we investigate the effects of hydrostatic pressure and porosity of the SS316 powder.
67 Further, we look at the relationship between the axial and radial stress to comment on the Janssen
68 coefficient as a function of changing particle size. Lastly, we are interested in the loading and
69 unloading behavior of the SS316 powder observing the changing stiffness degradation and material
70 behavior. By addressing these areas, advancements will be obtained to better understand the effects
71 of particle size on strength and part quality where stainless steel particles are used.

72 *1.1. Experimental Techniques*

73 *1.2. Material*

74 *1.2.1. Composition and Geometry*

75 SS316 powder was used in the experiment and the elemental composition was provided by the
76 manufacturer and is summarized in (Table 1). The elemental breakdown is typical for SS316. The
77 powder was sourced from Alfa Aesar by Thermo Fisher Scientific of Tewksbury, Massachusetts. The
78 average particle sizes used for the triaxial compaction were three different size ranges: 1. 127 ± 34
79 μm , 2. $309 \pm 88 \mu\text{m}$, and 3. $487 \pm 98 \mu\text{m}$.

80 The SS316 powders used, range from $127 \mu\text{m}$ to $487 \mu\text{m}$ in size and exhibit a rough irregular
81 shape. This is most likely due to the manufacturing process of using water atomization to create
82 the powder [25]. For visual aid, the powder geometry is illustrated in Figure 1 for sieved particle
83 sizes: $309 \pm 88 \mu\text{m}$ and $487 \pm 98 \mu\text{m}$. The images were taken with a Hitachi S-4800 Field Emission
84 Scanning Electron Microscope (SEM).

85 *1.2.2. Particle Distribution*

86 Three different sizes were sieved, microscopically analyzed, and tested to observe the effects
87 of mean particle diameter, $\bar{\phi}$ (μm), on hydro-static pressure as a function of porosity. The mean
88 and standard deviation was documented for each particle size for the range of: $127 \pm 34 \mu\text{m}$,

89 $309 \pm 88 \mu\text{m}$, and $487 \pm 98 \mu\text{m}$, summarized in Table 2. The particle sizes were chosen based on the
90 restrictions of the crucible design and access to materials with a higher Rockwell hardness. Every
91 specimen tested was separately analyzed and sampled following principal sampling techniques for
92 granular material outlined by Maynard [26].

93 The particle size distribution and shape was analyzed using the automated Malvern Morphologi
94 G3 microscope (G3). The De Broukere mean diameter (volume moment mean diameter) was
95 considered when analyzing the particle sizes. This parameter is most relevant when determining
96 the distribution of the bulk sample used in experiment and has been used by many in literature
97 [27, 28]. The definition of this relation is on the basis of introducing another linear term in diameter,
98 analogous to moments of inertia i.e. accounts for the center of gravity of the particle distribution.
99 The next point to underline is that the advantage of this method does not require the particle count.
100 This was done by taking the square difference of the upper and lower bin to determine D_i which
101 determined the mean particle distribution in the sample. The numerator is taken to the power of
102 four (the power of three to account for the equivalent volume and the additional power to account
103 for the center of gravity) and the denominator is taken to the power of three (where it accounts for
104 the mass of the particles). Seen below is the summarized equation:

$$D[4, 3] = \frac{\sum_1^n D_i^4}{\sum_1^n D_i^3} \quad (1)$$

105 To evaluate the variability of the powder, the span, Δ , was calculated for each sample, as was
106 similarly used by Engeli et al. [29]. The span was determined through the following and this takes
107 into account the tenth (D_{10}), fiftieth (D_{50}), and ninetieth percentile (D_{90}):

$$\Delta = \frac{D_{90} - D_{10}}{D_{50}} \quad (2)$$

108 The particle size based on (1) and the span of the distribution based on (2) are summarized
109 in Table 2. With an increase in mean particle diameter, the standard deviation increases while
110 the span decreases, highlighted in Table 2. These particle statistics are considered later when
111 mechanical behavior is investigated.

112 Particle shapes were also considered. Shown in Table 2, the average circularity, C , of the
113 particles decreases as the average particle size increase. Circularity refers to how close the particles
114 resemble a perfect sphere. Less circular particles have a larger deviation in the particle size. The
115 circularity of the material is directly related to the manufacturing process such that the method of
116 atomization dictates the final shape of the particle [25]. Here we computed particle circularity(C)
117 defined as:

$$C = 2\sqrt{\frac{\pi A}{P}} \quad (3)$$

118 where A , (m^2), and P , (m), are the area and perimeter of the particle, respectively. As the particles
119 increase in mean diameter the circularity decreases. The shape results are important because
120 as noted by Schade et al. [19] the decrease in circularity causes the particles to align in other
121 orientations and ineffectively fill in the voids. Schade et al. [19] determined that the difference
122 in circularity is related to the atomization process of the granular material i.e. gas atomization
123 produces spherical shaped particles and water atomization produces rough, irregular particles.

124 *1.3. Quasi-static compression*

125 *1.3.1. Configuration*

126 Illustrated in Figure 2 is a schematic of the experimental apparatus used to conduct confined
127 uniaxial compression experiments on the granular material to obtain a triaxial response. The top
128 punch, bottom punch, and platens are fabricated from D2 tool steel and were heat treated using
129 quenching and double tempering to reach a hardness of 62 HRC. The punches are used to press
130 together the platens that contact the granular material. The top and bottom punches are designed
131 around the accessibility of the MTS 810 machine used in this study, and are 130 mm in length.
132 The crucible was made out of 4340 steel and heat treated by quench and tempering followed by
133 gas nitriding. The maximum hardness achieved was 52 HRC with an inner and outer diameter of
134 6.3 mm and 22.2 mm, respectively. The inner diameter was machined with a tight tolerance so that
135 the sacrificial platens were able to seal in the granular material and to protect the punches from
136 the granular media. The supporting beam, designed out of aluminum, was placed to differentiate

137 between the normal forces on the top and bottom of the sample. The support beam is held up with
138 aluminum blocks that attach to the MTS. Aluminum supports were designed for attaching the linear
139 variable differential transformer (LVDT) so that relative displacement could be recorded to minimize
140 the deflection during compression. The operating parameters for the LVDT were ± 7.5 mm and
141 0.2% linearity error. The MTS machine outputs the applied load at a rate of 0.33 kN s^{-1} with a
142 resolution of $305 \mu\text{V}$ per analogue-to-digital converter count, placing the specimen in a quasi-static
143 stress state. Between the top punch and the crucible is a load washer from Omega Engineering Inc.
144 The operating parameters for the load washer was 10 kN with an accuracy of $\pm 0.5\%$.

145 For specimen preparation, one end of the crucible was closed using the platen and punch. In
146 order to limit the wall frictional effects occurring during confined compaction it is critical to maintain
147 an aspect ratio of height to width of < 1 when determining the sample size [30]. Granular material
148 was poured in and an aspect ratio of 0.76 ± 0.05 was achieved. The volume of material needed to
149 satisfy this constraint is $150.8 \pm 3.3 \text{ mm}^3$. Sample volume was not the same for all samples and the
150 uncertainty reflects the variation.

151 The supporting beam illustrated in Figure 2 creates a free floating state such that the compaction
152 of the material is not affected by the weight of the crucible. The top punch was lowered, rested
153 against the platen on the top of the specimen with the assumption that the specimen has not been
154 compacted. A linear variable differential transformer (LVDT) was secured to the beam, so that
155 the compaction depth could be measured relative to the displacement. The load washer was used
156 to measure the combined frictional effects during the uniaxial compression. Next, we outline the
157 theory used to interpret the measurements.

158 *1.3.2. Theory*

159 To better understand the response of the granular material during quasi-static triaxial compres-
160 sion, we investigate the relationship between the hydrostatic pressure and porosity by tracking the
161 volumetric strain and relative density in the confined crucible. First, we track the time-evolving re-
162 duction in porosity of our test samples through measurements of initial mass m (kg), initial packing
163 porosity ϕ_i (%), and the cross-head displacement δ (mm) of a plunger that is used to compress our

164 granular samples. The mass of the initial granular sample is measured by a digital scale with the
 165 precision of 0.01 g. The size of each samples was controlled by volume ($\sim 150.8 \text{ mm}^3$). Throughout
 166 the compression experiment, the change in volume, ΔV (m^3), is related to axial displacement:

$$\Delta V = A_0 \delta \quad (4)$$

167 where A_0 is the cross-section area of the void (m^2) and δ is the relative axial displacement dur-
 168 ing compression (m). From there, we can calculate the specimen density ρ as it evolves during
 169 compaction:

$$\rho = \frac{m}{V_0 - \Delta V} \quad (5)$$

170 where m is the mass of the specimen (kg), and V_0 is the initial specimen volume (m^3). The evolving
 171 porosity is calculated by normalizing the specimen density with the solid bulk density:

$$\phi = 1 - \frac{\rho}{\rho_s} \quad (6)$$

172 where ϕ is the porosity fraction (unit less) and ρ_s is the bulk solid density (kg m^{-3}). For stainless
 173 steel the bulk density is taken as 8000 kg m^{-3} , which provided by the manufacturer.

174 Next, the hydrostatic pressure is calculated by measuring the axial stress, σ_{zz} (MPa), from
 175 the MTS machine and the radial stress, $\sigma_{\theta\theta}$ (MPa), measured from mounted strain gauges on the
 176 crucible, see schematic in Figure 2. The equation for hydro-static stress, σ_{hyd} (MPa), in cylindrical
 177 coordinates is defined as [31]:

$$\sigma_{hyd} = \frac{1}{3} (2\sigma_{\theta\theta} + \sigma_{zz}) \quad (7)$$

178 We calculate σ_{zz} by dividing the axial force experienced by the sample, F_{zz} (N), and the cross-
 179 sectional area of the void A_0 (m^2), assuming that the area does not significantly deform during
 180 compression:

$$\sigma_{zz} = \frac{F_{zz}}{A_0} \quad (8)$$

181 F_{zz} was computed by subtracting the axial force outputted by the MTS machine, F_{MTS} , and
 182 the friction forces, F_f (N), induced by the platens and granular powder contacting the crucible
 183 walls:

$$F_{zz} = F_{MTS} - F_f \quad (9)$$

184 To determine the radial stress, we assumed the crucible was a thick walled cylinder (TWC).
 185 The TWC equation assumes that the crucible geometry is symmetric on θ ($^\circ$) and the stress is only
 186 a function of r (m). The problem is statically determinate and so only the equilibrium equations
 187 must be satisfied. The derivation of the equilibrium equations along with the solution have been
 188 computed extensively in literature and therefore will not be explicitly shown (see reference [31] and
 189 [32] for full derivation). Equation (7) can be rewritten to include the axial stress and radial stress
 190 to obtain an expression for hydro-static pressure as a function of measurable parameters:

$$\sigma_{hyd} = \frac{1}{3} \left(\frac{F_{zz}}{A_o} + E_{cr} \varepsilon_{\theta\theta} \frac{b^2 - a^2}{a^2 (1 - \nu_{cr}^2)} \right) \quad (10)$$

191 where E_{cr} (MPa) is the stiffness of the crucible, ν_{cr} is Poisson's ratio of the crucible, a is the
 192 inner radius (m), b is the outer radius (m), and $\varepsilon_{\theta\theta}$ is the radial strain as a function of thickness
 193 [31]. There are, however, limitations of using the TWC approach which has been identified by, for
 194 example, Kim et al. [33]. According to their research, shortcomings exists when measuring hoop
 195 strain. This is dependent on the inner diameter, cylinder thickness, and location of strain gauge
 196 along the axial direction. In our apparatus, the inner diameter is much smaller than the diameter
 197 tested by Kim et al. [33], further, when extrapolating from the two sizes that were compared, our
 198 void size is proportional to the maximum strain we measured. In addition, the strain gauge size
 199 and placement encompassed in our experiment the entire testing section and so we think that some
 200 of the uncertainty in the measurements is mitigated.

201 Next, we look at the bulk behavior of the material. This characteristic describes the compress-
202 ibility of the material and relates the change in pressure of the material with respect to volume.
203 This is given by:

$$B_{ep} = \frac{\sigma_{hyd}}{\frac{\Delta V}{V_o}} \quad (11)$$

204 where B_{ep} represents the bulk modulus taking into account elastic and plastic behavior (MPa), and
205 all the other variables have been previously defined. This parameter evolves during loading and is
206 an indicator of deformation in the granular sample. Similar calculations were done by Gustafsson
207 et al. [34] when conducting confined compression experiments with iron ore.

208 We also consider the axial-to-radial effects by calculating the Janssen coefficient of SS316 powder.
209 This parameter relates the radial and axial stresses during compaction which allows for simplifi-
210 cation when conducting computer modeling of the compaction [35]. The Janssen constant, K , is
211 given by:

$$K = \frac{\sigma_{\theta\theta}}{\sigma_{zz}} \quad (12)$$

212 Lastly, we look at frictional effects by monitoring the transmitted stress ratio, T (unitless). To
213 do this, we calculate the ratio of transmitted stress, σ_t (MPa), over applied stress, σ_a (MPa). The
214 applied stress is the stress that the compression machine outputs and the transmitted stress is the
215 stress that is interpreted from the load washer. The difference in applied and transmitted stresses
216 provides insight on how much energy is lost to friction in the crucible apparatus. This ratio is given
217 by:

$$T = \frac{\sigma_t}{\sigma_a} \quad (13)$$

218 To account for the uncertainty in the experiment, we conducted a systematic propagation of
219 error, taking into account the uncertainty of the sensors and measured geometries. This will help
220 in understand the accuracy of our results. Based on the guide outlined by Berendsen [36], Table 4

221 summarizes the relative uncertainty of critical material parameters that were calculated. The rules
222 for calculating uncertainty have been derived and computed extensively in literature and will not
223 be explicitly shown. Refer to [36] for full derivations.

224 **2. Experimental Outcome**

225 *2.1. Porosity*

226 As an outcome of the confined uniaxial compaction, we first investigate compressibility, loading
227 path, average particle diameter and geometry of the material. As a limitation to this work, addi-
228 tional tests should be conducted and a wider range of particle sizes should be included to account
229 for any outlying behavior. Shown in Figure 3 is the relationship between the hydro-static pressure
230 and porosity for three different particle size ranges; $127 \pm 34 \mu\text{m}$, $309 \pm 88 \mu\text{m}$, $487 \pm 98 \mu\text{m}$. The
231 initial porosity for each test was 60 % for size $127 \pm 34 \mu\text{m}$, 66 % for size $309 \pm 88 \mu\text{m}$, and 67 %
232 for size $487 \pm 98 \mu\text{m}$. From the initial porosity to a porosity of 30 %, the curve of each particle size
233 behaves differently for each sample. For porosities above 30 %, smaller particles have lower hydro-
234 static pressure. For porosities less than 30 %, the relationship between porosity and hydro-static
235 pressure collapses onto a single curve for each particle size range studied. Previous work by Heckel
236 [37] noted similar results when studying metallic powders such as iron, nickel, tungsten, and copper.
237 Linear-like trends were observed for compressive stresses applied $>135 \text{ MPa}$. In addition, no dif-
238 ference in load-unload behaviors are observed across all particle sizes for a given porosity. As the
239 hydro-static pressure is reduced, the effect of unloading results in incremental increases in porosity
240 (the linear curves back to the right). The uncertainty for the hydro-static pressure parameter was
241 calculated to be 11.3 % and 10.9 % for the small ($127 \pm 34 \mu\text{m}$) and larger particles ($487 \pm 98 \mu\text{m}$).
242 This deviation was consistent for the particle sizes investigated, which shows that the equipment
243 was consistent through out every test conducted.

244 *2.2. Bulk Modulus*

245 The bulk modulus represents the ability of the material to withstand compaction. Shown in
246 Figure 4 is the relationship between the bulk modulus (described in Equation (11)) as a function

247 of the applied stress. This plotting convention is typically seen in literature [38, 39]. This figure
248 accounts for load behavior carried out during compression. Notable in the figure is that the smaller
249 particle size range reach a higher bulk modulus (760 MPa for $127 \pm 34 \mu\text{m}$) while the larger particles
250 appear to reach a lower value (663 MPa for $309 \pm 88 \mu\text{m}$, and 648 MPa for $487 \pm 98 \mu\text{m}$). Similar
251 studies conducted in the past showed that granular aluminum followed linear trends when loaded
252 with a compressive stress of $>300 \text{ MPa}$ [40]. Interestingly, the aluminum particle trends observed
253 were that larger particles (150 - 212 μm) had a slightly steeper slope than the smaller particles (53
254 - 75 μm) when looking at the B_{ep} as function of pressure. The bulk modulus is also more sensitive
255 to the smaller particles (it increases at a faster rate). ($309 \pm 88 \mu\text{m}$, $487 \pm 98 \mu\text{m}$). Propagation of
256 error was likewise completed looking at the B_{ep} . Based on our calculations, the relative uncertainty
257 was 11.7% and 2.9% for the small particle size ($127 \pm 34 \mu\text{m}$) and large particle size ($487 \pm 98 \mu\text{m}$)
258 respectively.

259 *2.3. Janssen Coefficient*

260 Next, we investigate the Janssen coefficient which is used to relate the axial to the radial stresses
261 which helps in simplifying analytical models when simulating triaxial compression behavior. Shown
262 in Figure 5 is the relationship between the porosity and the Janssen coefficient for three particle
263 size ranges. At higher porosity there is more variability in the Janssen coefficient across and within
264 each particle size range. Specifically, larger particles have a higher Janssen coefficient for larger
265 porosities. As the porosity is crushed out, (i.e reduced), the Janssen coefficient for all particle size
266 ranges, converges to 0.23 (near 30% porosity). The typical range for the Janssen constant seen in
267 bulk materials has been noted to be 0.3-0.6 by [41]. These values are typically seen in round-like
268 particle shapes and so the interesting behavior seen in Figure 5 is dependent on the elongated and
269 rough edged particles. The relative uncertainty for the Janssen coefficient was 1.5% and 1.7% for
270 the small ($127 \pm 34 \mu\text{m}$) and larger ($487 \pm 98 \mu\text{m}$) particles respectively.

271 *2.4. Friction*

272 Next, the wall friction effects can be probed by calculating the transmission ratio, which relates
273 to the transmitted force through the material. Shown in Figure 6 is the transmission ratio as a

274 function of applied load for all particle sizes. For a given applied stress the transmission ratio
275 decreases for large particle sizes. For lower applied stresses, the values for lower applied stress is
276 related to the compliance of the system. As the applied stress increases, the transmission ratio
277 converges to 0.96. The small particles ($127 \pm 34 \mu\text{m}$) approach convergence faster in comparison
278 to larger particles ($487 \pm 98 \mu\text{m}$). To account for the systematic error in the experiment, the
279 uncertainty was calculated. For the small particles ($127 \pm 34 \mu\text{m}$) the relative uncertainty was
280 6.5% and for the large particles ($487 \pm 98 \mu\text{m}$) the relative uncertainty was 6.6%.

281 *2.5. Failure analysis*

282 Lastly, SEM images were taken of the consolidated SS316 to investigate failure features on the
283 compacted specimens surfaces. Attempts were made to cut and polish the consolidated pucks after
284 testing, but that only introduced further damage to the specimens. Shown in Figure 7a is the
285 failure surface of the $309 \pm 88 \mu\text{m}$. This demonstrates that these particles withstood significant
286 plastic deformation. Elongated laminate structures are noted (red arrows) on the surface, which
287 are believed to be generated from particle shear stress during compaction. Shown in Figure 7b
288 is the failure surface of the larger particles ($487 \pm 98 \mu\text{m}$). There is noticeably more fracturing
289 and cracking (red arrows). and the surface appeared to be more jagged in comparison to the
290 smaller particles. Such evidence has been noted before by Roberts and Rowe [42]. Based on
291 theoretical equations and experimental evidence, larger particles crack because the stress required
292 for brittle fracture is less than the stress required for plastic flow. Likewise, smaller particles undergo
293 plastic deformation due to the stress required is lower than the brittle deformation stress. These
294 observations seen in our experiments are linked with the experimental data next in the Discussion.

295 **3. Discussion**

296 In this paper, we explored the mechanical response of granular SS316 for size ranges of: 127 ± 34
297 μm , $309 \pm 88 \mu\text{m}$, and $487 \pm 98 \mu\text{m}$. To accomplish this, we adapted a uniaxial compression exper-
298 iment utilizing equipment and sensors: MTS 810, loader washer, displacement LVDT, and strain

299 gauges. Similar experiments have been performed in the powder metallurgy and defense indus-
300 tries [43, 44, 45, 46, 47, 48], but limited data exists for stainless steel, however, it has been noted
301 by Roberts and Rowe [42] that larger particles tend to fracture while smaller particles tend to
302 plastically deform. Limitations to this work have been observed and noted. The use of (10) is
303 limited based on the assumption of uniform radial stress and negligible friction. Kim et al. [33]
304 underlined the limitations of our approach, using the TWC method. The other limitation of our
305 design was the wall thickness of the crucible. The result of this would be decreased range in strain
306 measurements. To resolve the issue additional calibration tests could be done. Furthermore, our
307 design of the crucible was based on the assumption that the radial stress is uniform which would
308 result in having the specimen friction-free, and therefore (10) is limited to the case where friction
309 is negligible. However, we mitigated this issue by having the length of the strain gauge encompass
310 the entire height of the specimen. Additional thought is given to the design parameters and so
311 further calibration testing will be conducted to narrow the variability of the results and to validate
312 the simplification of (7) from 3D to 2D space as highlighted by Meyer and Faber [46]. To better
313 validate the material behavior, additional tests could be conducted to eliminate the outlying trends.

314 First we investigated the relationship between hydro-static pressure and porosity as a function
315 of particle size. Other researchers have also looked at hydro-static pressure effects on granular
316 material [44, 49]. In our experiments (Figure 3), it was observed that deviations in the hydro-static
317 pressure among the particle sizes were sensitive for porosities greater than 30 %. For porosities less
318 than 30 %, the behavior converges independent of particle size. This behavior has been observed
319 before [50] when studying low carbon SS316 and the result was determined that densification of the
320 material was sensitive to particle size. Cristofolini et al. [50] also demonstrated that initial porosity
321 of the material highly influenced the loading path and that force dissipation during loading mostly
322 originated from wall friction. These behaviors of granular stainless steel 316 have previously not
323 been greatly considered or linked back to global granular response. We believe these linkages
324 are valuable contributions in the additive manufacturing and other powder based industries. When
325 calculating the uncertainty in porosity for the smallest and largest average particle sizes, we obtained
326 values of 0.97 % for the $127 \pm 34 \mu\text{m}$ size and 0.96 % for the $487 \pm 98 \mu\text{m}$.

327 Despite general insensitivities in particle size influence on the relationship between porosity
328 and hydro-static pressure, it was observed that the bulk modulus was sensitive to particle size
329 (Figure 4). In our experiment, we observed smaller particle sizes had higher bulk stiffness, were
330 more sensitive to applied stress, and were more sensitive to unloading. Indeed, similar experiments
331 [51, 52] demonstrated that particle size and compaction pressure influenced the mechanical response
332 of the material. More importantly, the particle distribution with particles of different sizes required
333 less energy to compact. In our experiments the ranging particle sizes dictated the mechanical
334 response differences, which is observed to be associated with failure. Specifically, we note that
335 smaller particles ($127 \pm 34 \mu\text{m}$) undergo more plastic deformation (Figure 7a) as a consequence
336 of compaction, while the larger particle ($487 \pm 98 \mu\text{m}$) exhibits more fracture and micro-cracking
337 (Figure 7b). Note that these particle size-dependent trends have been noted in other fields [53, 54,
338 55], and in particular this brittle-ductile transition behavior has been noted by Roberts and Rowe
339 [42].

340 In addition to bulk response, we looked at the stress transmission ratio to investigate frictional
341 behavior in compaction of granular SS316. In order to maximize the stress transmission through
342 the material and minimize wall effects, the samples must maintain an aspect ratio of < 1 . From our
343 experiments, referring to Figure 6, the transmission ratio for the smaller particles ($127 \pm 34 \mu\text{m}$),
344 the curve begins a steep climb and later plateaus as the applied stress increases. Similar research
345 conducted by Fleck and Cocks [30], showed that increasing the aspect ratio beyond 1 resulted in
346 a significant decrease in the stress transmission. This leads us to better replicating an isotropic
347 compaction environment. Other studies conducted by Perez-Gandarillas et al. [56] underlined the
348 idea that lower axial transmission was observed with an increase in particle size. They concluded
349 that part of the loading energy was consumed by the breakage of larger particles. Tracking these
350 relationships allows for a better understanding of the complex behavior of fragmentation that could
351 lead to more accurate failure modeling. Expanding on the idea of friction, additional research
352 referencing frictional effects have been investigated by Staf et al. [57], where they determined the
353 frictional coefficient, of granular ceramics, was a function of pressure. They determined that particle
354 size and distribution were key factors in the compactability of a material. In contrary, the slope

355 in the curve for the larger particles ($487 \pm 98 \mu\text{m}$) in our experiments is shallower and does not
356 increase as abruptly.

357 We further explored triaxial effects by studying the Janssen-Walker theory [58]. The Janssen
358 coefficient is often used to simplify the relationship of granular stresses experienced in the axial and
359 radial directions, so that the complex behavior can be simplified and implemented, for example,
360 in existing material models to predict granular failure. Yousuff and Page [59] has shown this
361 radial-axial relationship when studying iron powders. In our experiments, we were able to conclude
362 for all three particle sizes tested the Janssen coefficient converged to a value of 0.23 (Figure 5).
363 This explains that the radial wall stress was approximately 1/5 of the axial load regardless of the
364 particle size. It was also observed that more variability exists for larger particle sizes ($487 \pm 98 \mu\text{m}$)
365 and porosities (50%+). Altogether limited research exists on particle size dependencies and the
366 observation highlights the importance of studying particle variability when researching granular
367 metals. The idea of incorporating length scales into failure modeling, reduces the variability in
368 predicting material behavior, but more importantly advances the forefront of powder material
369 design.

370 **4. Conclusion**

371 A confined uniaxial compaction technique was used to determine the triaxial characteristics of
372 granular SS316 as a function of ranging particle sizes. The results showed an influence of particle
373 size in the compaction curves where porosity is related as a function of hydro-static pressure.
374 In these experiments the bulk modulus was determined to be sensitive with respect to average
375 particle sizes. This is believed to occur due to the failure mechanism that is likely related to the
376 particle size, shape, and initial porosity. The smaller particles ($127 \pm 34 \mu\text{m}$) appeared to exhibit
377 higher flow stresses and underwent plastic deformation, while the larger particles ($487 \pm 98 \mu\text{m}$)
378 developed micro cracks which lead to fracture. Further research must be conducted to expand our
379 understanding of particle size effects on mechanical properties of the material, to better establish
380 failure regimes exhibited during triaxial loading conditions.

381 **5. Acknowledgments**

382 This research was sponsored by the Army Research Laboratory and was accomplished under
383 Cooperative Agreement Number W911NF-12-2-0022. The views and conclusions contained in this
384 document are those of the authors and should not be interpreted as representing the official policies,
385 either expressed or implied, of the Army Research Laboratory or the U.S. Government. The U.S.
386 Government is authorized to reproduce and distribute reprints for Government purposes notwith-
387 standing any copyright notation herein. The author(s) would also like to express their support to
388 Bernie Faulkner, Pouyan Motamedi, and Aidan Keaveny for their help in conducting the compres-
389 sion experiments, obtaining specimen images, and preliminary conceptualization respectively.

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Table 1: Chemical composition of Alfa Aesar SS316 powder

Element	% Mass
C	0.022
Cr	16.860
Mo	2.200
Mn	0.100
P	0.019
S	0.011
Ni	11.190
Si	0.730
Fe	0.0001

Table 2: Particle size characterization observing the diameter distribution and circularity of each particle

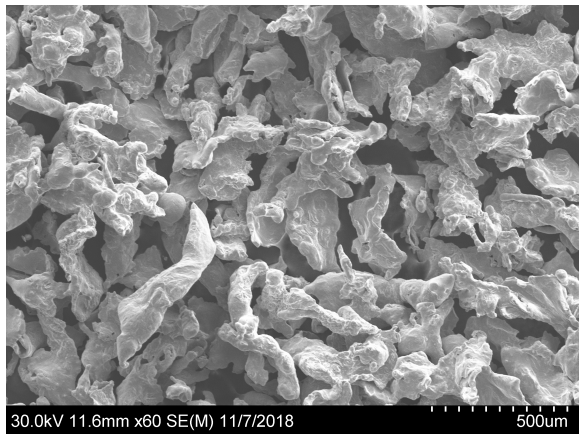
\emptyset [μm]	D_{10} [μm]	D_{50} [μm]	D_{90} [μm]	Δ	C
127 ± 34	94	125	169	0.601	0.765 ± 0.116
309 ± 88	233	314	364	0.417	0.643 ± 0.127
487 ± 98	413	480	549	0.283	0.626 ± 0.146

Table 3: Compression Experiment Parameters and Results

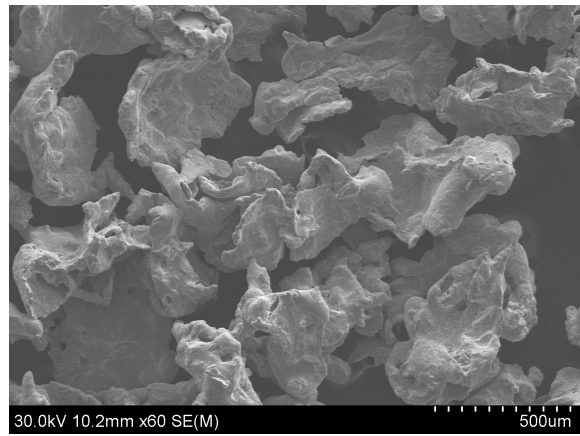
\emptyset [μm]	m [mg]	V_o [mm^3]	ρ_o [kg m^{-3}]	ρ_f [kg m^{-3}]	ϕ_i [%]	ϕ_f [%]
127	0.43	137	3132	6230	60	22
309	0.33	122	2699	6281	66	21
487	0.31	118	2632	6142	67	23

Table 4: Propagation of error

Relative systematic error	Particle size 127 μm	Particle size 487 μm
Initial volume (V_o)	0.77%	0.90%
Porosity (ϕ)	0.97%	0.96%
Hydro-static pressure (σ_{hyd})	11%	10%
Bulk modulus (B_{ep})	12%	2.9%
Janssen coefficient (K)	1.5%	1.7%
Transmission ratio (T)	6.5%	6.6%



(a)



(b)

Figure 1: SEM images of the SS316 powder at two different particle sizes: (a) $127 \pm 34 \mu\text{m}$ and (b) $487 \pm 98 \mu\text{m}$ respectively.

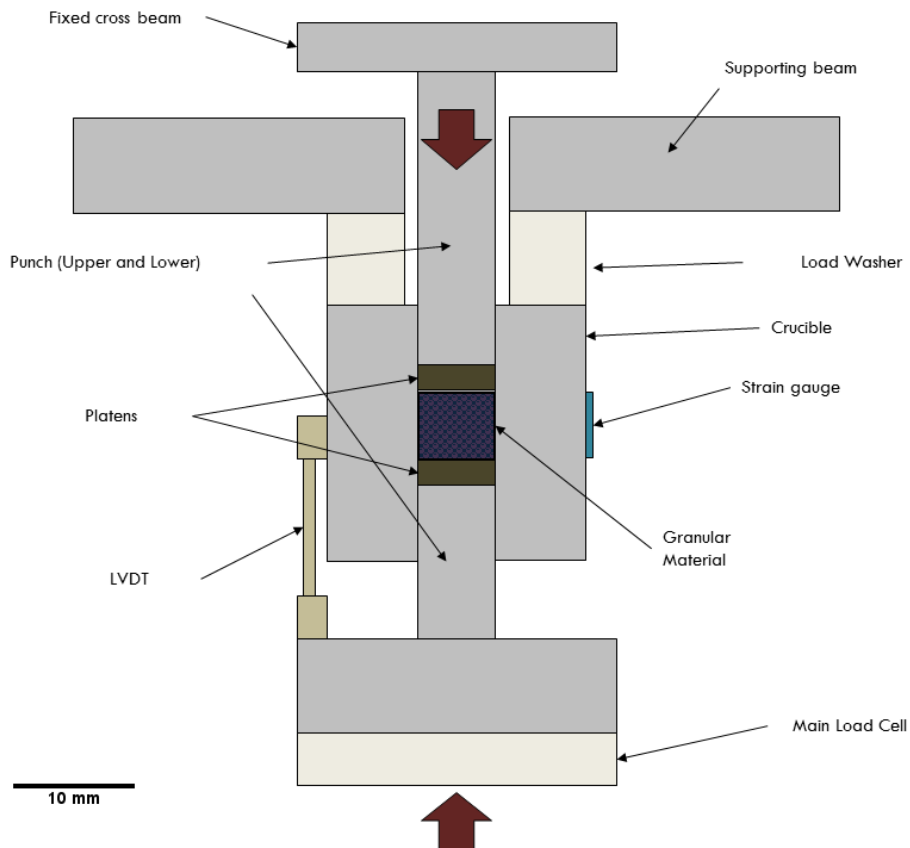


Figure 2: Schematic of the experimental set up illustrating key components and dimensions.

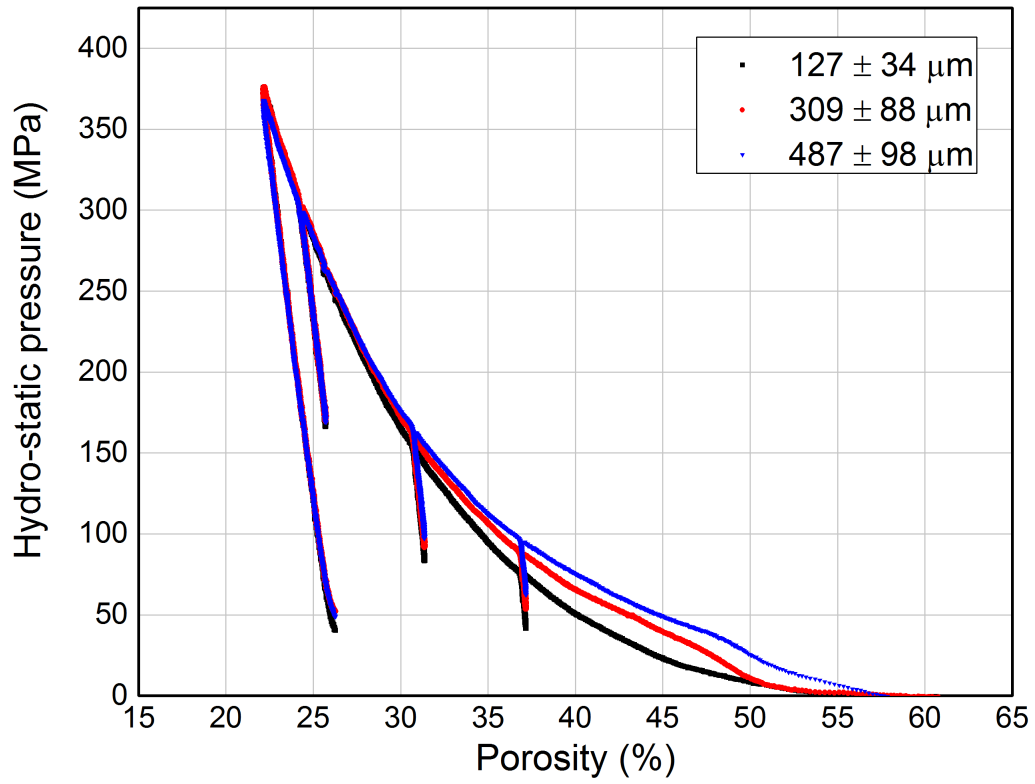


Figure 3: Quasi-static granular compaction response of SS316 powder for particle sizes of: $127 \pm 34 \mu\text{m}$, $309 \pm 88 \mu\text{m}$, $487 \pm 98 \mu\text{m}$.

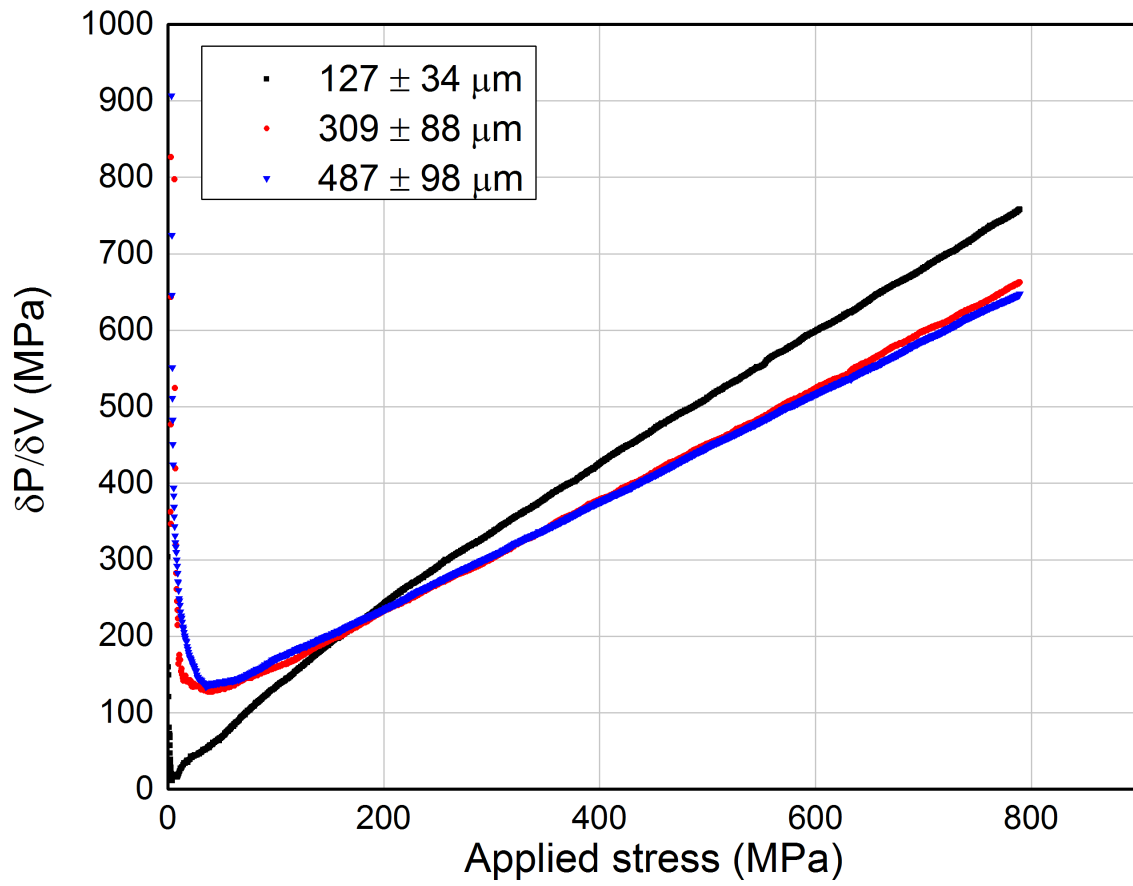


Figure 4: The bulk modulus as a function of applied stress for particle sizes of: $127 \pm 34 \mu\text{m}$, $309 \pm 88 \mu\text{m}$, $487 \pm 98 \mu\text{m}$.

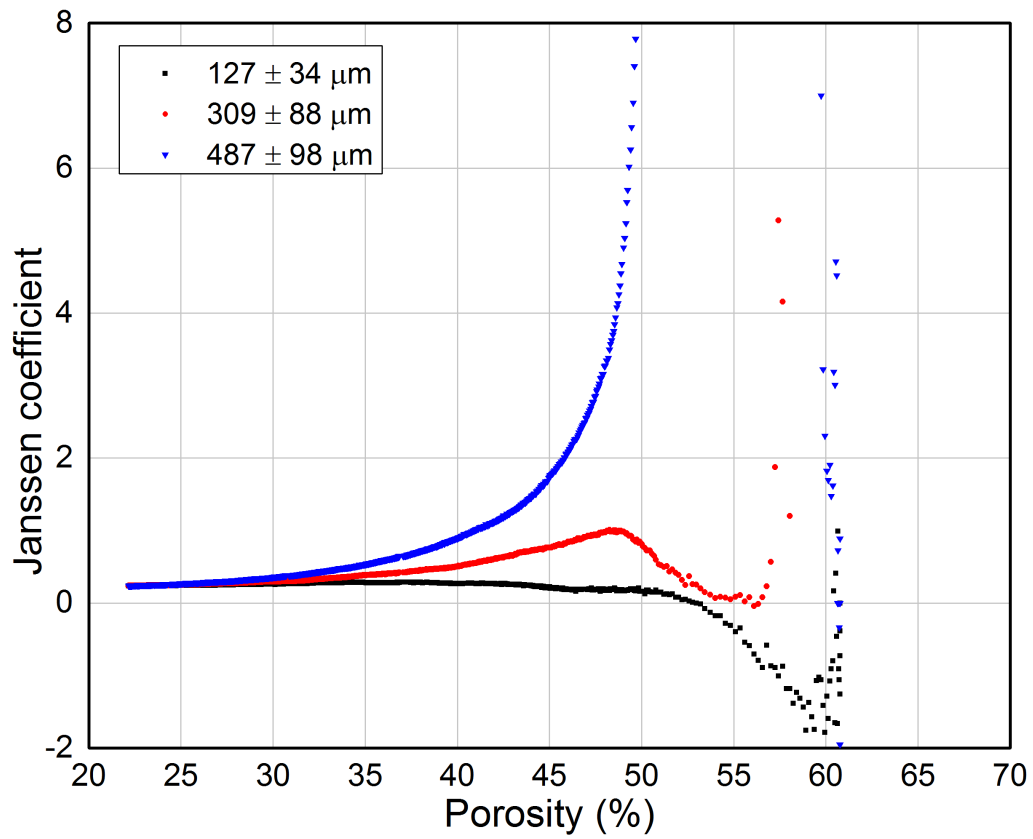


Figure 5: Janssen coefficient as a function of porosity for varying particle sizes: $127 \pm 34 \mu\text{m}$, $309 \pm 88 \mu\text{m}$, $487 \pm 98 \mu\text{m}$.

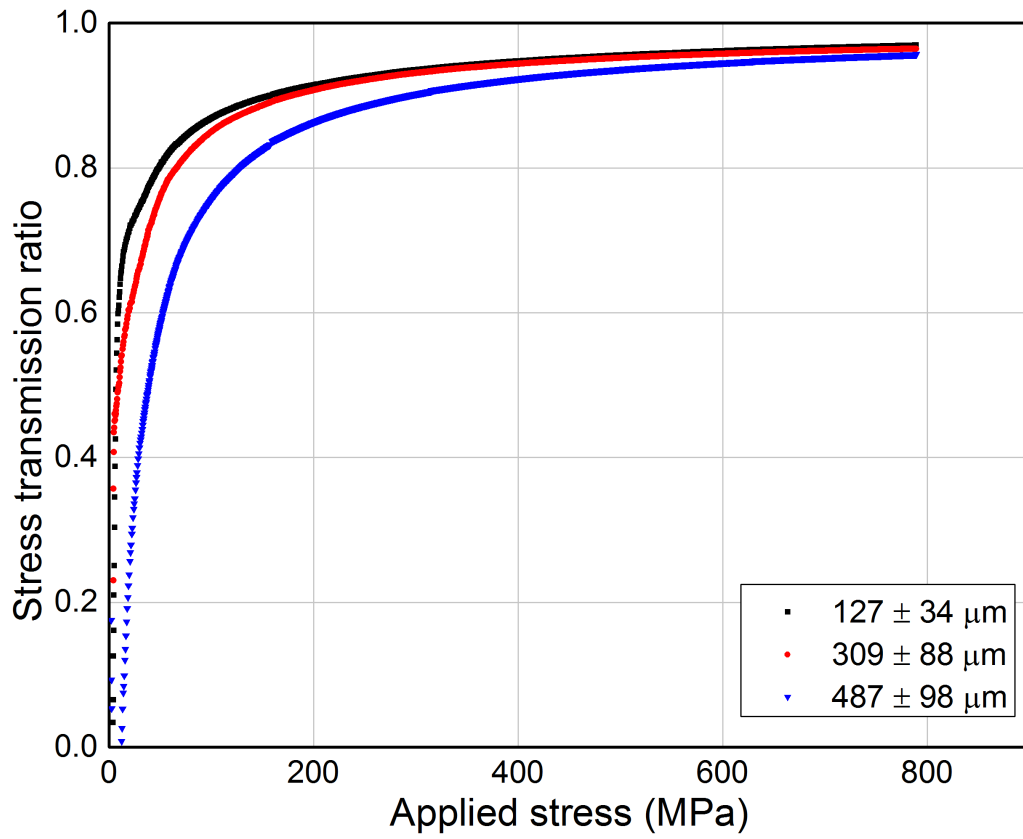
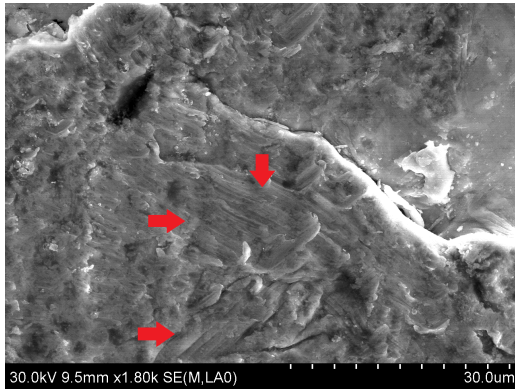
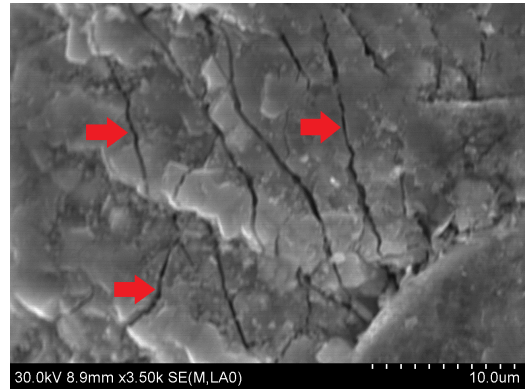


Figure 6: The transmission ratio relationship as a function of applied stress for different particle sizes: $127 \pm 34 \mu\text{m}$, $309 \pm 88 \mu\text{m}$, $487 \pm 98 \mu\text{m}$.



(a)



(b)

Figure 7: Post-experiment SEM images depicting failure mechanisms on the surface perpendicular to the applied load located at the top of the sample (in contact with the platen): (a) $309 \pm 88 \mu\text{m}$ particle size and (b) $487 \pm 98 \mu\text{m}$ particle size.

LaTeX Source Files

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