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#### THE UNIVERSITY OF ALBERTA

Design and Performance Evaluation of Tunnels and Shafts

by

Chik Kwong Wong (Ron)

#### A THESIS

SUBMITTED TO THE FACULTY OF GRADUATE STUDIES AND RESEARCH
IN PARTIAL FULFILMENT OF THE REQUIREMENTS FOR THE DEGREE

OF Doctor of Philosophy

IN

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Department of Civil Engineering

EDMONTON, ALBERTA
SPRING 1986

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# THE UNIVERSITY OF ALBERTA FACULTY OF GRADUATE STUDIES AND RESEARCH

The undersigned certify that they have read, and recommend to the Faculty of Graduate Studies and Research, for acceptance, a thesis entitled Design and Performance Evaluation of Tunnels and Shafts submitted by Chik Kwong, Wong (Ron) in partial fulfilment of the requirements for the degree of Doctor of Philosophy in Civil Enginhering.

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1.

### Abstract

The behawiour of tunnels and shafts in purely cohesionless and cohesive soils (elastic perfectly plastic) are studied by the use of the convergence confinement method (or ground reaction concept). The main objectives of this research work are to investigate the behavioral modes of tunnels and shafts in terms of yield initiation, propagation of yield zone and collapse mechanism, and to propose analytical techniques to predict the support stress-displacement relationship around these openings.

#### Tunnels

The excavation of a tunnel is simulated; in the two-dimensional plane strain model, by the proportional. unloading of the initial insitu stress around the opening. The conditions and consequences of ground responses around the tunnel are analyzed using numerical techniques, e.g., continuum mechanics, finite element method and limit state theory. Several possible modes of yield initiation and yield zone propagation are identified. It was found that the occurrence and development of each mode is governed by such parameters as depth ratio (H/a), strength properties ( $\phi$ , c), insitu stress coefficient (K\_) and internal support pressure  $(p_{\cdot}, \tau_{\cdot})$ . Two modes, Modes I (localized yield zones at the tunnel shoulders) and II (continuous yield domain around the opening), are identified. They are commonly encountered in soft ground tunnelling. Experimental results of model tests (Atkinson et al., 1974; Atkinson and Potts, 1977) confirm the existence of these modes (I and II). Furthermore,

analytical studies of Mode I and II behaviour reveal that a unique relationship exists between the support pressure and the displacement around the tunnel opening. The relationship can be interpreted by the use of the ground convergence curve. The approach proposed in this thesis is supported and verified by case histories of model tests and field measurements. It indicates a strong dependency among support pressure, mode of yielding and induced ground displacement, which is usually neglected in current design practice based on different semi-empirical approaches.

### Shafts

Ground deformations around circular shafts cannot be determined from currently available design approaches which are mostly derived on the basishof limit equilibrium techniques. The convergence confinement method (usually only applied to tunnels) with consideration of the mechanism of the shaft behaviour as a three-dimensional problem, is proposed as an analytical tool to predict the formation pressure on a shaft and the ground displacements. It was found that the behaviour of a shaft is governed by: (1) the mode of yield initiation dominated by the insitu stress. state and the soil strength, and (2) the extent of the yield zone generated by the wall displacements allowed during construction. Closed-form solutions for pressure-displacement relationships of cohesionless and cohesive grounds are presented. The results from the proposed technique compare well with those obtained from finite element analyses. The validity of the proposed

technique is also evaluated by comparison of the predicted response with the observed behaviour of shafts in model tests and actual field measurements. It was shown that the well established limit equilibrium methods provide minimum support pressures that are required to maintain stability, but these pressures are only encountered in the field if relatively large ground movements are permitted during construction with relatively poor ground control.

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#### 1 . INTRODUCTION

#### 1.1 Thesis Introduction

Due to rapid growth of population, underground construction has been active for the last 10 years in Edmonton, Alberta. Most underground projects have involved construction of shafts and tunnels for storm and sanitary sewer systems, power distribution systems, and light rail transportation.

The Geotechnical Section of the Civil Engineering Department at the University of Alberta have been aware of the benefits incurred from monitoring underground construction. This monitoring approach provides a method for validating empirical design techniques, evaluating stability and determining the mode of ground behaviour adjacent to the opening. Eisenstein and Thomson (1978) observed the geotechnical performance of a tunnel in till and predicted the field performance using the finite element method of analysis. Thomson and El-Nahhas (1980) measured the deformations of the temporary lining of two tunnels in till and clay shale. Medeiros (1979) studied the problem of magnitude and distribution of lateral pressure acting on a deep retaining structure excavated in till and sand by an approach integrating field measurement from a case history with finite element analysis. El-Nahhas (1980), Branco (1981) and Corbett (1984) carried out extensive field instrumentation to monitor the soil mass displacement and pressure formation around tunnel excavation. Kaiser et al. 1982) performed an investigation of a shaft in clay shale.

Korpach (1983) used stressmeters to measure stress changes near tunnel faces. Besides these field monitoring activities, Kaiser and Hutchinson (1982) studied the effects of construction procedure on tunnel performance using finite element methods. Eisenstein et al. (1984), Heinz (1985), and Eisenstein and Negro (1985) carried out finite element parametric studies on the ground responses near shallow tunnels.

Inspite of the above research work, there were few . quantitative analyses of the ground response that follows the correct sequence of stress relief of insitu stress, i.e., passage from elastic behaviour through yielding to collapse. Few attempts were made to relate stress relief with displacements for the whole spectrum of responses. Support pressure and displacement observations were often interpreted separately. In reality, soil pressures on tunnel and shaft linings are highly dependent on the soil displacements permitted during construction. This interrelationship has been recognized by many, e.g., Rabcewicz and Golser (1973), Egger (1975), Daemen (1975), in relating tunnel wall movement and support pressure, and more recently by Eisenstein and Negro (1985) for surface settlement and support pressure relationship. Current design methods usually treat with these two aspects individually, separating pressure prediction from displacement prediction. For example, limit equilibrium methods proposed by Terzaghi (1943), Berezanztev (1958) and Prater (1977) yield minimum support pressures to maintain stability of a shaft and do not predict displacements. In the field displacements are

usually controlled by selecting an appropriate factor of safety and through proper construction techniques. The current practice may lead to conservative or unsafe designs, as will be demonstrated later. Again, this aspect has been discussed for shallow tunnels by Eisenstein and Negro (1985).

In view of these limitations, a theoretical approach to solve this problem and to accommodate the dependency between stresses and displacements is attempted. This thesis is built on the concept of ground-support interaction or convergence confinement method which was originally proposed by Fenner (1938) and expanded later by Pacher (1964) and Rabcewicz and Golser (1973). This approach proposed and verified in this thesis is expanded to include various aspects as insitu stress, yield initiation and propagation, effect of gravity and collapse mechanism.

# 1.2 Convergence Confinement Method

#### 1.2.1 Introduction

The convergence confinement method (CCM) will be used throughout this thesis, as a conceptual framework for understanding the ground-support interaction in underground openings. This method is applied by constructing two characteristic curves of fictitious radial pressure versus radial displacement defining the ground response and the support reaction. The former describes the ground convergence in terms of the internal pressure relief (Ground Convergence Curve, GCC) while the latter relates the

confining pressure acting on the support to its deformation (Support Confinement Curve, SCC). The state of equilibrium of the ground support interaction is given, in a simplified fashion, at the intersection of the GCC and SCC, as illustrated in Fig. 1.1.

#### 1.2.2 Ground Convergence Curve (GCC)

The ground convergence curve gives the relationship between the radial displacements at the wall and a fictitious radial support pressure. This relationship at a specific point on the wall of the opening depends on (i) the initial insitu stresses, (ii) the strength-deformation properties of ground, (iii) the excavation method and (iv) the seguence of stress-relief for materials that are not linear elastic (see Fig. 1.2).

Fig. 1.1 shows a schematic ground convergence curve for a uniformly stressed circular opening in a homogeneous ground, along with the relationship between the extent of the plastic zone and the displacement. For K = 1, all GCCs originate at a single point of radial support pressure corresponding to the initial insitu stress. For a non-uniform biaxial stress field, separate GCCs can be developed for each point along the circumference of the opening. For small stress relief, the ground response is elastic, and the GCC is linear with the slope being proportional to the shear modulus of the ground. Further stress relief might induce yielding around the opening, and the GCC becomes nonlinear. For sufficiently strong ground, the opening remains stable without any support (Curve 1 for

elastic or Curve 2 for plastic ground). However, in weak ground yielding will propagate outward as the internal support stress decreases, and the pressure may reach a mainimum value (Curve 3) or start to increase (Curve 4) after an optimum support pressure at Point A.

Several modes of yielding can possibly develop near an underground opening: (i) localized yield zones (or shear bands), (ii) separate yield zones (diffuse yielding), (iii) a continuous (global) yield zone and (iv) combinations of the above. The plastic zone development depends primarily on the boundary condition and ground properties. The prediction of the mode of yielding is an unresolved area of research, but might be handled by the bifurcation theory, Vardoulakis (1985). One of the main objectives of this thesis is to identify typical modes of yielding that develop around tunnels and shafts, and to derive techniques to predict the behaviour of these modes.

Excessive yielding around an opening may induce a kinematically possible collapse mechanism under the effect of gravity, especially near the free boundary surface. A minimum support stress p due to this gravity effect is required to maintain the stability of the opening. The opening will remain stable as long as p is greater or equal to p. The gravity effect becomes dominant whem p reaches a critical value, p (Fig. 1.1). For strain-hardening ground (no reduction in strength during yielding), p will be approached asymptotically (Curve 3). For strain-weakening ground, the strength reduction in the yield zone will accelerate the gravitational effect, and

p increases with growth of the yield zone (Curve 4).

For ground with time-dependent strength or deformation properties, the displacement will increase with time. These time-dependent aspects have been studied by Ladanyi (1980) and Panet and Guenot (1982), and will not be pursued further in this thesis.

For the design of underground openings, we are concerned with both the support pressure and the ground displacement (e.g., for ground control). For the GCC of type (1), design and construction are simple and little difficulties are to be expected. For GCC type (2), we may reduce the support pressure at the expense of permenant plastic displacements, and hence ground control criterion may govern the design. For ground where the gravity effect becomes dominant, the support pressure will not decrease to zero and may even increase with large or "excessive" displacements. In this type of ground, the design must concentrate on the determination of the critical displacements (u\_\_) and the amount of gravity loading which depends on the type of collapse mechanism that is initiated. For good ground control, the displacements allowed during construction should be less than the critical in order to prevent excessive deformation that brings undesirable load increases and possibly destructive effect to surrounding structures.

Brown et al. (1983) presented a summary of currently available GCC formulations of different material models for the idealistic case of a circular opening in an isotropic stress field under two-dimensional plane strain condition.

Daemen (1975) included the gravity effect in his formulation. Application of these closed-form solutions to underground openings is very restrictive because of the inherent assumptions. For complicated boundary conditions and anisotropic stress fields as in shallow tunnels, the GCC can only be obtained by numerical simulation techniques (i.e., Eisenstein and Negro, 1985). However, severe problems are confronted in these approaches when excessive yielding takes place or the system approaches its limit state. Thus if the primary interest is the critical state (u ) or the limiting load, uses of these numerical simulation technique, e.g., FEM are not practically possible because the collapse of a geomechanical system is mostly associated with the formation and propagation of yield zones. Alternatively, the pressure required to prevent the collapse of an opening can be determined using analytical methods of the theory of limit state, but this approach requires the correct prediction of the formation of the yield zone and the mode of collapse. In this research, the mode of initiation and propagation of yield zones are studied and the collapse loads are determined.

Q

Thus, the strategy used to determine the entire GCC is to obtain the initial portion using continuum mechanics techniques (closed-form solutions) or finite element method, where appropriate, and to combine it with a terminal portion determined from theory of limit state (Fig. 1.1).

# 1.2.3 Support Confinement Curve (SCC) and Ground-Support Interaction

The support confinement curve defines the relationship between radial support pressure and radial displacement at the support ground interface (Curve 5, Fig. 1.1). The Location at which the curve originates depends on the time of installation, construction techniques and face effects. These three-dimensional effects can be accounted for in a simplified way by considering either an equivalent radial displacement (u) or an initial stress change (p) for the plane strain model. The shape of the SCC is associated with the strength-deformation properties of the support system and its activation. It is generally not a unique curve, except for K =1, circular opening and no gravity effects.

Hence, the way to obtain the equilibrium state of ground-support interaction as shown in Fig. 1.1 (intersection of the GCC and SCC) is over-simplified. The final equilibrium state depends on K, and the ground and support parameters (Pender, 1979). The study of this aspect is beyond the scope of this thesis. However, one can approximately locate the equilibrium state by considering two possible extremes in pressure formation on the support. For very flexible supports, symmetrical pressure distribution around the tunnel opening is appected. Averaging the pressure formation at ground support interaction points yields loci of the final equilibrium state (e.g., at the roof and floor). For very stiff supports, the final equilibrium state is comparable to the pre-support state. The uneven pressure around the support is

balanced by shear at the ground/support interface. This simplified approach suggested by Peck et al. (1972) was expanded by Einstein and Schwartz (1979).

Hoek and Brown (1980) presented solutions for a number of support systems. For simplicity, we assume linear-elastic behaviour for the support system.

### 1.3 Objective and Scope of this Thesis

The objectives of this thesis are (i) to investigate the mechanisms of ground behaviour near shallow and deep twenters and shafts, (ii) to develop an approach to assess behavioral modes of both shallow and deep tunnels and shafts, and (iii) to verify this approach by comparison with results of behavioral mode, support pressure and ground displacement from finite element analyses, field measurements and model tests.

The convergence confinement method, employed as a conceptual framework for this thesis, was introduced in this introductory chapter. Chapters 2 to 3 are concerned with the analysis of tunnels while Chapters 4 to 6 deal with shafts.

Chapter 2 is mainly concerned with the mechanism of tunnel behaviour and the formulation of the GCC. Numerical examples, generated by the finite element method, are compared with the proposed approach. Chapter 3 presents case histories from field and model test studies. The proposed technique is used to interpret some of the observations obtained in these case studies.

A review of limitations and assumptions of the currently available design methods for shafts is introduced

in Chapter 4 and followed by a detailed study of the ground deformation mechanisms near shafts. The CCM with the inclusion of gravity effect is suggested as a theoretical approach to shaft design. In this fashion it is possible not only to predict the formation pressure but also the ground displacements. Numerical examples generated by finite element analyses are compared with results predicted by the proposed technique.

Chapter 5 describes a well documented case history of an instrumented shaft constructed in Edmonton. Field measurements of displacements and pressures around the shaft are included in this chapter. Results of back analyses using the proposed technique and the FEM are compared with those obtained from field observations.

Chapter 6 presents case histories of model tests for shafts. The proposed CCM technique has been used to interpret these test results.

Conclusions on tunnel and shaft behaviour are given in chapter 7.

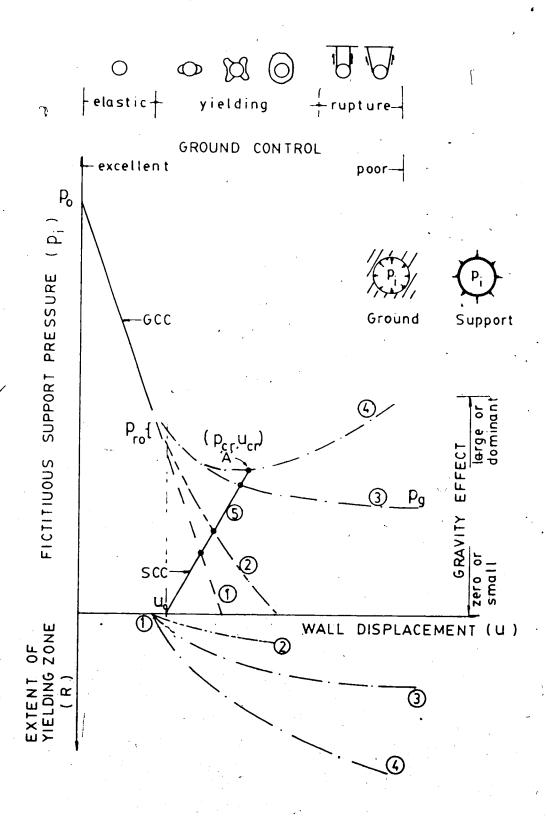


Figure 1.1 Convergence Confinement Method

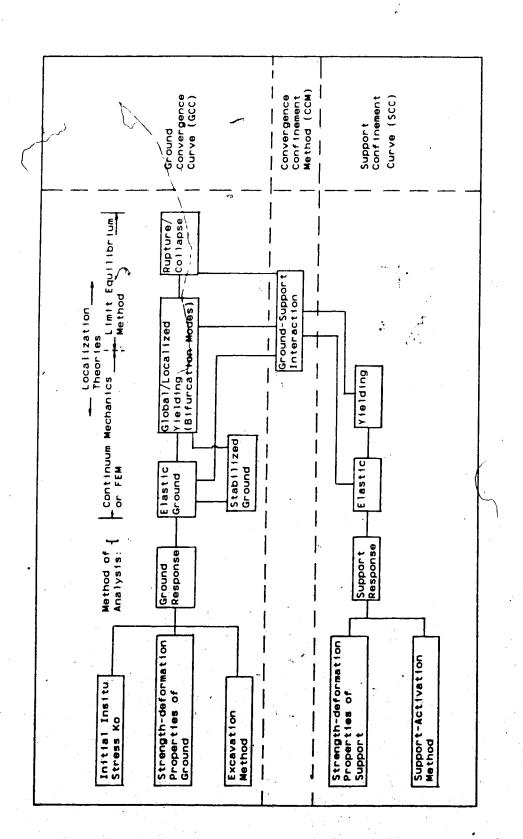


Figure 1.2 Conceptual Framework - Ground-Support Interaction

#### 2 . MECHANISM OF TUNNEL BEHAVIOUR

#### 2.1 Introduction

Driving a tunnel causes stress relief of initial insitu stress around an opening. The disturbed ground mass will displace and search for a new state of equilibrium.

Depending on the insitu stress, the strength of the ground and the construction method, this new state of equilibrium might be reached after undergoing elastic or plastic deformation, or collapse. Excessive displacements around a tunnel are prevented by installing a support system. Hence, the mechanism of tunnel behaviour must be governed by the distribution and magnitude of stress relief permitted near the opening during construction.

Ground movements due to tunnelling may cause destructive distortion or settlement to adjacent structures. For shallow soft ground tunnels, prevention of damage due to surface settlements becomes of prime concern and many design and construction decisions may be directed toward preventing excessive damage to structures or utilities near the surface.

The objectives of this chapter are to identify possible behavioral modes of shallow and deep tunnels and to propose analytical techniques to assess these modes. The proposed techniques will be verified by comparison with results of finite element analyses, model tests and field measurements.

#### 2.2 Ground Responses

#### 2.2.1 Simulation of Excavation

Consider the case of a circular tunnel, to be excavated to an initial radius a (Fig. 2.1a). The length of the tunnel is assumed to be much greater than its diameter and the section is remote from the excavation face. The tunnel can be treated as a two-dimensional plane strain problem. Excavation removes the insitu stress existing at the periphery of the tunnel opening. It is assumed that the ground response to this excavation process can be simulated by successively reducing a fictitious internal support pressure (radial and shear components). For a deep tunnel under hydrostatic insitu stresses, this fictitious support pressure may be approximated by a uniform radial stress equal to the overburden pressure at the tunnel axis. But for a shallow tunnel with non-uniform initial insitu stresses (K  $\neq$ 1), the fictitious support pressure is no longer uniform (p(roof)≠p(floor)) and is composed of radial (normal) and shear stresses,  $p_i$  and  $r_i$  ( $r_i=0$  at the roof, springline and floor).

# 2.2.2 Stress Relief during Excavation and Modes of Yielding

As the fictitious support pressure decreases, the wall will converge and stress redistribution involving a decrease in radial stress and an increase in tangential stress will take place. Arching action develops around the opening, as indicated by arrows in Fig. 2.1a. Yielding at the wall will be induced with a further decrease in support pressure if

the stress difference between tangential and radial stresses violates the failure criterion.

... The mode of yielding which is initiated at the wall depends solely on the initial insitu stress state, K . Fig. 2.1b shows two modes of yield initiation for (I) K  $\pm 0.5$  and (II) K = 1.0. For Mode I, the initial radial stress is greater than the tangential stress at the roof, and vice versa at the springline. As p decreases, the stress differences ( $\sigma$  - $\sigma$  ) becomes smaller at the roof (R) and larger at the springline (S). The soil element at the roof or floor always remains elastic while the soil element at the springline will fail. Thus the plastic zone will initiate at the springline for K=0.5. For K=1.0, the stress difference  $(\sigma - \sigma)$  at the roof and springline will grow approximately at the same rate. Hence, yielding will initiate around the periphery simultaneously unless localization occurs. For K >1.0, the plastic zone will initiate at the roof and floor because of high initial horizontal stress. This mode of yielding corresponds to Mode I rotated by 90°.

Once yielding is initiated at the wall, further stress relief in p will cause propagation of the plastic zone. How the plastic zone propagates depends on the displacement boundary conditions. For a deep tunnel where the external boundaries are remote from the opening, the problem can be simplified as an opening subjected to a confined insitu stress field. The plastic zone wilk propagate symmetrically.

Near'a shallow tunnel the free surface boundary affects the propagation of the plastic zone. Again, two

K conditions must be considered separately; for example, K=0.5 and K=1.0, as shown in Fig. 2.2. For K=0.5, the behaviour of the circular tunnel with plastic zones at the springline are comparable to a rectangular slot (trap door) of a width equal to the distance between two extremes of the plastic zones (w=2R). The weight of the soil block above the opening is supported by the shear resistance in the unyielded so. Thus, resisting shear stress will be mobilized along some inclined surfaces due to the downward movement of the soil block. The state of stress on a soil element along the inclined plane is plotted in a Mohr diagram in Fig. 2.2. The induced shear stresses due to gravity increase the stress difference between two principal stresses (the diameter of the Mohr circle) until failure occurs. Thus, localized plastic zones propagate upward from the springline and toward the free surface. The direction of the propagation of the yield zones, denoted by the inclined angle  $(\beta)$  (e.g., Fig. 2.5) depends on the distance to the free surface (H/a), K and the stress boundary at the opening. This will be proved later (Section 2.2.3.3). It is interesting to notice that the soil elements at the roof may still remain in an elastic state when a limit state is approached. Roof and surface settlements will increase at a comparable rate at this point.

The tunnel behaviour for K = 1.0 is different. After one initial yielding, the opening becomes completely surrounded by the yielding ground. The maximum extent (w) of the plastic zones at the springlines is smaller than that for K = 0.5, and the stress concentration is less severe.

Furthermore, because of relatively high initial horizontal stresses, the shear resistance against downward movements of a soil block is enhanced. Hence, the soil block above the opening has less tendency to displace downward (because the state of stress is farther from the yield surface) and high localized shear stresses do not develop to cause localized yielding. The plastic zone around the opening tends to propagate evenly outward. Its shape may be influenced by gravity causing more rapid propagation towards the free ground surface (Mode II).

The above illustration only shows two basic and idealized mechanisms of tunnel behaviour (I) formation of localized yield zones and (II) formation of a concentric (or an egg-shaped) plastic zone. However, other modes may also occur. Fig. 2.3 shows some possible modes which are of practical interest. It is assumed that in all these processes, the stress relief around the periphery of the tunnel opening is in proportional unloading of the insitu stress (p<sub>1</sub>, r<sub>1</sub>) and yielding is induced by an active failure mode. Modes I and II have already been discussed in detail. A critical K-value (K<sub>C</sub>) should exist between 0.5 and 1.0 at which yielding may initiate at the springline and two localized shear zones may coalesce to form a continuous domain around the opening as further stress-relief is allowed, i.e., Mode I develops into Mode II.

A third mode, Mode III occurs at K > 1.0 when yielding is initiated at the roof and floor, and a fourth mode, Mode IV similar to yielding in the 'trap-door' model test (Terzaghi, 1943), is observed if only vertical stress-relief

or displacement is allowed.

These four modes of yield initiation (Modes I to IV) can develop into other sub-modes, depending on the boundary conditions and ground parameters. In this thesis, only Modes I and II are of the primary interest. For Mode I, two separate yield zones at the springline may propagate in forms of localized lobes, and they may (Mode I-1) or may not (Mode I-2) reach the free surface before yielding is initiated at the roof. Mode I-1 usually dominates at deep tunnels whereas Mode 1-2 at shallow tunnels (Atkinson and Potts, 1977). The extent of yield zone at the roof propagates with further stress relief (i.e., Modes I-1.1 and I-2.1). At this stage, kinematic collapse may possibly occur at the roof. This collapse mechanism is governed by the formation of slip lines within the yield zone (i.e., upper bound solution). Mode I can also develop into a special Mode I-3 in which the yield lobes merge to form a ring with an unyielded core before the surface is reached. Mode II is common in situations where K is close to 1.0. The continuous yield zone will expand as stress relief proceeds, and two ear-shaped zones may be formed at the shoulders of the tunnel (El-Nahhas, 1980). Similarly to Mode I, two ear-shaped zones may or may not reach the free surface (i.e., Modes II-1 and II-2, respectively). In Modes I and II, the minimum fictitious support pressure around the tunnel opening is the one that prevents the kinematically possible collapse mechanism.

Various bifurcation processes discussed above govern the behaviour of tunnels in terms of their mode of yielding,

required support load (e.g., collapse load), and soil displacement (e.g., surface and subsurface settlements). Hence, it is essential to investigate what conditions and parameters govern each possible mode before other aspects like load and displacement development can be quantified.

### 2.2.3 Bifurcation Phenomena in Tunnelling

## 2.2.3.1 Introduction

The main purposes of this section are:

- 1. to identify parameters governing the modes of yielding near tunnels;
- 2. to locate boundaries and regimes of each possible mode; and
  - 3. to propose analytical techniques formulating the ground responses near the tunnel for different modes of yielding.

Formation of localized yield zones at the springline (e.g., Mode I) or a continuous plastic domain around the tunnel opening (Mode II) can be considered as a bifurcation problem. Bifurcation of a deformation process means that under certain critial states the deformation process may turn from one mode of behaviour to another entirely different mode.

Vardoulakis (1985) presented an excellent review of recent theoretical and experimental work done in the area of stability evaluation and bifurcation theory in soil mechanics. Some of the studies relevant to our objectives are briefly mentioned here. The study of bifurcation in

continuous mechanics was pioneered by Hill (1962). Following this fundamental work, Palmer and Rice (1973) formulated conditions for localization in connection with shear-band development in overconsolidated clays. Rudnicki and Rice (1975) also derived conditions for localization in pressure-sensitive, dilatant materials. Bifurcation analyses in deformable materials under tests of ideal stress conditions were performed by Hill and Hutchinson (1975), Young (1976), and Vardoulakis (1978, 1979, 1981). Needleman (1979) and Molenkamp (1985) investigated the effects of constitutive relationships of materials on bifurcation processes.

Bifurcation modes are governed by the constitutive properties of the soil, boundary conditions (stress and displacement), and geometrical and/or material imperfections. The mode of yielding near a tunnel, as discussed in Section 2.2.2, is dictated by the boundary conditions, e.g., insitu stress, internal support pressure and the free surface. Transition of one mode to another mode is due to the bifurcation process. Formulation of this bifurcation process near a tunnel is difficult because of the complexity of the boundary conditions. Vardoulakis (1985) suggested to use finite element methods or semi-inverse methods to solve this complex boundary value problem.

# 2.2.3.2 Strategy

While the three-dimensional conditions near the tunnel face may actually dominate, only the two-dimensional plane

strain conditions of a tunnel will be considered. Modes of yielding near a tunnel are linked by (i) strength parameters  $(\phi,c)$ , (ii) geometry and proximity of the free surface boundary (H/a), and (iii) boundary stress (K and  $p_i$ ,  $\tau_i$ ). The main objectives of the following analyses are to investigate under what combinations of the above governing conditions each mode dominates and to categorize them into typical modes of yielding. Instead of deriving exact solution for each possible mode which involves complicated numerical formulations, a semi-theoretical approach coupled with finite element analyses and interpretation of results from model tests and field observations will be used to establish boundaries among various modes. Continued research will be needed to formulate the method of analysis for each individual mode.

It was shown in the previous section that two basic modes can be identified under different K conditions: (I) localized yielding from the springline for K < K , and (II) continuous yielding around the opening for K > K . Mode I was observed in model tests (Atkinson et al., 1975; Potts, 1976; Cording et al., 1976). Reyes and Deere (1966) and Smith (1973) investigated the mode of yielding around a circular cavity under biaxial stresses and found that localized yield zones developed for K < 0.4. Lo et al. (1984) made use of Mode II in predicting the crown settlement for tunnels in soft clay, and found that their predictions compared well with field observations. Atkinson and Potts (1977) also carried out model tests in overconsolidated kaolin (at

strains were concentrated at the springline, i.e., Mode I. Their findings seem to contradict the proposition that Mode II occurs at K 1.0. However, it will be shown (Fig. 2.4) that because of the bifurcation process the mode may change depending on how and to what level the boundary stress p is reduced. Hence, from these observations in model and field tests, it is postulated that two modes are possible near a tunnel where K is less than or equal to 1.0. It is further postulated that these two modes are mutually exclusive and independent events under given boundary conditions, i.e., there is no overlapping of the two modes. Transition from one mode to another is, however, possible if changes in boundary conditions occur.

All important parameters governing the bifurcation process near a tunnel are included in Fig. 2.4 which is a plot of p/p versus K for a given  $\phi$  and H/a. Fig. 2.4 shows schematically regimes of possible modes. Steps in constructing the boundaries of the various regimes for different modes of behaviour are given below and followed by some numerical examples.

# Regime A

 $\mathcal{A}_{j}$ 

Line a-b is the limiting value to K, equal to the active pressure coefficient, K. The K values within this regime are inadmissable if the initial state of the ground is assumed to be elastic (i.e., the initial deviatoric stress of the K condition must be less than that of the

#### Regime B

This regime corresponds to conditions where no yielding occurs on elastic unloading. Line c-d denotes the limit for initiation of yielding. Since we are primarily concerned with cases of K <1.0, yielding will be initiated at the springline. Line c-d is determined by comparing the tangential-radial stress difference at the springline with the available strength of the ground. Schmidt (1926) derived equations for stress distribution around a circular tunnel in a stress field where the pressure increases with depth (Appendix B.1). The internal support pressure is assumed to be uniform around the circumference instead of percentage of the original insitu stress. This assumption leads to an underestimation of  $p_1/p_0$  for K <1.0. The underestimation magnifies as K decreases to K a.

# Regimes C. D. E. F. G and H

When the internal support stresses are slightly below those of Line c-d, yielding occurs at the springline.

Regimes C and H correspond to conditions where initial yieldings are localized at the springline. But there exists a critical K -value, K (Line g-j) which separates two different modes of yielding around the opening.

Determination of Line g-j will be explained later.

For K < K (Regime C) the two localized plastic zones of cr at the springline will develop into pairs of localized yield zones (Model I-1, Regime D). This development occurs at a p /p value defined by Line p-q which cannot be quantified at this stage. Further stress relief causes the top pair of

localized yield zones to propagate to the free surface.

These yield zones may (Regime F) or may not (Regime D) reach the free surface before yielding is initiated at the roof.

Line e-f separating Regime F from Regime D is determined by searching for the most critical case of Timit equilibrium state around the tunnel opening as demonstrated in Fig. 2.5 for Mode I. In this critical state, the two localized yield zones intersect the free surface. The assumption that the states of stress along the inclined planes AB and CD are at ultimate strength (or residual for strain-weakening ground) may be made. Thus, the required support pressure (pyserf, Fig. 2.4) to allow the localized yield zones to propagate to the free surface is determined using the limit equilibrium method. This approach will be given in details along with numerical examples in Section 2.2.4.1.

Line h-k signifies the point of yield initiation at the roof. Values of p<sub>i</sub>/p<sub>o</sub> for Line h-k could be determined using the same method as that for Line c-d, except for the roof instead of for the springline. Alternatively, the method presented by Detournay (1983) may be used. Hence, as a simplification it may be assumed that the elastic stress field around the tunnel is not affected by the yielding process at the springline. This effect of the yielding process is minimal as K<sub>0</sub> is close to unity where the extent of yielding at the springline is narrow (e.g., Regime H). In Regime E (Mode I-1.1), yielding has been initiated at the roof but the yield zones at the shoulder do not reach to the free surface. In Regime G (Mode I-2.1), the mode of yielding is similar to that of Mode I-1.1, except that the yield

zones at the shoulders have reached the surface.

#### Regimes J and K

In Regime H yielding only occurs at the springline.

This localized yielding will develop into a continuous domain (Regime J, or Mode II) as p decreases below support pressures of Line h-k. Further stress relief may again induce ear-shaped yield zones at the shoulders of the tunnel, as shown in Regime K. (Model II-1).

Line i-l represents the boundary of the transition from Mode II into Mode II-1 where ears develop. Within Regime J, the opening is completely surrounded by a continuous plastic zone. The solution to determine the extent of this continuous plastic zone comprises (i) determining elastic and plastic stress distributions satisfying all boundary conditions and (ii) searching the elastic-plastic interface by stress continuity (Oda and Yamagami, 1979; Detournay, 1983). Within the plastic zone stress equilibrium and failure criteria must be satisfied. These requirements impose conditions to seek for values along Line i-l. Steps to obtain Line i-l are given in Appendix B.2.

# Regime L

Once yielding occurs around the tunnel, kinematical collapse mechanisms bounded by slip lines within the yield zone can develop under the effect of gravity. The collapse mechanism becomes more obvious at the roof. Line m-n defines the minimum support pressure, p required to prevent any kinematically possible collapse at roof. If

 $p_i$  is reduced to support pressure below  $p_i$ , the collapse mechanism at the roof will be triggered and will propagate toward the free surface. Then, the state of stress above the roof has been disturbed and a new final equilibrium state must be reached. The required support stress to maintain this equilibrium state is designated as  $p_i$ . This mode is not shown in Fig. 2.4, and will be explored in detail later (Section 2.2.4.1).

#### 2.2.3.3 Examples

The approach described above (Fig. 2.4) is now used to investigate the modes of yielding near a tunnel in cohesionless soil. Governing parameters such as H/a, K and p /p are varied to investigate the modes of behaviour, and the results are shown in Figs. 2.6 to 2.8. The support stresses at the periphery of the tunnel are non-uniform, distributed in proportion of the insitu stress.

Comparisons among Figs. 2.6 to 2.8 reveal that areas for each regime are a function of H/a and  $\phi$ . The area of the elastic unloading Regime B increases in size as expected with increasing ground strength  $(\phi)$ . Regimes denoting the first mode (Mode I) contract in general as  $\phi$  increases. Regimes F and G diminish for most cases at H/a=18, i.e., for deep tunnels the localized yield zones will not reach to the free surface. The tendency of transforming from Mode II to Mode II-1 (area of Regime J) or the development of ears in the yield zone becomes less pronounced as  $\phi$  and H/a increases. This implies that the behaviour of a tunnel is

not properly approximated by the "hole-in-plate" theory unless the values of governing parameters fall within the area of Regime J.

On Fig. 2.7 ( $\phi$ =30°) are also plotted results of some numerical examples calculated by FEM, which are presented in Section 2.2.5. As expected, the yield initiation point for K =0.5 obtained from FEM (empty circles) are higher (by 5-18%) than the predicted owing to the different stress distributions of internal support pressures assumed for two cases (percentage of p in FEM and uniform in closed-form solution). However, the FEM results agree better with the closed-form solution as K approaches unity. With regard to aspects such as mode of yielding and yield initiation at the springline, the FEM and the predicted results compare well. Similar consistence between results obtained from the FEM and the predicted are also found for the case of K =0.82.

Results obtained from field and model test monitoring (Chapter 3) are also plotted in Figs. 2.7 and 2.8 for comparison, which will be discussed in detail later.

# 2.2.4 Formulation of Ground Response

This section mainly explores analytical methods to depict the ground responses in stages from the state of yield initiation to collapse under the progressive stress relief in the excavation. We are interested not only in the ground convergence curve at the tunnel wall but also in the relationship among crown displacement, surface settlement and stress relief. Since two types of modes in ground behaviour, I and II have been identified, it is appropriate

to investigate them separately. Although the following analyses are based on elastic perfectly plastic cohesionless ground, basic principles and reasonings can be applied to other types of ground with different constitutive laws.

## 2.2.4.1 Mode I

Fig. 2.9 presents a typical ground convergence curve at the roof for Mode I (Mode II also). On this figure are also plotted modes of yielding and methods of analysis. The changes in states of stress above the roof described by K-coefficient are also shown in Fig. 2.9. K is the average ratio of horizontal stress to vertical stress at soil elements above the crown. Each portion of the ground convergence curve is depicted in detail in the following.

#### Portion o-b

Portion o-a represents the elastic unloading response of the ground and Point a indicates the initiation of yielding at the springline. Portion a-b reflects the nonlinear ground response when the yield zones propagate outward with stress-relief.

Portion o-b may be determined using a continuum mechanics approach. But the influence of such factors as free surface boundary, gravity, non-hydrostatic loading and localized plastic zones must be considered. They impose numerical difficulties in obtaining closed-form solutions. Duddeck (1980), Eisenstein et al. (1984) and others suggested the finite element method as an analytical tool to

evaluate the ground response to tunnel excavation. Eisenstein and Negro (1985) developed a two-dimensional finite element program to simulate the tunnel behaviour for ground of non-linear hyperbolic elastic stress-strain response with time-independent properties. For shallow tunnels in cohesionless soils, they encountered severe convergence problems in equilibrium iteration when the stress relief was more than 60% of the original insitu stress. This is one of the numerical inadequacy in finite element analysis when the limit state is approached. In reality, however, support pressure much less that 40% of insitu stress have been reported from field measurements (Eisenstein et al., 1981; Branco, 1981; Corbett, 1984). Hence, it is fruitful if one can explore the ground response beyond this state.

#### Portion b-e

This portion describes the ground responses due to the stress relief from yielding to ultimate behaviour. Point c represents the state at which the yield zones just intersect the free surface, and the support pressure at this state is denoted by p. At Point c, the soil block above the roof may or may not remain elastic. The support pressure can be further reduced to a lower p. (Point d) at which the collapse occurs at the roof. There exists a point between Points c and d corresponding to the point of yield initiation at the roof (i.e., Line h-k in Fig. 2.4). From the initial Point o to Point d, the tangential arching action above the roof (indicated by K-value) increases with

stress-relief. But, beyond Point d the roof collapse mechanism is triggered and will propagate to the free surface. This mechanism will induce stress redistribution above the roof and the resultant K-value will drop. Hence, a sudden increase in support pressure from Point d to Point e results mainly from the decrease in tangential arching. For strain-weakening ground the whole portion a-e will be further shifted upward because of strength reduction associated with plastic straining in residual state.

With correct predicted modes of yielding, the support pressures at Points c, d and e can be determined by use of theory of limit state. Currently, no solution exists for the corresponding displacement at each point. Thus, some simplifications are required to construct Portion b-e. This can be done by extrapolating Portion a-b to Point d, i.e., following the negative slope at Point b. This assumption may lead to conservative design because Portion b-d should become flattened for yielding ground (i.e., the slope (negative) increases with displacement).

Portion b-e is of practical importance to tunnelling design because it provides information on marginal safety from collapse mechanism and also the ground control.

Point c (p )

At Point c, the yield zones just intersect the free surface forming a wedge as illustrated in Fig. 2.5a. The required support pressure p can be estimated by limit equilibrium method.

Several assumptions regarding the failure surface and the stress destribution along it must be made such that an overall equation of equilibrium, in terms of stress resultants, can be written. Consider the equilibrium state of a soil block (ABCD) above a tunnel (Fig. 2.5a). Planes A-B and C-D inclined at  $\beta$  to the horizontal describe two failure surfaces giving the most critical condition of maximum required support pressure. Due to symmetry, no shear stress exists along Plane E-F. Considering the horizontal and vertical force components acting on the soil block, the equilibrium gives:

$$F + p_a - (W - p_a) \tan(\beta - \phi) = 0$$
2.1

where: F - force acting on the plane E-F  $p_{h} - \text{equivalent uniform support pressure in} \\ \text{horizontal direction } (= \text{K} \left[ (\text{H-a/2}) \gamma \right] \eta \text{ where } \eta \text{ is } \% \\ \text{of insitu stress})$ 

p - equivalent uniform support pressure in vertical direction ( =  $[(H-a/2)\gamma]\eta$  where  $\eta$  is % of insitu stress)

Differentiation of Eqn. 2.1 with respect to variable  $\beta$  yields a maximum value for p<sub>1</sub>. The angle  $\beta$  calculated in this manner should correspond to the direction of yield zone propagation from the springline to the surface.

Solution of Eqn. 2.1 requires knowledge of the force

F. This force is equal to the resultant of the horizontal insitu stresses before tunnel excavation plus an increase in stress due to tangential arching resulting from

stress-relief (see Fig. 2.5b). Thus  $F_r$  can be expressed as:

$$F_r = 1/2(K_r)\gamma(H-a)^2$$
 2.2

where:  $K = K + \Delta K$ .  $\Delta K$  and K are coefficients used to approximate the increase in tangential stress and the resultant respectively (Figs. 2.5 and 2.9). It is important to note that the ratio of horizontal stress to vertical stress (K) actually varies along the depth above the crown.

Assuming even stress redistribution around the opening (Fig. 2.5b), the coefficient K can be calculated in terms of the variables H/a and K (see Appendix B.3). It is plotted in Fig. B.2. These K values depend on the percentage of stress-relief of the initial insitu stress. Thus, the maximum K is governed by the minimum support pressure, p. Theoretically, K can reach the passive resistance coefficient, K (i.e., the whole soil block above the roof are in passive yielding). From Fig. B.2, it can be seen that the case of K approaching K only occurs near shallow tunnels and the maximum K is usually dictated by Pic.

Eqn. 2.1 provides information not only on the support pressure, p but also on the direction of the propagation of the yield zone,  $\beta$ . A computer program (TUN5, Appendix F) is written to calculate p and  $\beta$ . Fig. 2.10 shows the angle  $\beta$  for ranges of H/a and  $\phi$  values. It can be seen that this mode of yielding (Mode I-2) is only admissible at certain values of H/a, K and  $\phi$ . The K-value has a limiting value of active pressure,

K. At high H/a (deep tunnel) and  $\phi$  values (strong ground) the tendency of yield zone to propagate to the free surface is less pronounced. With K approaching to unity, the mode of yielding is Mode II. The angle  $\beta$  depends on K and H/a as weld. For a shallow tunnel with high K, the yield zones are vertical.

The German Tunnelling guidelines (e.g., Duddeck, 1980 and 198/2) differentiate between 'shallow' and 'deep' tunnels for H/a<5 and 'deep' tunnels for H/a>7. For 'shallow' tunnels, an equivalent continuum model (no embedding allowed at the crown, e.g., Duddeck and Erdmann, 1982) without reduction of ground pressure at the crown is recommended for lining dimensioning. For / deep' tunnels, a continuum model (fully embedded, e.g., Muir-Wood, 1975) with some reduction of ground pressure may be appropriate. Careful examination on these guidelines reveals that one important factor governing the tunnel behaviour is neglected; the displacement permitted during construction, u been proven that the ground pressure on a support and the behavioral mode of a tunnel depend on u and hence the structural model. Therefore, the structural models proposed in these guidelines are only valid under certain conditions. It is of practical importance to explore these conditions and to locate the limitations of these guidelines.

Limits between 'shallow' and 'deep' tunnels are plotted on Fig. 2.10 as 'A' and 'B' respectively. If Mode I-2 takes place (i.e., yield zones reach the ground surface),  $\beta$  are in ranges of 75° to 90° for 'shallow' tunnels and 60° to 80° for deep tunnels. The boundaries of these ranges depend on

φ. The structural model for 'shallow' tunnels recommended by the German guideline is conservative in terms of ground pressure. From Fig. 2.9 (GCC), full overburden pressure is only experienced to no displacement is permitted during construction. This condition is seldomly encountered in the field. Ground pressure may increase for u Point d, Fig. 2.9), but the induced pressure is still less than the overburden pressure. A partially embedded model (no embedding at the crown) is justified because in Mode I-2 the soil block bounded by two yield zones intersecting the ground surface has tendency to displace downward and exert pressure on the liner at the crown. Hence, the soil block above the crown should not be treated as part of embedding. For structural model of 'deep' tunnels, the ground pressure can be reduced with the increased displacement. The reduction depends on u and GCC. From Fig. 2.10, it is observed that the fully embedded model is only valid for Modes I-1 and II. If yield zones intersect the ground surface (e.g.,  $\phi=20^{\circ}$ , K\_<0.60;  $\phi=30^{\circ}$ , K\_<0.40;  $\phi=40^{\circ}$ , is greater than u. (i.a., roof collapse has been initiated), then the partially embedded model instead of the fully embedded should be used.

It is interesting to observe that K of normally consolidated soils given by  $(1-\sin\phi)$  (Brooker and Ireland, 1965) falls within limits 'A' and 'B' in Fig. 2.10. This suggests that the German Tunnelling guidelines are appropriately designed for normally consolidated soils provided that conditions stated above are satisfied.

Fig. 2.11 is a plot of normalized support pressure  $(p_y/\gamma h)$  versus normalized depth (H/a) when the yield zone reaches the ground surface. This figure indicates that  $p_y$  is a function of K and  $\phi$ .  $p_y$  decreases with increasing K and  $\phi$ , implying that a small amount of stress relief of the insitu stress will cause the yield zones to reach the free surface in weak grounds under low K. Hence, tunnelling in these conditions is much more risky as intuitively expected. On Fig. 2.11 is also plotted the support pressure  $p_i$  corresponding to the initiation of roof collapse mechanism. For conditions of  $p_i > p_j$ , Model I-1 is impossible, and Mode I-2 or Mode II-1 takes over depending on K (see bifurcation modes, Fig. 2.4).

#### Poínt d (Pic)

At this point yielding at the roof has occurred and possible collapse mechanisms defined by slip-lines within the yield zone may develop. The critical mechanism can be found by selecting any possible modes and performing an appropriate work rate calculation. The accuracy of this calculation (upper bound solution) will depend on the proximity of the assumed mechanism to the real one.

Atkinson et al. (1975) observed the configuration of the collapse mechanism in model tests and derived an upper bound solution.

$$p_{ic}/\gamma H = (a/2H\cos\phi)(1/\tan\phi + \phi - \pi/2)$$

provided that H/a is equal to or greater than  $1/\sin(\phi)$ . This

equation is true only for material with an associated flow rule and  $\psi=\phi$  where  $\psi$  is the angle of dilation (Hansen, 1958).

The values of  $\vec{p}$  normalized to overburden pressure ic ( $\gamma H$ ) are plotted in Fig. 2.11. Actually,  $\vec{p}$  are independent of the depth of overburden.

# Point $e(p_f)$

Beyond Point d the roof collapse mechanism has been initiated. The extent of the propagation of this collapse mechanism depends on the displacement allowed at the roof. At the ultimate behaviour, the mechanism will reach the free surface and the soil block above the roof will be in the limit state. The support pressure  $p_{fc}$  at this state can be found by the use of the soil arching theory (Terzaghi, 1943). The arching action will be developed against both sides of the wedge (Fig. 2.5a). The mathematical treatment on arching is given in Appendix B.4.

Normalized support pressures  $p_f$  are plotted against the depth ratio (H/a) for different shape of wedges ( $\beta$ ) and friction angle ( $\phi$ ). The  $p_f$  values also depend on the assumed earth pressure coefficient at the sides,  $K_s$  (Handy, 1985). Fig. 2.12 with  $K_s = K_s$  shows that the  $p_f$  becomes fairly constant for H/a>5. This agrees well again with the guidelines given in the German Tunnelling Handbook (suggesting full overburden at the crown if  $p_f$  is  $p_f$  and  $p_f$  and  $p_f$  is  $p_f$  and  $p_f$  and  $p_f$  are plotted against and  $p_f$  and  $p_f$  are plotted against agai

It is interesting to realize that two different types of arching action occur before and after this initiation of the roof collapse (Point d). For the state before Point d,

the horizontal (tangential) stress is the major principal stress and the vertical (radial) is the minor within the yield zones above the roof. In this soil arching, the catenary must be convex upward as illustrated by Fig. 2.13a. In contrast to this, an inverted arch develops after excessive movements (Fig. 2.13b). This mode beyond Point d was well recognized and explained by Handy (1985). In this sagging state, soil arching may be depicted as a trajectory of minor principal stress that approximates a catenary (the vertical stress is major principal stress). For arching action to be supportive, the catenary must dip downward. The configuration of the catenary will change from convex upward to concave downward after passage of Point d. This transition signifies an abrupt increase in settlements above the roof.

From the GCC the relationship between the support stress and displacement at the roof can be obtained except for the ultimate state where the total displacement cannot be predicted. The next essential step then is to relate this to the roof settlement profile above the tunnel. This aspect will be discussed in detail together with Mode II.

#### 2.2.4.2 Mode II

If yield initiation and propagation take place during unloading under K equal or close to unity, the tunnel opening becomes completely surrounded by a continuous, approximate concentric plastic zone. This situation follows the mechanism of a "hole-in-plate" model, and permits application of continuum mechanics and theory of plasticity

to predict the internal pressure-convergence relationship.

Daemen (1975) adopted this approach and included the gravity effect within the plastic zone to determine the ground convergence curves for deep rock tunnels. Lo et al. (1984) also used the same approach to predict the roof settlements due to tunnelling in soft clays.

In spite of the above studies, the continuum mechanics approach has not been used as an analytical tool to predict the behaviour of shallow tunnels, which are dominated by the gravity effect and complicated by the free surface boundary. In this thesis the technique proposed by Daemen (1975) will be adopted and expanded to include the effect of the free surface.

Considering a case of a tunnel as shown in Fig. 2.14a, i.e., a two-dimensional plane strain problem under K = 1.0, a plastic zone of radius R will develop around the opening if the ground strength is exceeded. There exists a relationship between the internal support pressure p and the plastic zone R at the periphery of the opening.

For a shallow tunnel the true conditions can be replaced by a simplified model as shown in Fig. 2.14b. The stress distribution within the plastic zone are governed by the equilibrium equations.

$$\sigma_{t} - \sigma_{r} - r(d\sigma_{r}/dr) - d\tau/d\theta + r\gamma\cos\theta = 0$$
 2.4

$$d\sigma_{t}/d\theta + 2\tau + r(d\tau/dr) - r\gamma \sin\theta = 0 \qquad 2.5$$

Assuming  $\tau=0$  at the roof, springline and floor, Eqns. 2.4 and 2.5 can be reduced to:

$$\sigma_{t} - \sigma_{r} - r(d\sigma_{r}/dr) \pm r\gamma = 0$$
2.6

where: + for floor, - for roof and  $r\gamma=0$  at springline.

For the ground that can be characterized by a cohesionless Mohr-Coulomb yield criterion:

$$\sigma_1/\sigma_3 = \sigma_1/\sigma_r = \tan^2(\pi/4 + \phi/2)$$
 2.7

Integration of Eqn. 2.6 along with Eqn. 2.7 and continuity of radial stress at the elastic-plastic boundary leads to the following relationships for (p<sub>i</sub>).

$$p_{i} = p_{0}[(1-\sin\phi)(a/R)^{2\sin\phi/(1-\sin\phi)}] \pm a\gamma[(1-\sin\phi)/(1-3\sin\phi)],$$

$$[1-(a/R)]^{(3\sin\phi-1)/(1-\sin\phi)}$$
2.8

where: + for floor, - for roof and last term reduced to zero for springline.

In Eqn. 2.8, the parameter p is a function of R and varies at the roof, springline and floor as shown in Fig. 2.14b. Also, in Eqn. 2.8 it can be seen that the first term is primarily concerned with strength mobilization while the second term describes the effect of gravity within the plastic zone. This term is more significant in strain-weaking ground because less strength can be mobilized over the total extent of the plastic zone.

In the roof area the gravity acts against the internal pressure thereby increasing the internal pressure required to achieve equilibrium. Hence the gravity term is positive. In the floor area, gravity acts as a stabilizer and thus reduces the internal pressure. Hence the gravity term is negative. There is no gravity effect at the springline where the direction of the gravity is perpendicular to the radial pressure.

Integration of Eqn. 2:6 with known stress boundaries given by Eqn. 2.8 provides stress distributions at the roof, springline and floor. When these stress distributions are coupled with material deformation characteristics (Brown et al., 1983), one can obtain separate ground convergence curves for the roof, springline and floor. These expressions for the GCCs depend on material properties and are quite lengthy. Formulations proposed by Ladanyi (1974) are considered in this thesis and details are given in Appendix B.5.

Fig. 2.15 presents a schematic plot of the pressure-displacement and pressure-extent of plastic zone relationships. Each GCC originates from its initial insitu stress state, i.e., the floor has the highest value. After an initial linear response it becomes non-linear with yield initiation. As the plastic zone increases, the gravity effect becomes dominant, reduces support stress at the floor and increases it at the roof. The floor may be left without support depending on soil strength. The same reasonings on the gravity effect can be applied to the p<sub>1</sub>-R relationship (given by Eqn. 2.8) but the extent of plastic zone at the

roof will stop when it reaches the ground surface at R=H.

Before the above technique is applied to determine the GCC, it is important to understand the limitations inherent in this method.

- Simplifications have been made in the derivat@ons of Eqn. 2.8. Based on the "hole-in-plate" theory, the equation is obtained from its own boundary stresses, independent of other boundary stresses as shown in Fig. 2.14b. In addition the assumption of τ=0 is only valid along vertical axes at the roof and floor, and at a horizontal axis through the springline. This assumption leads to an overestimation of p at the roof, but an 'underestimation at the floor.
- Eqn. 2.8 is formulated assuming a plastic zone of 2. constant radius formed around the opening. The derivation of the p.-u relationship follows the path (1) shown in Fig. 2.14. At excessive yielding, the support pressure distribution may result in a maximum at the roof and a minimum at the floor, which is in contrast to the assumed stress distribution in excavation simulation (stress-relief in a percentage of the original insitu stress). Hence the use of Eqn. 2.8 is restricted to a certain limited range of p, values. This is further confirmed by the studies of the modes of yielding near tunnels. From Figs. 2.6 to 2.8, the mode of yielding with a continuous plastic zone is limited only to Regime J (Mode II). The area of this regime increases, as H/a increases, implying that the proposed technique will yield better results for the

behaviour of deep tunnels.

3. It is of practical interest to find the extent of the plastic zone around the tunnel. For a rigid liner (Case 1), the plastic zone is fairly concentric, resulting in a differential support pressure between at the roof and at the floor. This unequal pressure may be transferred to the ground through shear at the support-ground interface. This shear effect will affect the shape of the GCC (Curtis, 1976), which is beyond the scope of this thesis. If flexible liners are used in tunnel supports, no bending action will be developed within the liner and symmetrically distributed support stress will result. This condition is similar to Case 2 in Fig. 2.15. Symmetrical support pressure distribution induces an egg-shaped plastic zone, i.e., large yielding occurred in the roof area.

From the above discussion, it can be seen that the application of the proposed technique (Mode II) is restricted by its inherent assumptions. Formulation of ground responses beyond the regime of Mode II requires techniques similar to Mode I (Fig. 2.9).

2.2.4.3 Vertical Settlement Profiles above Tunnel
(Modes I and II)

Analytical prediction of the settlement profile above the roof is complicated by the proximity of the surface boundary. Finite element methods seem to provide the only approach to predict the displacement pattern around the tunnel, but several attempts using continuum mechanics have

revealed that the surface settlements are normally underestimated, e.g., Atkinson and Potts (1977). Lo et, al. (1984) used a semi-empirical approach to estimate the surface settlement and succeeded to predict the crown displacements using continuum mechanics. Because of these problems, a semi-theoretical approach supported by case histories is used here to predict the settlement profile above the crown.

So far discussed are the relationship between the crown convergence and support pressure. There exists also a similar relationship for the surface settlement, i.e., the surface settlement is governed by the crown settlement S and thus the internal pressure p. The p-S-S relationship is a steady changing function until kinematically possible collapse mechanism takes place (e.g., Point d, Fig. 2.9). The surface settlement beyond Point d increases abruptly and become excessive. A prediction of settlement after a collapse mechanism has been initiated is not dificult. Hence, the following analysis is concentrated only to states before collapse.

Considering only the vertical ground displacement at the ground surface, it is well established that the settlement profile may be represented by a Gaussian distribution curve of form (Fig. 2.16a)

$$S = S_{S} (exp)[-(1/2)(x/i)^{2}]$$

2.9

where: i is the point of inflection.

For a given tunnel, the surface settlement profile (Eqn. 2.9) depends on p and S. Assuming the surface displacement profile outside a distance of w (between 2.5i and 3i, Cording et al., 1976) remains uninfluenced by the progressive stress relief a relationship of S -i is derived by putting S=S at x=w: (where w=H(cotanβ')+ a)

$$S/S = \exp\{1/2[(1+(H/a)\cot an\beta')/(i/a)]^2\}$$
 2.10

In Eqn. 2.10, the variable S is a measure of  $\overset{\circ}{S}$  or  $\overset{\circ}{p}_{i}$ , and the variable i is of the settlement profile. The parameter  $\beta$ ' (approximately equal to  $\beta$ ) defining the trough width is dependent on the mode of yielding or K as discussed earlier (Modes I or II).

Fig. 2.16b is a plot of normalized S versus i for H/a=3 to 10,  $\phi$ =20° to 40° and  $\beta$ =70° to 80°. The  $\beta$ -values are extracted from Fig. 2.10 of Mode I. For shallow tunnels i decreases to a limiting value with a small increase in S, i.e., small crown displacement or stress relief causes the yield zone to propagate to the surface. At deep tunnels i also decreases with S, but at a slower rate. Fig. 2.16b indicates that the shape of the settlement trough for a given cohesionless ground is dependent on H/a and p (or S) as well. The  $\phi$  parameter has less extent of influence.

It must be expected that Modes I and II display distinctly different and unique features in their settlement profiles because they represent different modes of yielding. Consider Modes I and II and allow a constant crown settlement

 $S_{c}$  in both cases (Fig. 2.17). If  $S_{c}$  is small, the extent of the yield zone is small. The vertical settlement profiles of Modes I and II are initially very similar, but the magnitude of settlement is larger in Mode I than in Mode II. Small displacement occurs in the elastic zone and large plastic straining within the yield zone (Fig. 2.17a and b). For excessive S (i.e., the yield zon'e reaching the surface), Modes I and II exhibit distinct differences in vertical settlement profile above the crown. In Mode I, two localized shear planes develop and the soil block displaces toward the opening as a rigid body. The soil block remains elastic so that the differential strain and displacement between the crown and the surface is small. Hence, the ratio of surface to crown displacement  $(S_S/S_C)$  is closer to and tends toward unity (Fig. 2.17d). In Mode II, a plastic zone develops around the opening and is surrounded by the elastic ground, the elastic zone area is small, and most of the straining will occur within the plastic zone. Thus, the settlement profile above the crown looks like the one plotted in Fig. 2.17b and the settlement ratio  $S_{2}/S_{3}$  must be much less than unity (Fig. 2.17d).

The shape of the surface settlement trough is defined by S and i. Eqn. 2.10 indicates a strong dependency between i and p (or S') which is shown diagrammatically in Fig. 2.17c. On this figure are also shown schematically two pairs of GCCs for a shallow and a deep tunnel, and S /a-i/a curves of Modes I and II for each tunnel. For the shallow tunnel the yield zones (Mode I-2 or II-1) will propagate to the surface with a comparatively small displacement. For Mode

I-2, i will displace toward the centre of the trough in a faster rate. The width of the settlement trough for Mode I must be related to the inclined angle β of the shear plane to horizontal (Fig. 2.10) and will be narrow for the shallow tunnel. For Mode II-1 the settlement trough will be wider because the strength mobilization and stress redistribution are distributed around the opening, instead of concentrated at localized zones. The inflection point will remain fairly stationary about at the initial state (i.e., no ear-shaped yield zone developed (Mode II)), and will decrease with progressive stress relief near collapse. Same reasoning is applicable to the deep tunnel. When the extent of yield zone around the tunnel is small, i for both modes may remain stationary. As the yield zones reach the surface, the i will drop significantly.

From Fig. 2.17d, the  $S_S$  ratio is fairly constant for Mode I, and seems to be almost independent of  $p_i$ . This is in contrast to Mode II where the  $S_S$  ratio is initially small and later increases with decrease in  $p_i$ . As the collapse mechanisms are approached the settlement ratio will tend toward unity for both cases, but at a faster rate for Mode I.

From the above analysis, if the same displacements are allowed in tunnel construction larger surface settlement and a narrower trough will be induced in Mode I than in Mode II. Ground with low K imposes more difficulties in ground controls than ground with K close to unity.

### 2.2.5 Numerical Examples (FEM)

### 2.2.5.1 Approach

Numerical examples generated by the finite element method are used to understand the tunnel behaviour and to investigate the applicability and validity of the proposed technique described in the previous section. Eight analyses were performed and the input data for each case are listed in Table 2.1. The behaviour of purely cohesionless and cohesive soils were investigated. For cohesionless soil, one set of typical soil parameters and a range of K -values from 0.5 to 1.3 were chosen to observe the effect of  $K_{\perp}$  on the yield zone localization. Behaviour near shallow and deep tunnels was studied by varying the free surface boundary parameter, H/a. It is important to realize that the cohesionless model with the assumed associated flow rule predicts larger dilation and displacement than the real soil model of non-associated flow rule. For the cohesive soil, only shallow tunnel behaviour was considered. For simplicity, the unconfined compression strength of the ground was assumed to be constant with depth and thus the K -value had to be set to unity. It was assumed that the excavation process could be simulated by proportional unloading of the original insitu stress (denoted by P\* in Table 2.1). A case of an air-pressurized tunnel (Case AP1, Table 2.1) was studied by applying a uniform internal pressure to the periphery of the tunnel opening.

The finite element program SAFE, developed by Chan (1985), was used for all analyses. Fig. B.3 shows the finite

element mesh for H/a=18, composing of 8-node quadrilateral iso-parametric elements outside the tunnel and 6-node triangular iso-parametric elements inside the tunnel. Zero lateral displacements are imposed along the two sides and zero vertical displacements at the bottom boundary. For H/a=5, the same mesh was used, but elements at the four top rows were deleted. Each analysis involved the following steps:

- 1. Apply gravity stress by switch-on-gravity method, i.e.,  $\sigma_{y} = \gamma y, \ \sigma_{b} = \nu/(1-\nu)\gamma y,$
- 2. Amend the Poisson's ratio,  $\nu$  to a constant value.
- 3. Remove elements inside the tunnel and reinstate the initial insitu stress by applying equivalent nodel forces so that no displacemnts are induced.
- 4. Simulate the excavation process by reducing the nodal forces around the periphery of the tunnel in increments.
- 5. Continue step (4) until no numerical convergence in equilibrium iteration is possible for stress relief increments of 0.1% of insitu stress.

# 2.2.5.2 Results

#### Modes of Yielding

The extent of the yield zones at intermediate and final stress levels are plotted in Figs. 2.18 to 2.22 for eight cases. Table 2.2 summarizes the mode of initial yielding, the mode of yielding propagation, the applied stress at the yield initiation and at the last step (expressed in % of insitu stress). Some of these results are plotted in Figs.

2.7 to 2.8 for comparison, and they agree well with the predictions. Several important aspects on mode of yielding near tunnel are confirmed by these results of the finite element analyses:

- (a) Cohesionless Soil (Figs. 2.18 to 2.21)
- 1. For K <1.0, yield initiation occurs at the springline.

  Depending on K, these localized yield zones at the springline develop into (1) localized lobes at both shoulders (K =0.5) or (2) continuous plastic domain around the opening (K =0.82). Fig. 2.18a (H/a=5.0 and K =0.5) shows the propagation of localized zones at the shoulder toward the free surface. On Fig. 2.18b

  (H/a=5.0, K =0.82), it can be seen than two ear-shaped lobes start to emerge from the continuous plastic domain, indicating the development of yield zone localization (Mode II-1). For K =1.3, yield zones are initiated at the roof and floor (Mode III) and merge to form a continuous zone.

Yielding is observed near the ground surface. This is explained by the fact that Mohr circles of elements near the surface are ill-defined because of small overburden pressure. Small stress changes could induce yielding. This problem is one of numerical, inadequacy of finite element analysis.

2. The influence of the free surface boundary on the yield zone can be observed by comparison of results of analyses of shallow and deep tunnels. For deep tunnels, the shape of plastic zones is symmetrical about the horizontal tunnel axes. In shallow tunnels, the free surface has a tendency to displace downward, which accelerates the yielding process at the roof. Thus at equal percentage of stress-relief, the extent of the plastic zone at the roof is greater in shallow than in deep tunnels.

3.

From Fig. 2.20 (Case AP1), it can be observed that the boundary stress around the periphery of the tunnel has a significant influence on the mode of yielding.

Initially, application of a uniform pressure equivalent to overburden pressure (at the level of tunnel axis) induces yielding at the roof and floor. At 60kPa yielding occurs only in the roof and floor, and after a further pressure reduction to 20kPa (20% of overburden) the mode of yielding is similar to Case ST1. However, the extent of the yield zone at comparable stress level is much reduced. This finding is of great practical signifance. Most excavation modelers simulating pressures neglect this important aspect. It is also extremely important for back analysis or interpretation of field measurements.

mechanisms are dominant in the six FEM cases. Thus the minimum support stresses are the p (upper bound solution). The required p for roof collapse are calculated: 8% and 3% of overburden pressure at tunnel axis for H/a= 5 and 18 respectively. Comparisons between the applied stress at the last step (Table 2.1) and

p imply that the tunnels are still far away from the ic kinematic collapse state. However, severe problems have been encountered at this stage in modelling, indicating the limitation of the FEM in exploring the limit state of this particular geomechanical system.

- b) Cohesive Soil (Fig. 2.22)
- 1. For both cases (STC1 and STC2), under K = 1.0 localized yielding initiates at the floor and develops into a continuous zone around the tunnel propagating downward rather than toward to the free surface. This response is a result of the method by which the stress is relieved at the periphery of the tunnel.
- 2. The tunnel in ground with higher strength (Case STC 2) remains stable without any support. For Case ST1 the yield zone reaches the surface at 40% support stress. Fig. 2.22 shows that the configuration of the yield zone is influenced by the free surface, i.e., the shape of the yield zone is no longer circular. Various kinematically possible collapse mechanism could occur within the extensive yield zone. Davis et al. (1980) studied such possible collapse mechanisms of shallow tunnels in cohesive material and derived solutions of collapse loads using the upper bound theorem of plasticity. The collapse load for Case STC1 was calculated: 38% of insitu stress which is close to that from the FEM.

Stress Distribution (Figs. 2.23 and 2.24)

Fig. 2.23 shows the tangential and radial stress distributions above the roof for a shallow tunnel (H/a=5) under K =0.5 and 0.82 (i.e., cohesionless cases ST1 and ST2). For both cases, an increase in tangential stress is observed, denoting arching developed at the roof due to stress-relief inside the tunnel. The equivalent coefficients K<sub>r</sub>, defined in Section 2.2.4 were calculated: K<sub>r</sub>=0.81 for K<sub>r</sub>=0.5 at 65% stress-relief, and K<sub>r</sub>=1.31 for K<sub>r</sub>=0.82 at 75% stress-relief. These K<sub>r</sub> values compare well with those given by Eqn. B.8. implying that the earlier assumption of equal amount of stress redistributed to the roof and floor is valid.

For K =0.5, the mobilized tangential-radial stress difference is much less than the soil capacity and the soil above the crown remains elastic. However, for K =0.82, yielding has occurred and the tangential stress decreases to cope with the decrease in radial stress in order to satisfy the failure criterion.

Similar reasonings are applicable to cohesive soils, Cases STC1 and STC2 (Fig. 2.24). The K values calculated are: 1.38 for STC1 (60% stress-relief) and 1.62 for STC2 (unsupported), and they compare well with the predicted values given by Eqn. B.8. The tangential-radial stress difference within the plastic zone is governed by the unconfined compressive strength of the soil which is independent of overburden pressure.

It is interesting to observe the horizontal to vertical stress ratio K along the depth above the roof. For K = 0.5 (at 65% stress-relief) the K-value is close to unity above 3.0m depth and increases to 2.0 at the crown. For K = 0.82 (at 75% stress-relief) the K values from 1.5 to 3.0.

Terzaghi (1943) derived a solution for the vertical stress distribution above a tunnel in cohesionless soil using the soil arching theory, and recommended that the K-value should be approximately unity. Hence, Terzaghi's case follows closely with Case ST1 (K =0.5). However, the K-value also depends on the amount of stress-relief or the displacement. Thus Terzaghi's case is only valid for a special case of tunnel behaviour.

## Ground Responses (GCC)

Ground responses at the roof (Rf), springline (Sl) and floor (Fl), expressed in GCCs of normalized fictitious support stress versus normalized displacements, are shown in Figs. 2.25 to 2.28 for all eight cases.

For all cases except AP1, the ground initially responds to elastic unloading, and the displacement is proportional to the amount of stress-relief. For K =0.5 and 0.82, the amount of stress relief at the floor is greater than at the roof and springline, and thus the slope of the convergence curve is initially steeper for the floor. However, once yield zones are developed, the plastic deformation starts to dominate. Once yielding occurs, the support stress at the roof will remain fairly constant with a further increase in displacements. Similar reasoning can be applied to the cases

with  $K_2 = 0.82$  and 1.3.

The ground responses for Case AP1 (uniform internal pressure) are very different. The wall at the springline first displaces outward and then starts to converge as the uniform pressure decreases. The responses at the roof, springline and floor are almost linear even though plastic zones are formed first in the roof and then near the springline and the shoulders. This shows clearly that the approach of application of a constant fictitious p in simulation of excavation does not lead to a reasonable solution of ground response for K different from unity.

For case ST2 where K is close to unity, Eqns. 2.6 to 2.8 can be used to approximate the p /p -u/a relationship (Mode II). The results are plotted on Fig. 2.25 for comparison. These predicted values fall reasonably close to those of the FE analyses. The discrepancies can be attributed to the fact that Eqns. 2.6 to 2.8 assume Ko=1.0 and different stress boundary condition. A similar comparison for a deep tunnel (H/a=18, Fig. 2.27) shows that the gravity effect becomes less dominant, i.e., the three predicted GCC (roof, springline and floor) are close together.

The FEM only provides the results for the initial portion of the GCC up to the applied stress levels stated in Table 2.2. The subsequent portion of the GCG, where localization processes and transition to collapse modes dominate, cannot be currently predicted analytically. However, it is possible to estimate the required support stress for several states

 $(p_{ys}, p_{ic})$  and  $p_{fc}$ ) which are of practical interest. These support pressures  $p_{ys}$ ,  $p_{ic}$  and  $p_{fc}$  were calculated for each case and are listed in Table 2.2.

For the two cases with cohesive soil (Fig. 2.28), Egns. 2.6 to 2.8 (with cohesive strength criteria) are used to calculate the predicted GCC. The predicted response corresponds well with those calculated by FEM for the stable case of STC2. It is interesting to note that for this case the GCC of the floor are higher than those of the springline and the roof. With proportional unloading greater deviatoric stresses are induced at the floor than at the roof and springline. This causes a deeper plastic zone and larger deformation at the floor because the strength of the cohesive soil is independent of confining pressure. From a comparison with Case STC1, agreement is only observed initially up to u/a<0.4% because the plastic zone influenced by the free surface is no longer circular. The GCC for the roof becomes horizontal, indicating that the collapse mechanism is approached. Eqns. 2.6 to 2.8. do no longer apply. The calculated collapse load of 38% is , however, close to the applied stress at the last step (40%). Despite these discrepancies it can be concluded that the models proposed earlier can be applied to predict the ground behaviour with reasonable accuracy.

## Ground Deformation near the Tunnel

Displacement vectors around the tunnel are plotted in Figs. B.4 to B.8 for all eight cases (Appendix B). The

directions of these vectors are influenced by the insitu stress state K and the proximity of the free surface boundary H/a. For low K values, radial displacements toward the openings are clearly restricted to areas above the roof, and below the floor. Near the springline, soil elements immediately adjacent to the opening exhibits inwards movements, but elements at about one radius away from the wall are not influenced much by the stress relief. For K close to unity, radial displacements are distributed evenly around the tunnel to a distance of about four radii from the wall. This implies that this can be approximated by the "hole-in-plate" model.

One feature of practical importance commom to Figs. 2.25 to 2.28 and B.4 to B.8 is the overall configuration of wall convergence. For K < 1.0, initial unloading induces larger displacements at the roof and floor than at the springline. But after yielding has initiated, the plastic deformation at the springline becomes dominant and the displacements at the springline are larger than those at the roof and springline. If a liner is installed at this stage and deformed according to the configuration of wall convergence as those shown in Fig. B.4, tension will be induced at the inner face of the liner at the springline. This is in contrast to the results of conventional continuum design method, like the relative stiffness method. In these methods (Muir-Wood, 1975; Einstein and Schwartz, 1979), the liner will displace inward at the roof and floor and outward at the springline for (K <1.0). This causes compression instead of tension at the inner face of liner at the

springline. Hence, for the determinaton of thrust and bending moments in a liner, one has to consider not only the initial external stress distribution but also the configuration at the time of liner installation.

From Figs. B.4 to B.8, it is of interest to observe that Case AP1 has a unique displacement field around the tunnel. Even at large stress relief, the magnitude of displacement is much smaller than for the other cases. The technique of simulating stress relief not only affects the extent of the yield zone but also dominates the ground displacements.

## Settlement Profile above the Tunnel

Subsurface settlement profiles above the crown for various stress-relief levels are plotted in Figs. 2.29 to 2.32 for all eight cases and should be compared chosely with the schematic Fig. 2.17.

For shallow tunnels in cohesionless soil (Cases ST1 and ST2), two distinct profiles can be identified. For Mode I (ST1) where localized yield zones form and the soil block above the roof does not yield, the differential settlement with depth is only due to the elastic deformations, and is small. Thus, the vertical displacement increases only slightly and gently with the depth (Fig. 2.29a). In Mode II where the opening is surrounded by a continuous yield zone, the settlement increases gently in the elastic zone and accelerates rapidly within the plastic zone (Figs. 2.29b and 2.30a). An abrupt change in settlement gradient can be observed at the elastic-plastic boundary (at

 $p_{i} = 25\%$ ).

For deep tunnels (Cases DT1 and DT2) the settlement profiles are similar because the extents of the yield zone for both cases ( $K_0=0.5$  and 0.82) are relatively small. However, at equal stress relief (say 40%) larger settlements are experienced in Case DT1 ( $K_0=0.5$ ).

The behaviour is less well defined for the cohesive soils (Fig. 2.32) because the numerical model assumed that no volume change is associated with plastic deformation. However, it is interesting to observe (Fig. 2.32a) that a decrease of support stress from 50% to 40% induces a sudden, almost constant, increase in vertical displacement indicating initiation of a collapse mechanism ("plug" failure).

The relationship between surface settlement (S<sub>S</sub>) and crown settlements (S<sub>C</sub>) at intermediate stress-relief levels are plotted in Fig. 2.33 for all cases. Fig. 2.33 confirms the earlier discussion on the slopes of S<sub>S</sub>/S<sub>C</sub> for Modes I and II. It is greater for Mode I than for Mode II. Fig. 2.34 shows the relationship between the settlement ratio S<sub>S</sub>/S<sub>C</sub> and the stress ratio p<sub>1</sub>/p<sub>0</sub>. From Fig. 2.34, it can be seen that S<sub>S</sub>/S<sub>C</sub> is a function of p<sub>1</sub>/p<sub>0</sub>, especially for cohesive soil of low strength. However, for the range of support pressure typically encountered in the field, the S<sub>S</sub>/S<sub>C</sub> can be reasonably assumed to be constant with p<sub>1</sub>/p<sub>0</sub>. Atkinson and Potts (1977) observed in model tests that the S<sub>S</sub>/S<sub>C</sub> is independent of p<sub>1</sub>/p<sub>0</sub> and only a function of soil properties.

In tunnel design, one is concerned with not only the maximum surface settlement but also the differential settlement, i.e., the gradient at the inflection point and the trough width. Figs. 2.35 to 2.37 show the surface settlements profiles for all cases. At a given stress-relief level, Case ST1 is the most critical with the largest total settlement and the narrowest trough. The large settlement is due to a low K, i.e., low tangential arching to resist downard movement. The narrow trough can be attributed to the formation of two near vertical yield zones which restrict the downward movement of the soil block above the roof. As K increases, the tangential arching above the roof increases. Also, the development of localized yield zone is suppressed and Mode II takes places resulting smaller settlement above the tunnel and wider trough. For K = 1.3 (Case ST3), the surface settlement is negligibly small because of high arching. For deep tunnels, the free surface boundary is less influenced by the tunnel, i.e., the trough is wider and gradient is gentler (Fig. 2.36).

settlement profile of Case AP1 (Fig. 2.35), which is similar to a case of an air-pressurized tunnel. The effect of uniform pressure inhibits the propagation of the localized yield lobes to the surface and thus reduces the surface settlement significantly even at relatively low support pressure levels. This demonstrates that application of air pressure (constant p<sub>1</sub>) is very beneficial for the two purposes of controlling the displacement distribution and the extent of the yield zone. Excellent field measurements

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that support these conclusions were presented by Kovari et al., 1979.

2				ч	>	>		•	-re1181
2		φ (deg	ф q <sub>u</sub> (deg.) (кра)	(MPa)	(MPa) (kN/m3)	•			•
Sff1	Perfectly Elastic- Plastic, Associated, Flow Rule,	30	0 40 20 (constant with depth)	40 11 w 11	20 1 depth)	4.0	O .5	6. 0	ه .
ST2	as above	90	0 40 20 (constant with depth)	4 t th	20 depth)	4.0	0.82	5.0	•
ST3	as above	30	0 40 20 (constant with depth)	t ¥1 th	20 depth)	O 4.	£. 3	5.0	å
AP 1	as above	30	Constant with depth)	. 40 t vith	20 depth)	0.4	S . O	رة .0	໋
DT 1	as above	30	, 0 40 20 (constant with depth)	¥ 40 th	20 depth)	0 4.	s. O	18.0	
012	ayoda sa	30	0 40 20 (constant with depth)	4 3 6 ±	20 depth)	0 4	0.82	18.0	<b>å</b> .
STC1	£ 5.	. •	30	0	50	0.495	0.	5.0	•
erc.	Von Mises	c	(constant with depth)	¥ . 6	depth)	, 0.4 0.4	و جانب	ر ت	٥
; ;		>	(constant with depth)	# *	depth)		<u> </u>	<b>&gt;</b>	

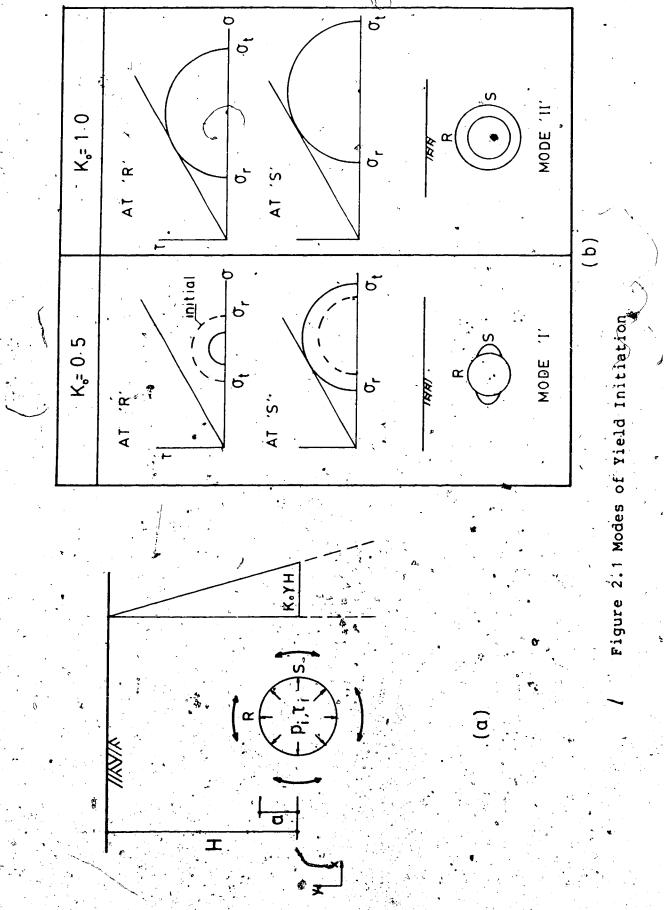
Ana lysts No.	н/а	, X	Initial Yielding at	Yielding Propagation	Applied Stress at last step	Type of Mode
STI	ហ	0.5	\$1 (81%)	Ø,	35%	1.2
ST2	ம	0.82	51 (56%)	0	75%	. 1-11/ 11
513	ស	e -	Rf/(62%)	<b>©</b>	26%	
AP 1	េ	6.0	Rf/F1(80kPa)	<b>X</b>	20k P.a	· <u>:</u>
DT.1	<del>6</del>	0.5	S1 (76%)	X	35%	1-1
DT2	8	0.82	\$1 (57%)		25%	1-11 / 11
STC1	ហ	0.	F1 (85%)	· (i)	40%	<b>:</b> .
STC2	ហ	0.	F] (66%)	0	%,	
Note: S1-spring Stress-rel	ng]ine relief		or, Rf-roof lapse load ex	pressed in % o	F1-f1oor, Rf-roof and collapse load expressed in % of insitu stress	

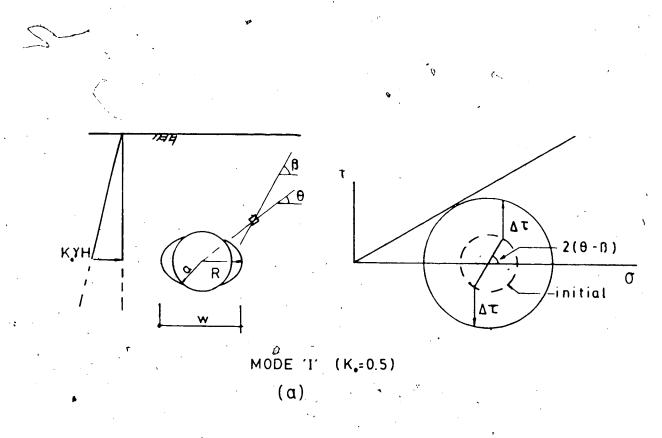
Table 2.2 Summary of FE Analyses (Tunnel)

Table 2.3 Summary of pys, pic, pfc (Tunnel)

Cases	Pys	Pic	_ Pfc-
ST1 (H/a=5, Ko=0.5)	24 %	8 %	67 %
ST2 (H/a=5, Ko=0.82)	-	8 %	61 %
ST3 (H/a=5, Ko=1.3)		8 %	-
7AP1 (*(H/a=5, Ko=0.5)	14kPa	8kPa	67kPa
DT1 (H/a=18, Ko=0.5)	-	3 %	67 %
DT2 (H/a=18, Ko=0.82)	- ;	3 %	41 %
STC1 (H/a=5, Ko=1.0, qu=30)	-	<b>-</b> ·	38 %
STC2 (H/a=5, Ko=1.0, qu=60)	· <del>-</del> ,	-	-

Support pressures at roof Pys, Pic and Pfc are expressed in % of overburden pressure to tunnel axis except for Case AP1





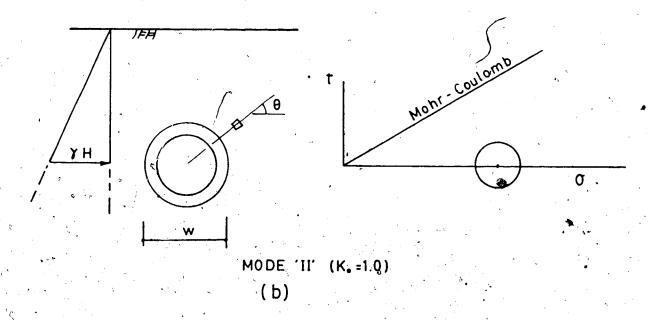


Figure 2.2 Propagation of Yield Zone

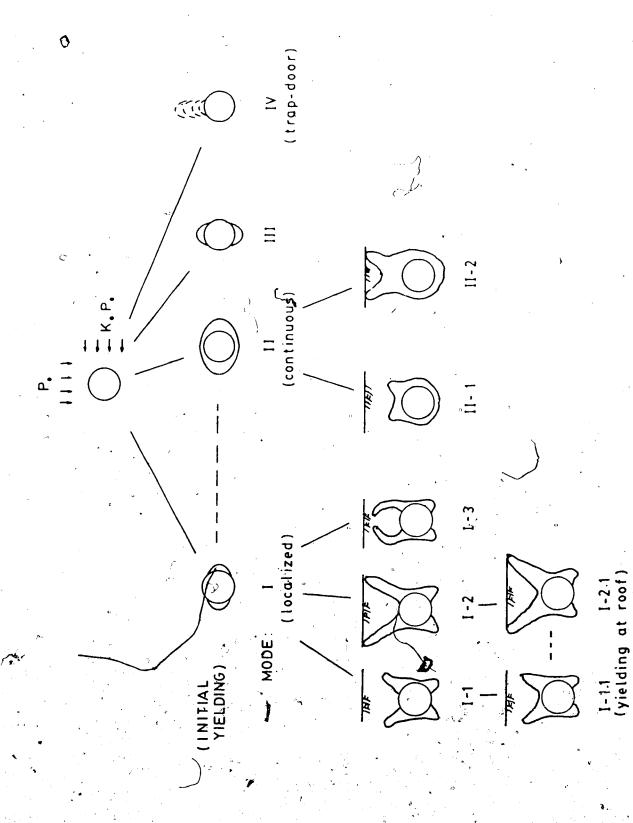


Figure 2.3 Possible Modes of Yielding near a Tunnel

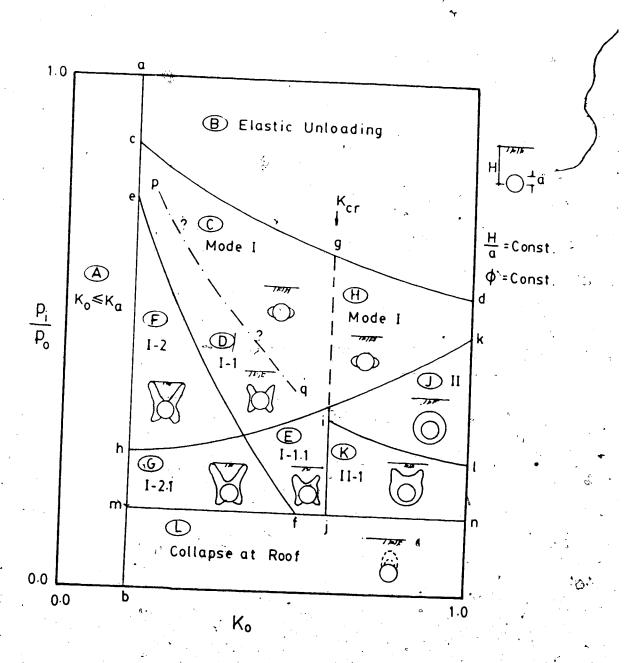


Figure 2.4 Modes of Yielding near a Tunnel

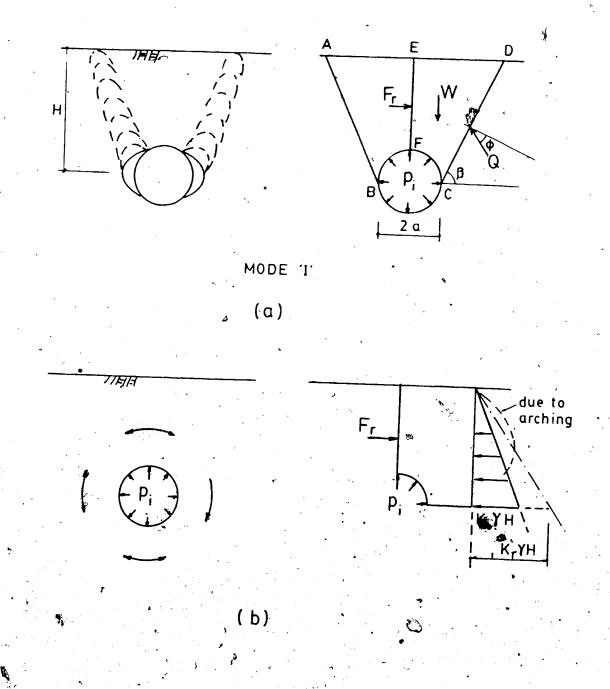


Figure 2.5 Mode I

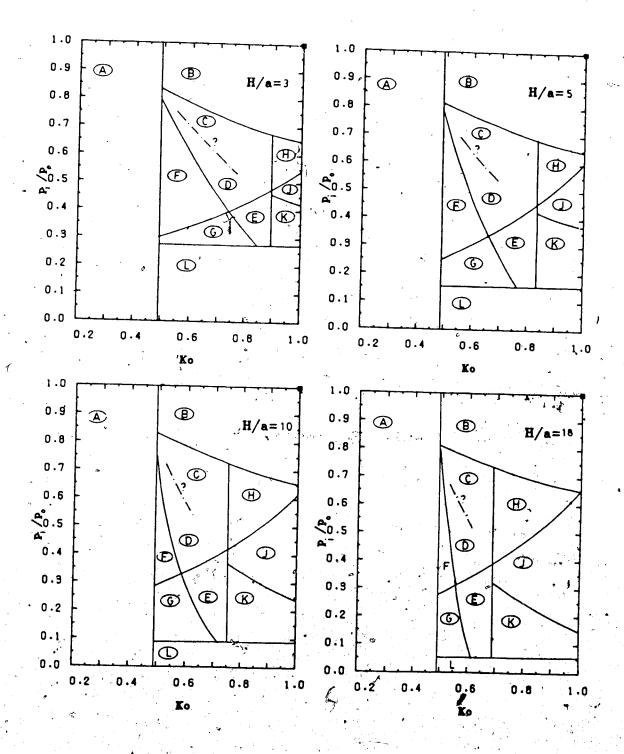


Figure 2.6 Modes of Yielding near a Tunnel  $\phi=20^{\circ}$ , c=0

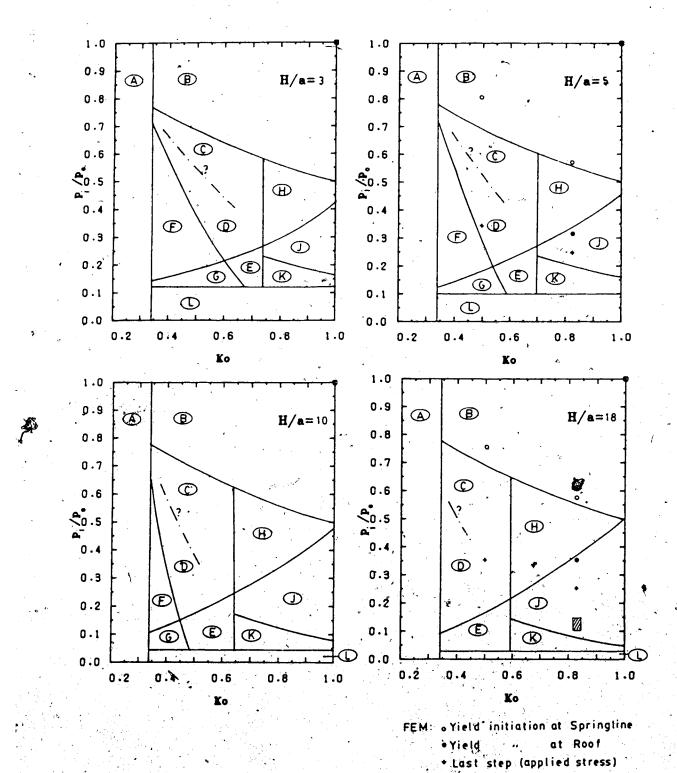


Figure 2.7 Modes of Yielding near a Tunnel  $\phi=30^{\circ}$ , c=0

Exp Tunnel:

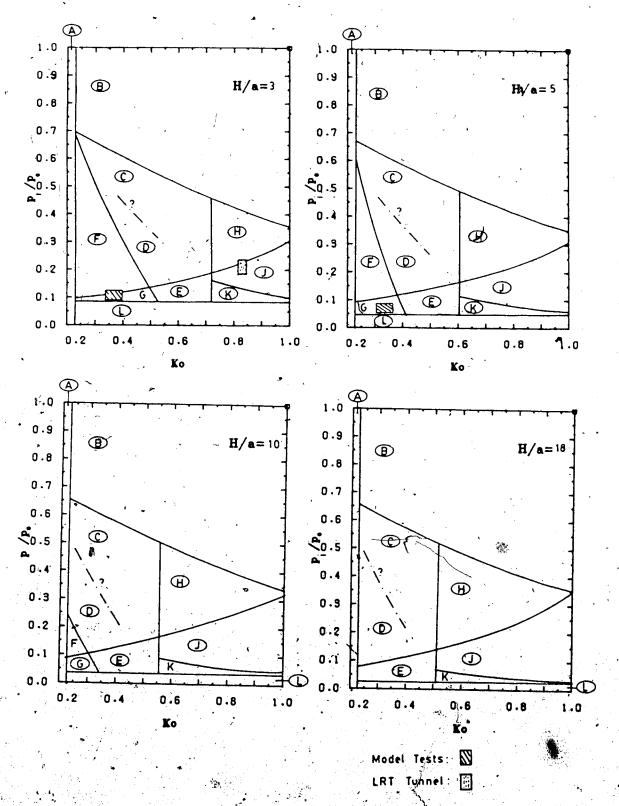


Figure 2.8 Modes of Yielding near a Tunnel \*=40°, c=0

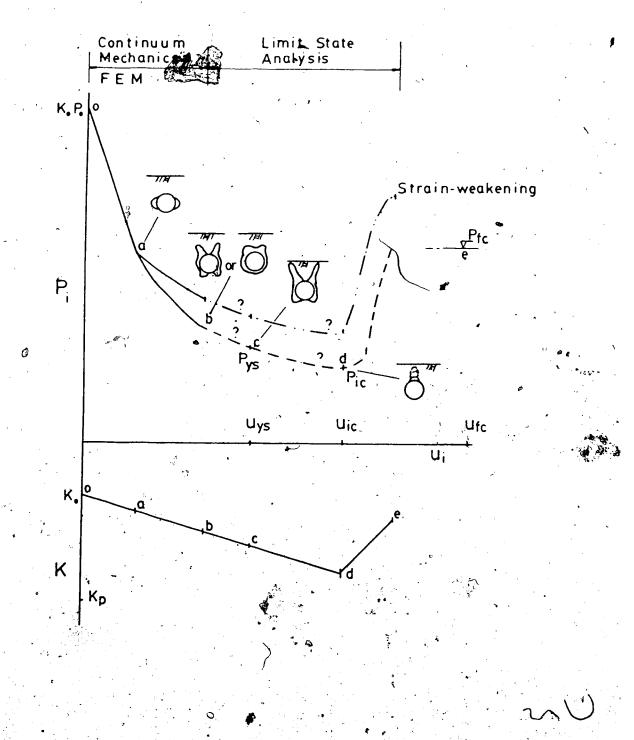


Figure 2.9 Formulation of GCC (roof) for Modes I and II

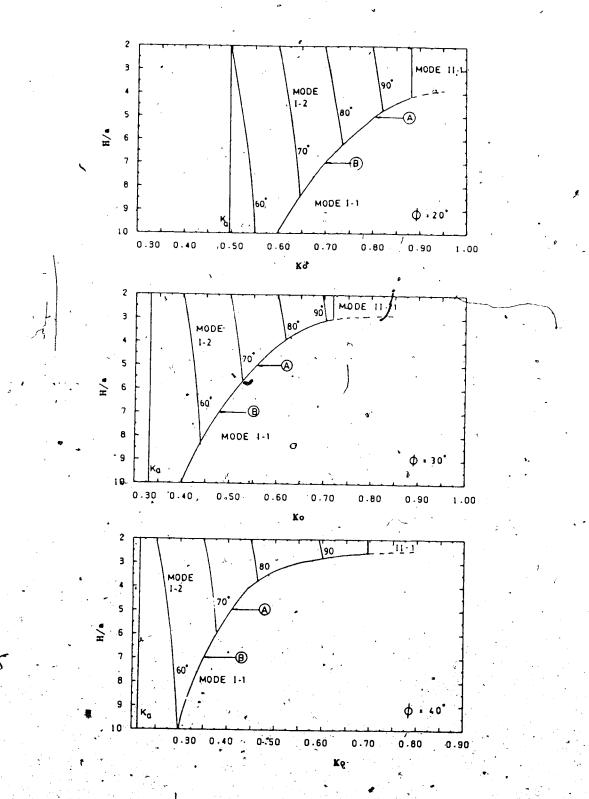


Figure 2.10 Direction of Propagation of Yield Zone (Mode I)

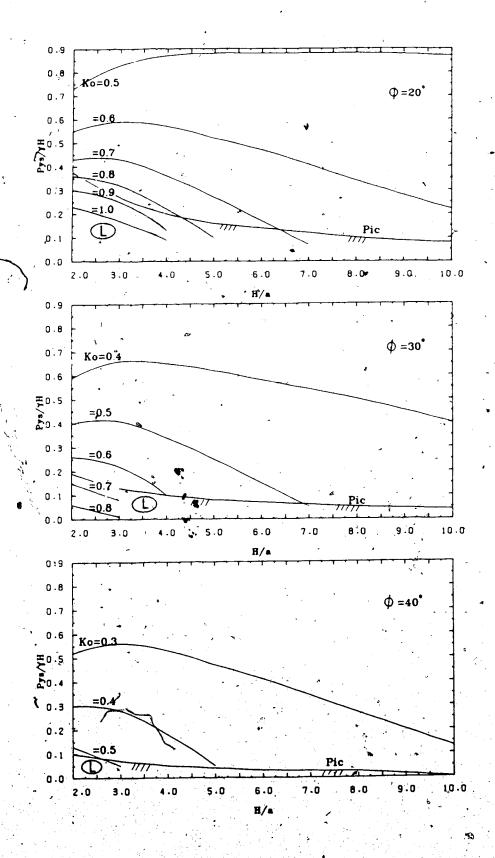


Figure 2.11 p (and p ) as function of H/a and  $\phi$ 

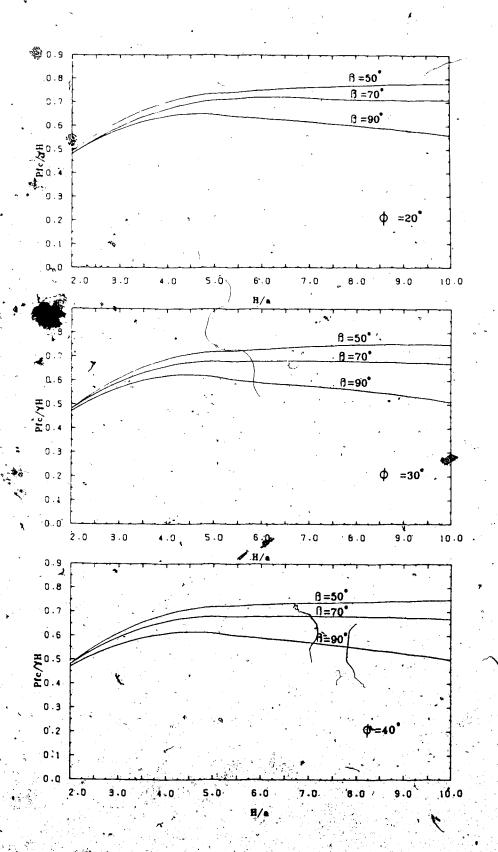


Figure 2212 p as function of H/a and  $\phi$ 

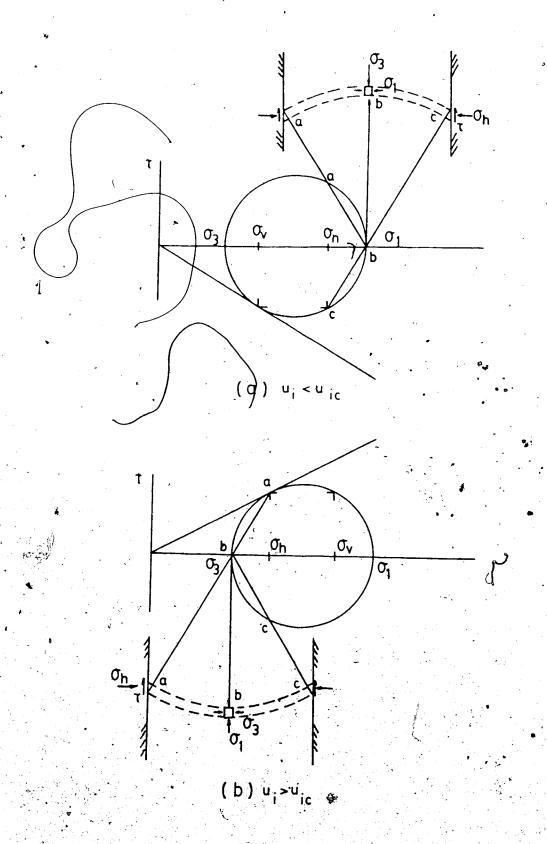
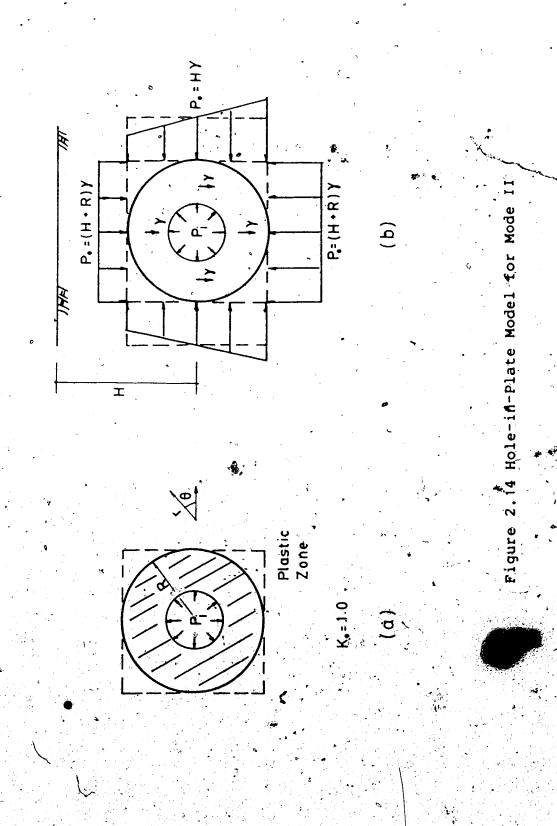


Figure 2.13 Types of Arching (a) Convex and (b) Inverted



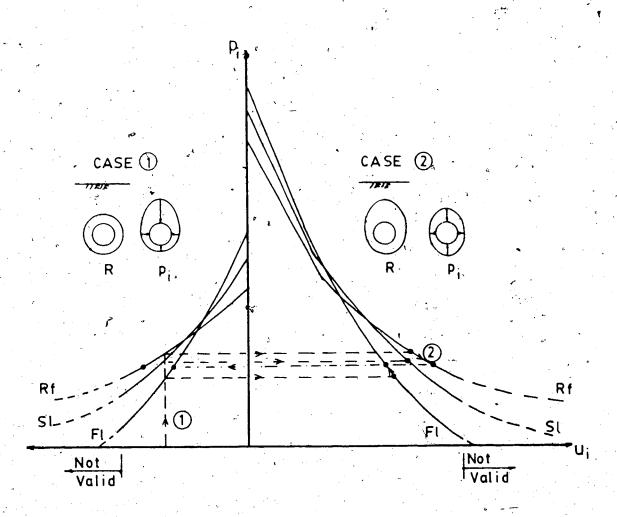


Figure 2.15 p -u-R Plot for Mode II

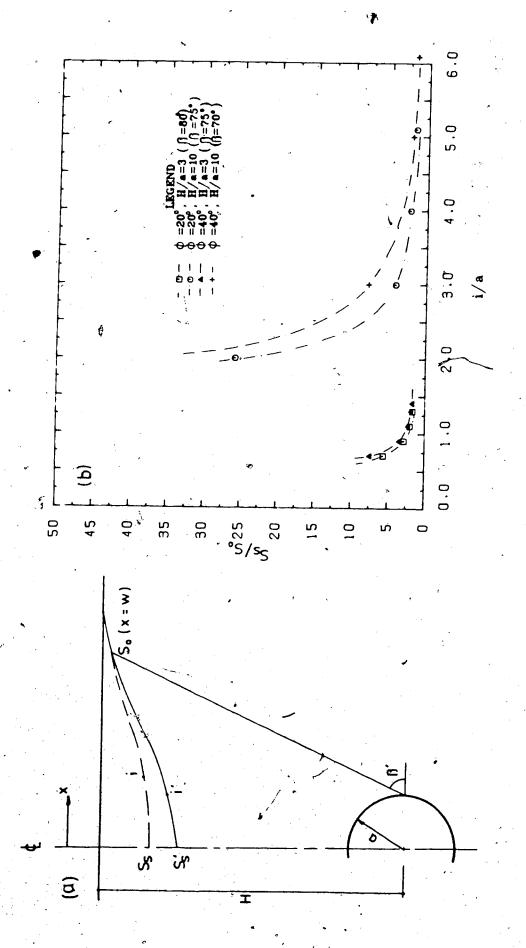


Figure 2.16 Surface Settlement Profile

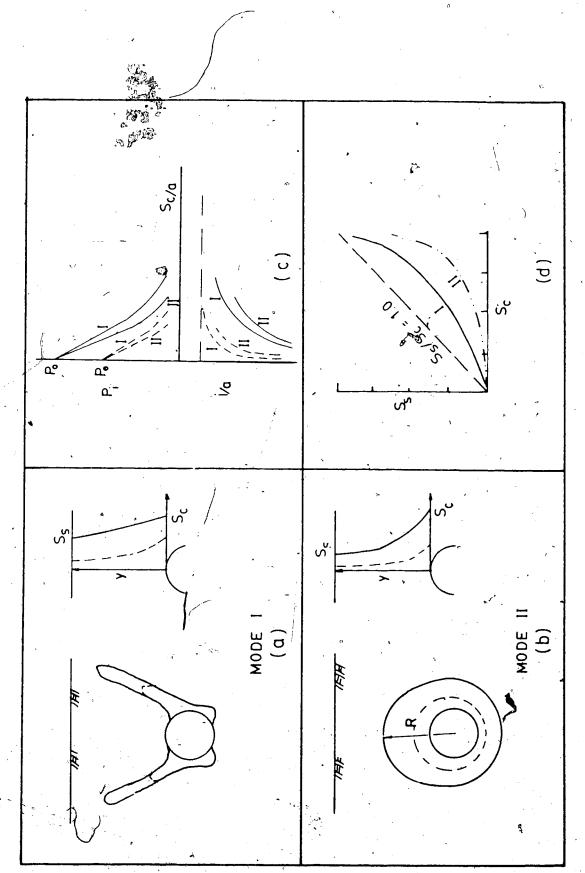
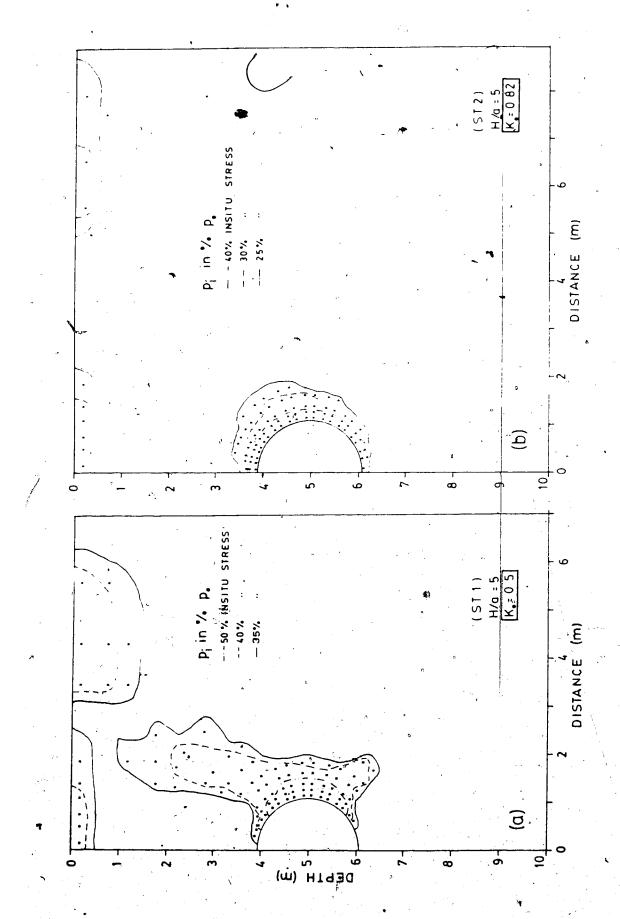


Figure 2.17 Vertical Settlement Profiles above Tunnel for

Modes I and II



- Cohesionless (ST1,ST2) Figure 2.18 Extent of Yield Zone

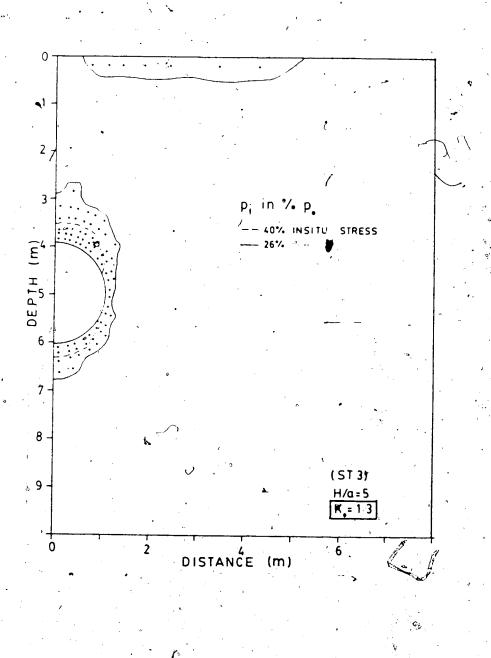
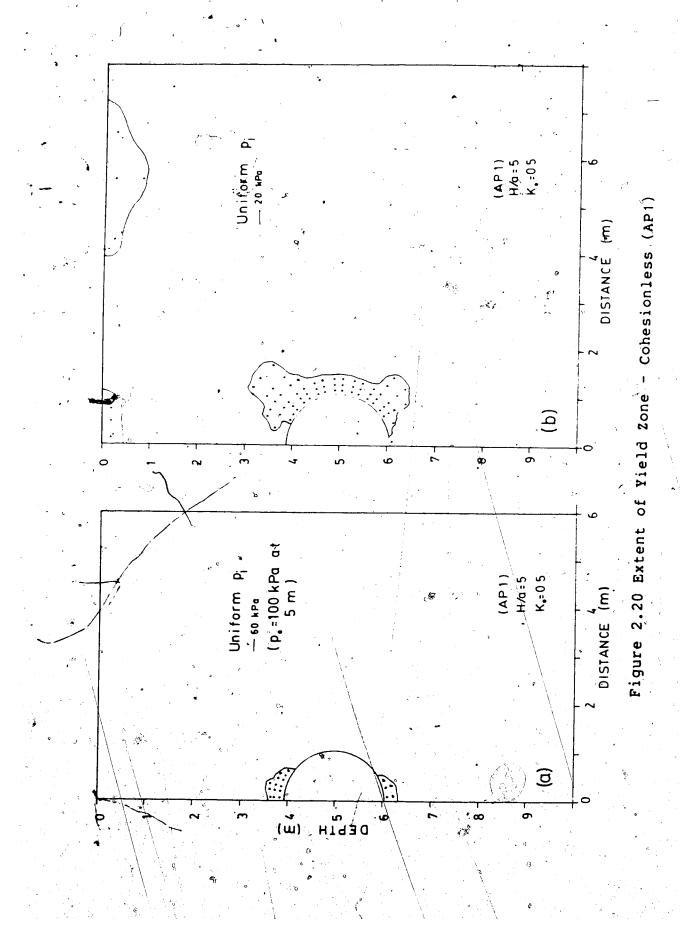


Figure 2.19 Extent of Yield Zone - Cohesionless (ST3)



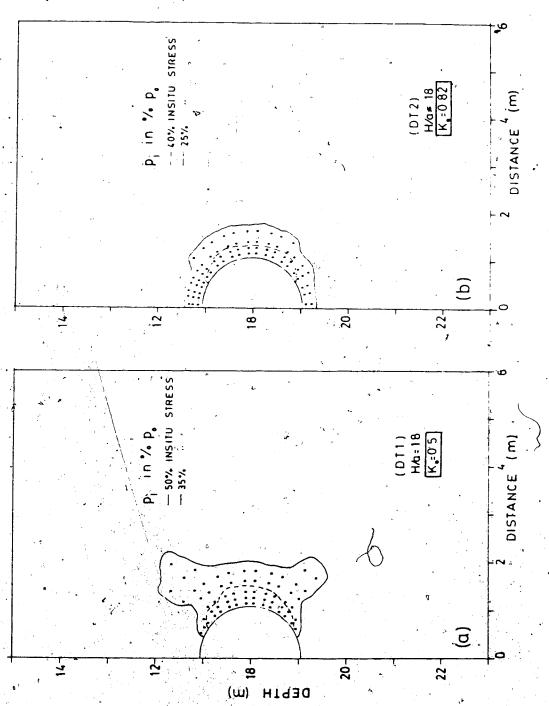
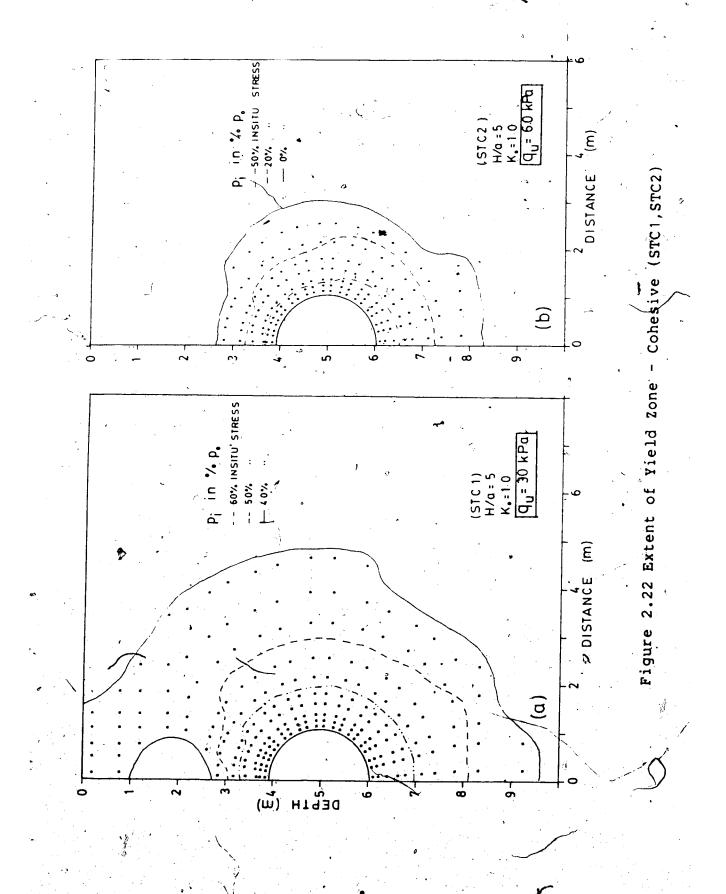


Figure 2.21 Extent of Yield Zone - Cohesionless (DT1, DT2)



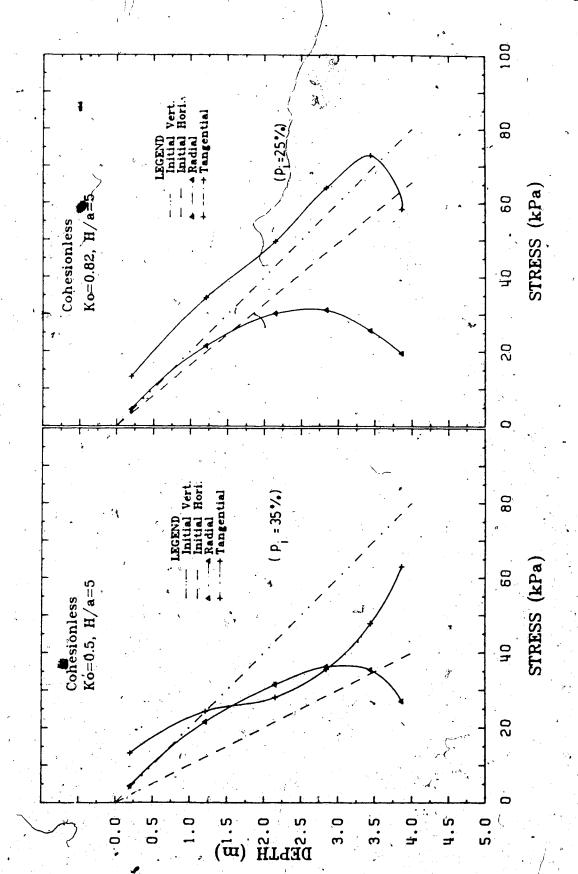


Figure 2.23 Stress Distribution above Roof - Cohesionless (ST1, ST2)

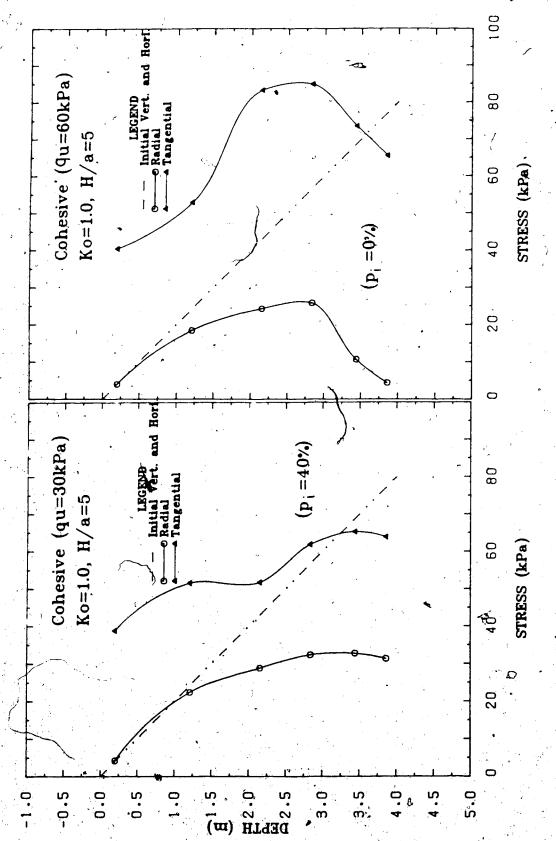


Figure 2.24 Stress Distribution above Roof - Cohesive (STC1, STC2)

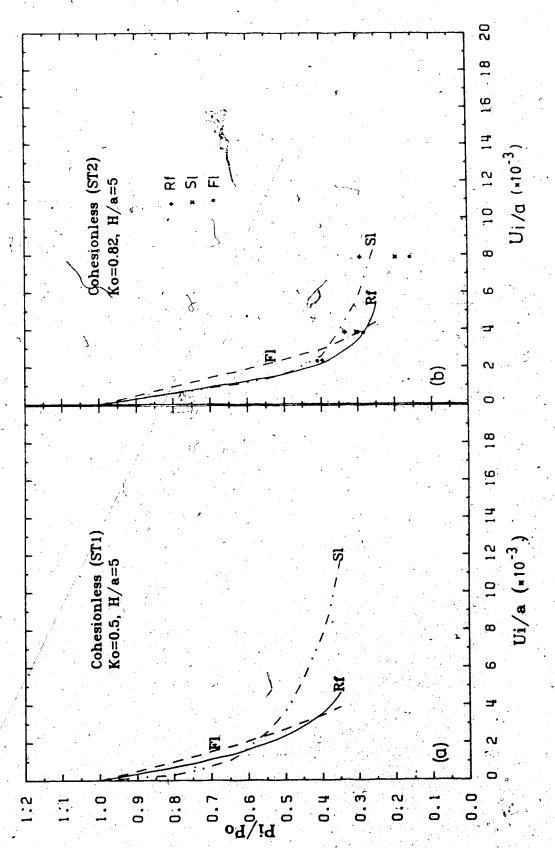


Figure 2.25 GCC - Cohesionless (ST1,ST2)



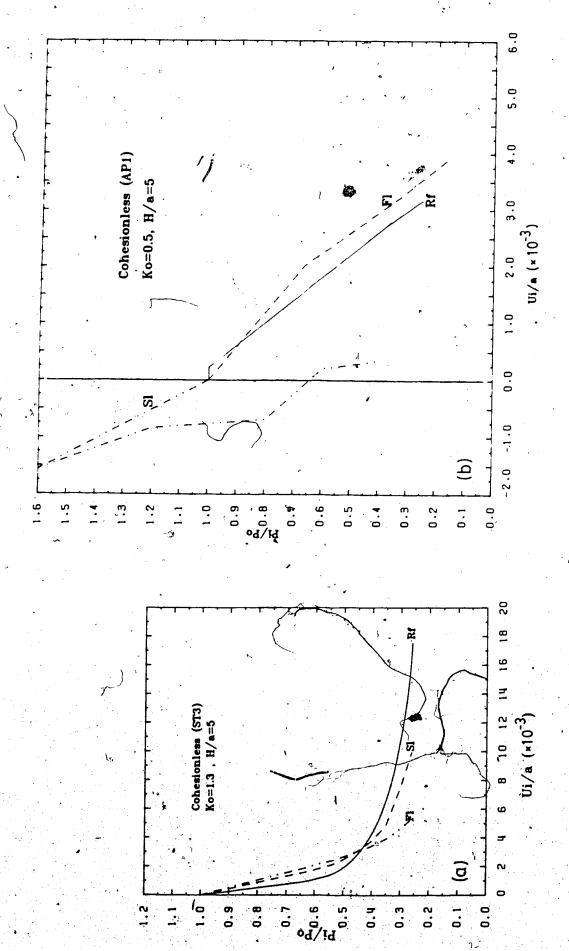


Figure 2.26 GCC Cohesionless (ST3, AP1)

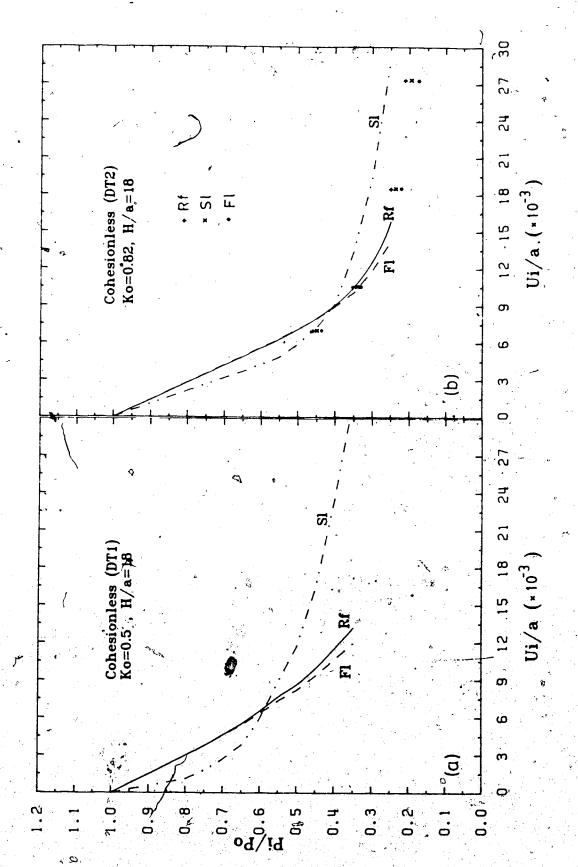
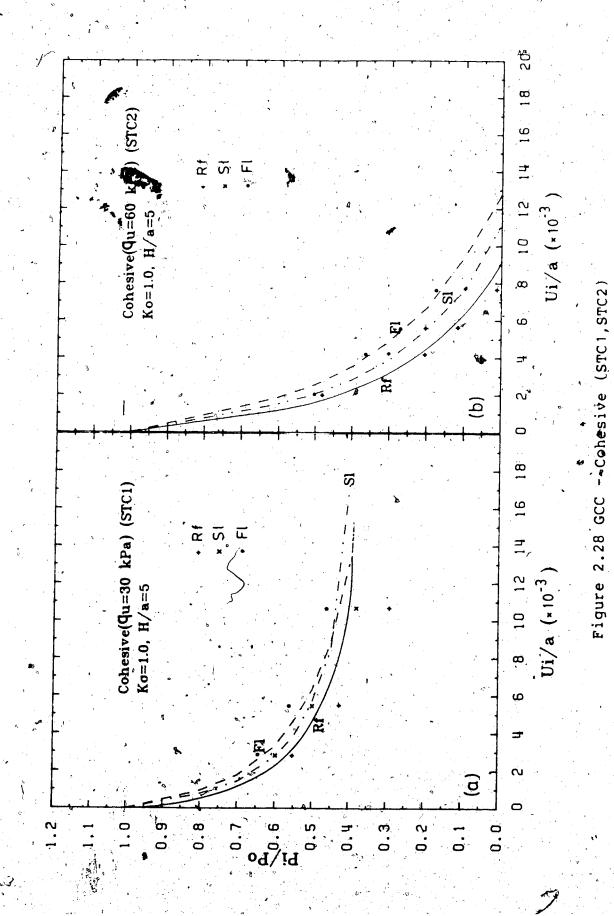


Figure 2.27 GCC - Cohesionless (DT1,DT2)



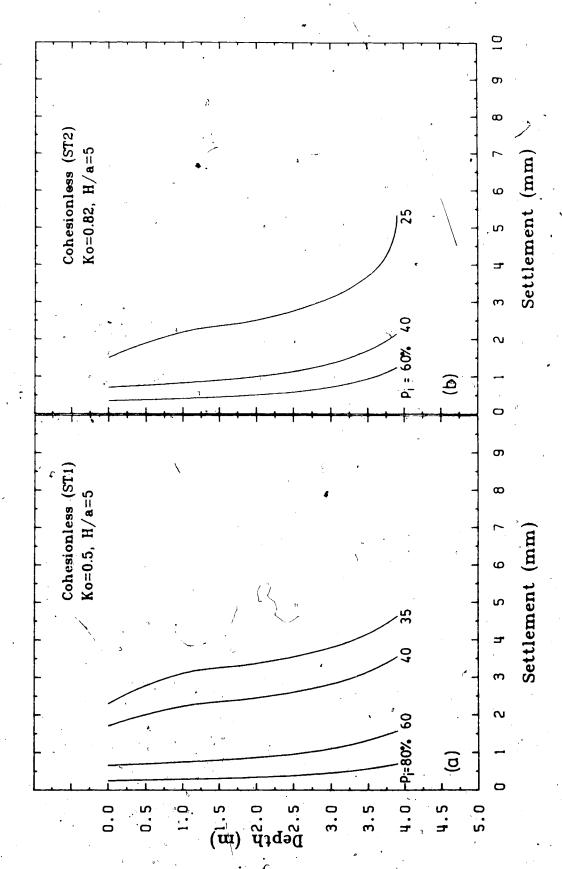
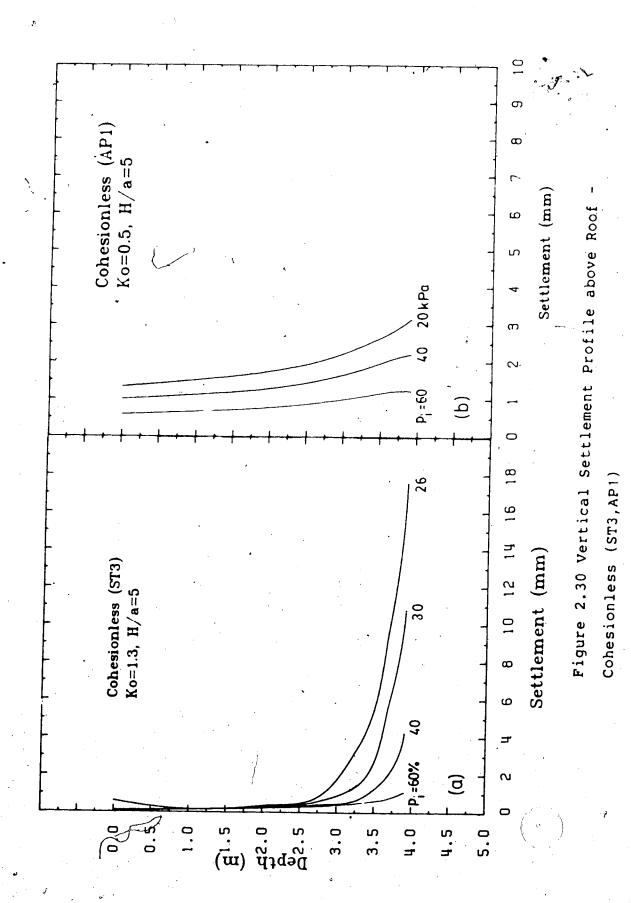


Figure 2.29 Vertical Settlement Profile above Roof - Cohesionless (ST1,ST2)



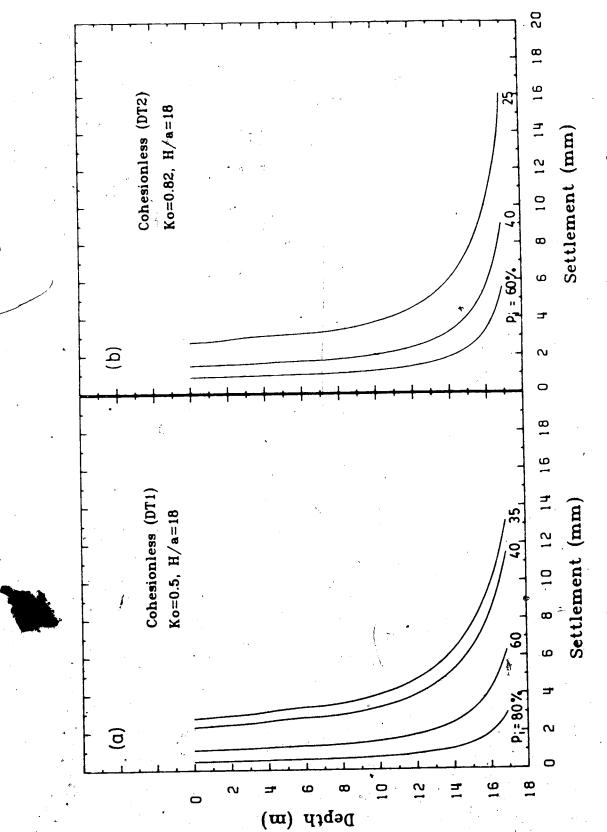
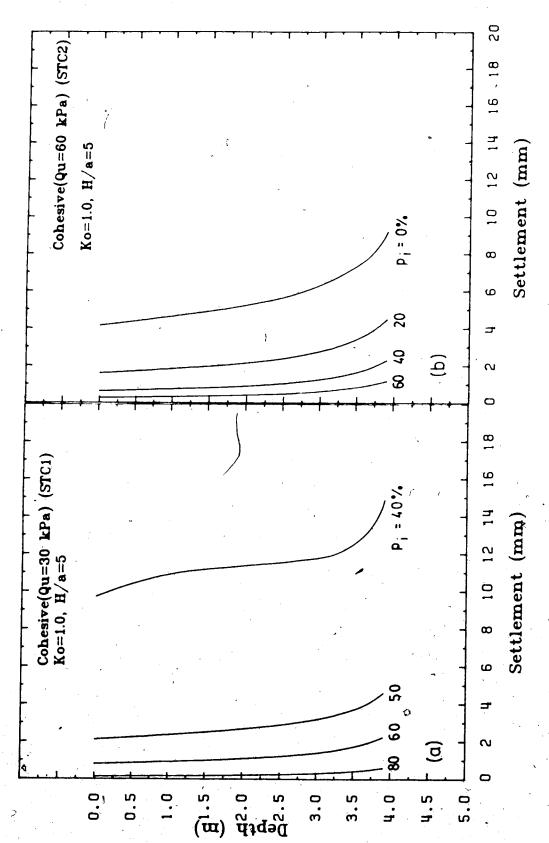


Figure 2.31 Vertical Settlement Profile above Roof - Cohesionless (DT1,DT2)

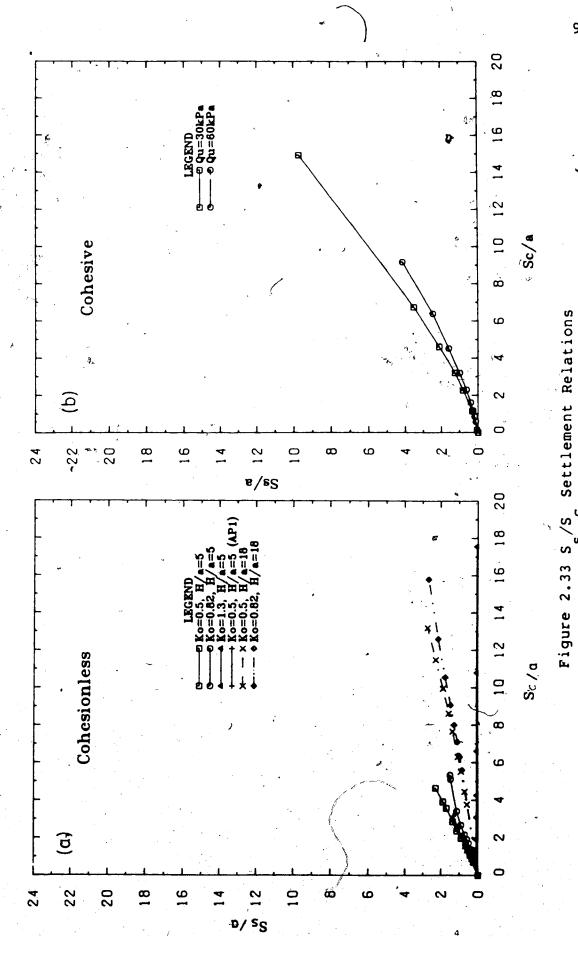
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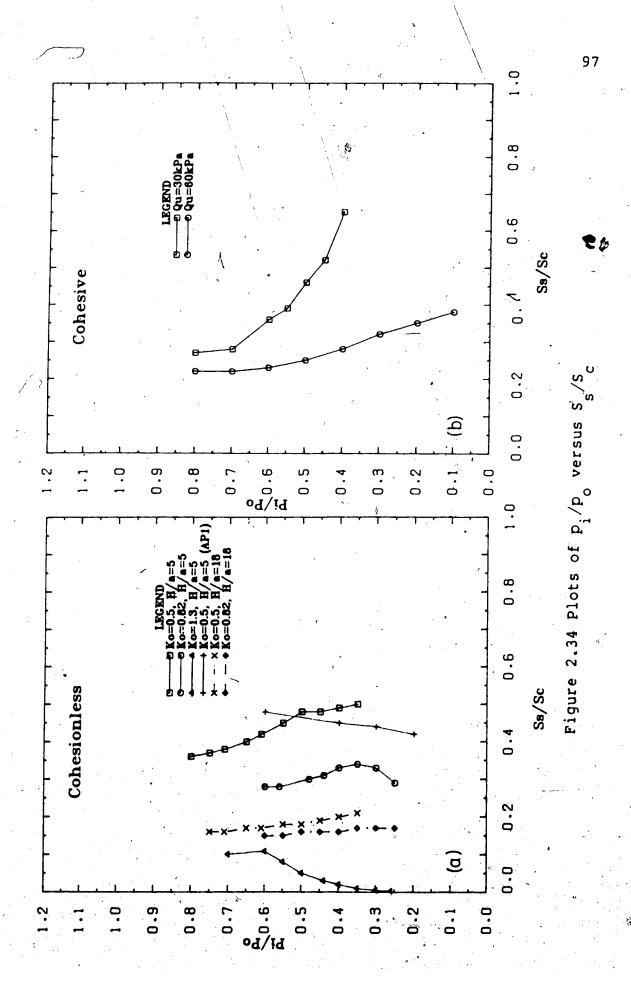


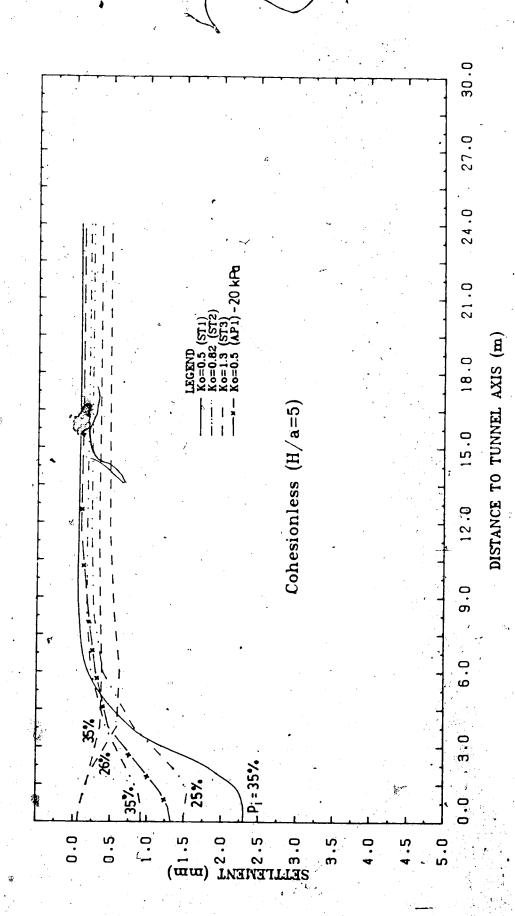
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Figure 2.32 Vertical Settlement Profile above Roof - Cohesive (STC1,STG2)

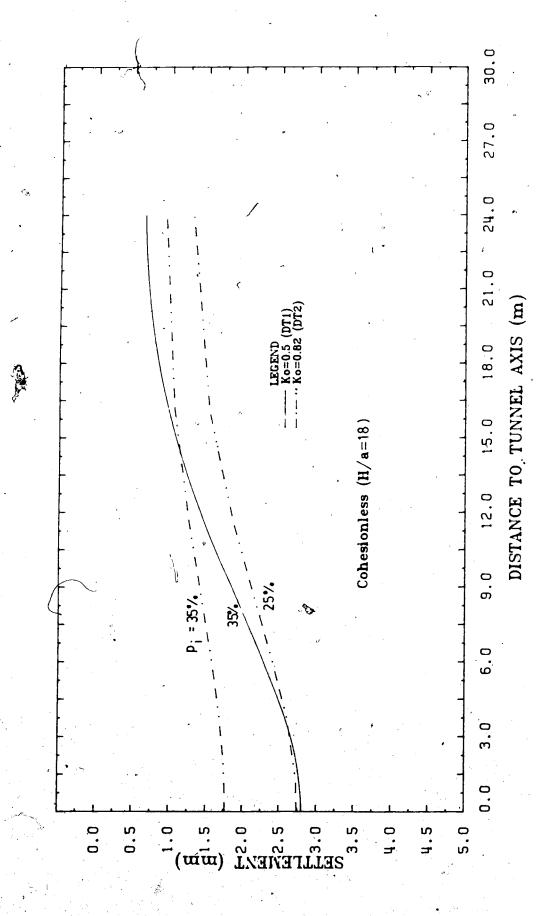
Settlement Relations



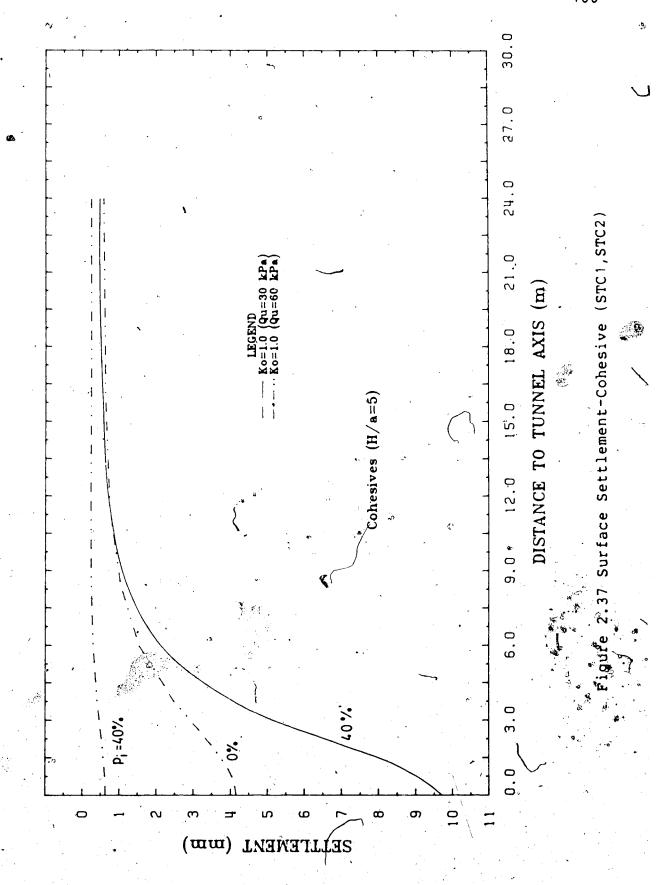




2.35 Surface Settlement-Cohesionless (ST1,ST2,AP1



 $^{*}$  Figure 2.36 Surface Settlement-Cohesionless (DT1,DT2)



#### 3 . TUNNEL - CASE STUDIES

#### 3.1 Introduction

The last chapter explores the mechanism of tunnel behaviour in detail and indicates that the behaviour of the ground around a tunnel is best described by the concept of ground convergence curve. This concept signifies the dependency between stress-relief and displacement of the ground. In this chapter an attempt is made to extend the application of this concept to explain results from some case histories, including model tests and field measurements.

#### 3.2 Case Studies

#### 3.2.1 Model Tests

The behaviour of shallow tunnels in sand and clay has been studied extensively by use of model tests at Cambridge University, England. Two types of tests were performed: (1) static tests and (2) gravity or centrifugal tests (Atkinson et al., 1974; Schofield, 1977). In static tests the soil sample was subjected to an external biaxial stress field in which the K -value could be varied, but the self-weight of the soil became negligible and did not contribute to instability modes. In the gravity or centrifugal test the stress field was due to the gravitational force and the K ovalue depended on the mobilized strength which could be approximated by the relationship (K =1-sin¢) (Brooker and Ireland, 1965; Lade et al., 1981).

Atkinson et al. (1975) performed a series of model tests (Type 2 - gravity test) in dense sand to measure the support pressure and deformation at the collapse state. The tunnel was supported by a pressurized flexible rubber membrane, and collapse was produced by reducing this internal pressure. The depth-ratio (H/a) was varied from 1.9 to 5.0. As described by the authors, the soil movements were initially restricted to a wedge zone immediately above the crown. The measured support pressure at this state was close to the pressure calculated from the upper bound solution (p., Eqn. 2.3). As the tunnel pressure was reduced further, this wedge grew upward toward the surface (Regime L in Fig. 2.4). A higher support pressure was measured to maintain the equilibrium at this final state. These observed modes correspond well with the prediction in Fig. 2.9 (Point d to Point e). The higher pressure at the final state should have corresponded to the pressure p. Back analyses of p are not possible because the horizontal to vertical stress coefficient (K<sub>e</sub>) and the residual angle ( $\phi_r$ ) were indeterminate from the tests. The support pressure p actually depends on the displacement allowed at the crown. The  $p_{fc}$  values shown in Fig. 2.13 were calculated assuming that the soil is in limit state (i.e., at lowest  $K_1 = K_1$ ). Thus the measured p in the model tests should be and are smaller than the predicted p fc.

Atkinson and Potts (1977) investigated theoretically and experimentally the stability of a circular tunnel in sand. The experimental investigation consisted of Types (1) and (2) model tests. In the gravity tests (with an assumed

 $K = 1 - \sin \phi$ ) the initial collapse was observed at the crown. The measured collapse pressures compare well with the pressure (p ) predicted by the theoretical upper bound theorem based on the observed collapse mode. For the static tests (where K conditions were not reported) the observed collapse pressures also bracketed closely the pressure range predicted using the lower bound theorem. However, their formulation (lower bound solution) assumes that a concentric yield zone develops around the opening and that the collapse state is reached when the yield zones intersect the free surface. It has been shown in Chapter 2 that this mode of . yielding, i.e., Mode II is only admissible if Ko is close to unity. If the model tests were performed under  $K = 1 - \sin \phi$ , the localized yield zones (Mode I) should develop, implying that the mode of yielding in the theoretical solution was not consistent with the one observed in the test.

Fig. 3.1 shows contours of volume and shear strains in sand near collapse for the gravity tests  $(K_0=1-\sin\phi)$ , (Atkinson and Potts, 1977). The distribution of strain around the tunnel is self-explanatory. Concentrations of shear strain with magnitude up to 20% occur between the crown and springlines, and the soil in these regions must have reached states of limiting stress. Taking shear strain as a yield criterion, the observed mode of yielding agrees well with Mode I. It is important to notice that the shear strains shown in Fig. 3.1 were observed near collapse at large displacements, probably between up and up in Fig. 2.9.

The shear strain at yield initiation can also be observed from the results of model tests (gravity type) performed by Cording et al. (1976). In their model tests a specified displacement pattern instead of uniform pressure was imposed at the tunnel perimeter. Fig. 3.2 shows shear strain from their finite element simulation of these model tests. It can be seen that fairly high shear strains are concentrated along a plane through the springline inclined approximately 65° to horizontal. Asumming  $\phi=30^\circ$ , the predicted  $\beta$  from Fig. 2.10 is about 67° which agrees with the observed one (note the boundary condition at tunnel wall are different in two situations).

All these model tests of gravity type (Atkinson and Potts, 1977, Cording et al., 1976) were carried out in plane strain condition. Hence, the K value is presumably given by the relationship  $K = 1 - \sin \phi$ . With  $K = 1 - \sin \phi$ , the mode of yielding should be Model I-1 or I-2. The propagation of high shear strain zones is comparable to that predicted by Mode I. These results therefore support the finding that Mode I takes place for K < K.

Potts (1976) performed a series of model tests (static test type) in sand for K =0.5 and 1.0. The development of shear strain around the tunnel was not reported and only the result near the collapse was recorded. Unfortunately, it is not possible to make use of these results to illustrate the occurrence of Modes I and II for K =0.5 and 1.0 respectively. However Mode II can be recognized very well from model tests of tunnels in over-consolidated Kaolin at K =1.0 (Atkinson et al., 1974). The tunnel pressure was

reduced from 140 kPa to zero while the surface surcharge 'pressure was held constant at 140kPa. Contours of shear estrains at p = 27 and 7 kPa are shown in Fig. 3.3a and b. Comparison of the contours of shear strain demonstrates a marked change in the pattern of deformation. At p = 27kPa, the shear strain contours are approximately concentric with the tunnel axis indicating radial inward deformations (Mode (1). At this stage,, contours of shear strain (say 1.0%) extend farther near the floor than the roof. This response is explained by the fact that under uniform internal pressure the deviatoric stress is higher at the floor than . the roof due to the gravity stress increasing with depth. At p = 7kPa, the shear strain contours have become non-uniform and a pair of zones of high shear strain extends upward toward the surface (Mode II to Mode II #I). The above observation verifies that Mode II develops for K = 1.0.

The initial insitu stresses (K) not only govern the mode of yielding but also influence the displacement of soil around the tunnel. Potts (1976) reported surface settlement profiles for a set of model tests in sand with K = 0.5 and 1.0 (see Fig. 3.4). The surface settlement for K = 1.0 is much smaller than that for K = 0.5 even a smuch lower support stress. This difference is attributed to the fact that tangential arching is higher with K = 1.0 as explained in Section 2.2.4.3. This causes an increase in resistance against the downward movement of the soil block above the crown and reduces the surface settlements. This implies that the potential for damage due to surface settlement becomes more critical for Mode I than for Mode II.

### 3.2.2 Field Measurements

Two cases histories of tunnelling in Edmonton till (Eisenstein et al., 1981, Branco, 1981) are available where sufficient measurements of the displacement field around the tunnel and the pressure on the lining system were made. These two cases are now examined in detail by use of the ground convergence concept. Other cases documented by Attewell (1977) and O'Reilly and New (1982) will also be studied and their results along with those from model tests and FE analyses then will be synthesized.

## Tunnelling in Edmonton Till

Two tunnels were built using very similar construction methods and in comparable soil conditions the first tunnel (EXP tunnel, Eisenstein et al., 1981) is a small diameter tunnel (D=2.56m) driven by a full face TBM at a depth of H=24m at the test section. The primary lining comprises segmented steel ribs and timber lagging. Fig. C.1 shows the contours of measured vertical and horizontal displacements . around the tunnel. The pressure pattern on the flexible lining was found to be almost uniform with ah average radial pressure  $\phi f$  12% of full overburden pressure. The second tunnel (LRT tunnel/) is a large diameter tunnel (D=6.1m) driven at a depth of 10m at the test section. The excavation method and installation of the primary lining are similar to those used in the EXP tunnel. The measured displacement at the crown was about 16-22mm. The load on the support system ranged from 0.18 to 0.24 of full overburden pressure.

The geotechnical properties of Edmonton Till in both locations are studied by Matheson (1970), May and Thomson (1978) and El-Nahhas (1980). The ground parameters used in the ground convergence calculations are shown in Table 3.1 for each case. The K -value is assumed to be close to unity  $(K_0=0.8)$ .

Imposing all parameters such as  $p_i/p_i$ , H/a, and  $K_i$  in the diagrams of modes of yielding (Figs. 2.7 and 2.8) shows that two tunnels lie within the Regime J of Mode II. A continuous yield zone must have developed around the tunnel. This implication permits one to use the "hole-in-plate" model (Section 2.2.4.2) to approximate the ground responses near the tunnel. The ground convergence curves for expected ranges of Young's moduli are calculated for two tunnels and plotted in Figs. 3.5 and 3.6 respectively. On these figures are also shown the relationship between the support pressure and the extent R of the yield zone. For the EXP tunnel an almost uniform radial pressure is predicted in the range of  $0.18\pm0.05p$  for the measured displacement range u/a=0.027 to 0.035. This pressure and the distribution correspond closely with the measured (0.12p). The extent of 'yield zone should be slightly oval-shaped with a maximum extent at the roof (R/a=1.8 $\pm$ 0.2). If the same procedure is applied to the LRT tunnel, the predicted pressure  $(p_i/p_i=0.12 \text{ to } 0.15)$  is comparable to the measured (p./p.=0.18 to 0.24). A reduced E value to 80MPa may yield better agreement. The average extent of yield zone is expected to be about R/a=1.4. However, at this stage GCC (Rf) indicates that the pressure at the roof starts to remain constant at

0.11p for u /a≥0.01. Further increase in roof displacement would not compensate with decrease in ground pressure, but could accelerate the propagation of yield zones to the ground surface and thus induce large settlements.

It is of practical importance to evaluate marginal safety against the collapse state, u for these two tunnels. The p values calculated from Eqn. 2.3 are 0.03p and 0.09p for EXP and LRT tunnels respectively. Imposing these values on the GCCs determines the u a values in ranges of 0.09 to 0.13 and 0.013 to 0.018 for the two tunnels. Comparisons between u and and u a clearly indicates that EXP tunnel has a higher marginal safety against the collapse state than the LRT tunnel. Because of small marginal safety (as expected for shallow tunnels), it justifies the recommendation proposed by the German Tunnelling guidelines that for shallow tunnels ground pressure of full overburden at the crown should be used in design.

#### Other Cases

Semi-empirical approaches for predicting the surface settlement profile (e.g., Peck, 1969) are based largely on field measurements of various case histories. In these approaches the location of point of inflection (i) is assumed to be dependent on only the depth ratio (H/a). Since some of the important factors governing the tunnel behaviour (e.g., strength-deformation properties of ground, stress-relief and boundary conditions) are not included in these approaches, å wide scatter of the observed data is

expected.

The surface settlement-support stress relationship (i-S/S-p/p) has been introduced in Section 2.2.4.3. Case histories have been collected and presented in Tables C.1 and C.2 (Appendix C). These case histories, grouped in cohesionless and cohesive soils are used to verify the validity of the dependence between the p/p and i/a discussed earifler (Figs. 2.15 and 2.16c).

#### a) Cohesionle\$s

Cohesionless soils have a smaller variation in strength-deformation properties as compared to cohesive soils. Hence it is reasonable to assume that the strength-deformation aspect does not exert significant influence on the field measurements in case histories listed in Table C.1. These results are plotted in Fig 3.7a and b, and the following aspects are observed.

- 1. Fig. 3.7a is a plot of H/a versus i/a similar to the one given by Peck (1969). It can be seen that a significant number of case histories (especially, H/a>5) lies outside the boundaries predicted by Peck (1969). This implies that i/a depends on H/a as well as other factors.
- The same case histories were grouped in four depth ranges and the results are replotted in Fig. 3.7b to observe the influence of the H/a ratio. On Fig. 3.7b is also drawn boundaries for tunnels of H/a<5 and H/a>5.

  From comparison between Fig. 2.16b and Fig. 3.7b, similar features are identified: (i) for shallow

tunnels i decreases to a limiting value with a small increase in S , and (ii) for deep tunnels i also decreases with S but at a slower rate. The results from model tests (denoted by squares) tend to yield the lower bound of i/a values because the measurements were taken near collapse. For the depth range of H/a<5 the ratio i/a approaches its limiting value when S /a  $_{\rm S}$  exceeds 1.0%. This is explained by the fact that a small displacement at the crown (S /a) will cause the yield zone to propagate to the ground surface and i is predetermined by  $\beta$  (e.g., S /a=0.4% for yield zone reaching the ground surface in FE analyses Case ST1). For deep tunnel the dependency between i/a and S /a becomes more pronounced, i.e., i will reach its limiting value through large S.

3. The influence of K cannot be properly assessed because the K value was seldom recorded in the case histories. However, for normally consolidated sand K is approximated by  $(1-\sin\phi)$ . Small variations of  $\phi$  in sand limits the range of K. From the FE Analyses it was found that at a given p / p the i/a will be larger in high K than low K.

# b) Cohesive Soils

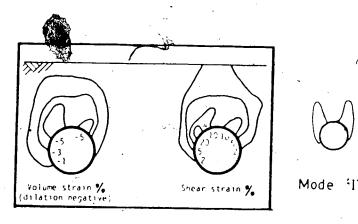
Cohesive soils exhibit a wider variation in strength-deformation properties (E and c) than cohesionless soils. Thus in order to isolate the strength-deformation effects for better interpretation of

the field measurements, the case histories are subdivided into categories of (1) c = 0 to 100kPa and (2) c > 100kPa. Their results are plotted in Figs. 3.8 and 3.9 respectively. Several aspects of practical interest are dissected.

- 1. For  $c_u$ =0 to 100kPa (Fig. 3.8) similar observations as for cohesionless soils are found. For shallow tunnels (H/a<5) the i/a becomes independent of S/a when S/a exceeds 1.2%. Deep tunnels show strong dependency between i/a and S/a.
- 2. From Fig. 3.9 (c >100kPa) the surface settlement is relatively small (<1.0%) and this may be attributed to the high strength (or stiffness). There seems to be a linear relationship between H/a and i/a, i.e., i/a increasing with H/a.
- 3. Generally the cohesive soils have a less critical settlement trough than the cohesionless soils. This may be due to higher K (close to 1.0) in overconsolidated cohesive soil.

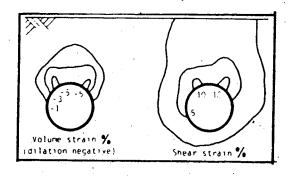
Table 3.1 Input Data for Constructions of GCC for EXP and LRT Tunnels  $\alpha$ 

	EXP Tunnel en El-Nahhas, 1980)	LRT Tunnel (after Branco, 198
Soil Model Per		astic. Cohesionless I flow rule)
Angle of (Internal, (deg.)	30	40
Young's Modulus	40 - 80	100 - 150
Poisson's Ratio	0 4	0.4
Depth to Tunnel axis	24 m ,	10 m
Tunnel Radius	1.28 m	3 05 m
Depth/Radius (H/a)	18 8	3.4
Surface Settlement	7 - 12 mm	9 - 10 mm
Crown Settlement	40 - 45 mm	16 - 22 mm





(a) Gravity test





Mode 1

(b) Centrigual test

Figure 3.1 Contours of Volume and Shear Strains in Sand (modified from Atkinson and Potts, 1977)

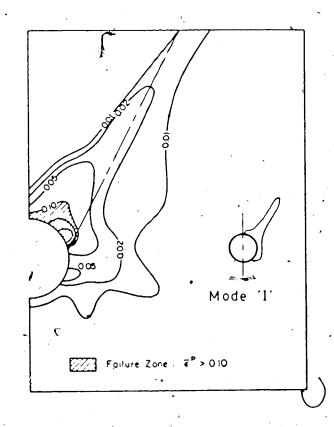


Figure 3.2 Contours of Shear Strains in Sand (modified from Cording et al., 1976)



Figure 3.3 Contours of Volume and Shear Strains in Kaolin (modified from Atkinson et al., 1974)

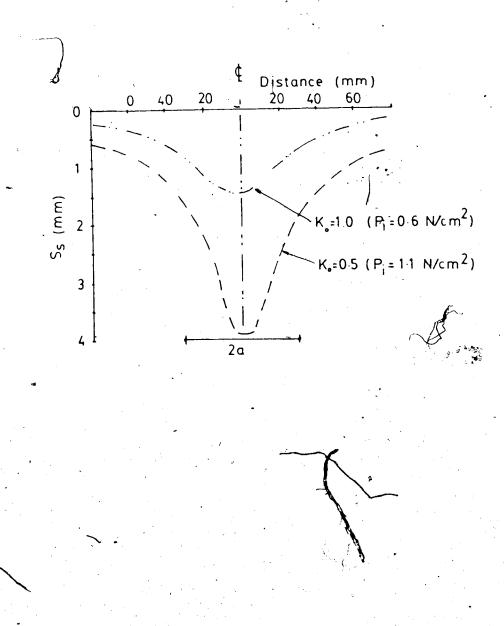


Figure 3.4 Surface Settlement Profiles in Sand (modified from Potts, 1976)

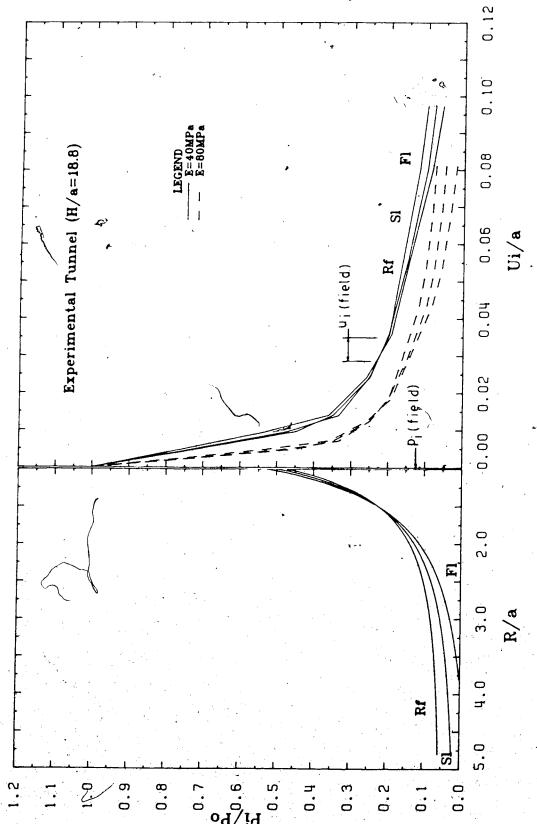


Figure 3.5 GCC and Support Pressure-Extent of Yield Zone

(EXP Tunnel)

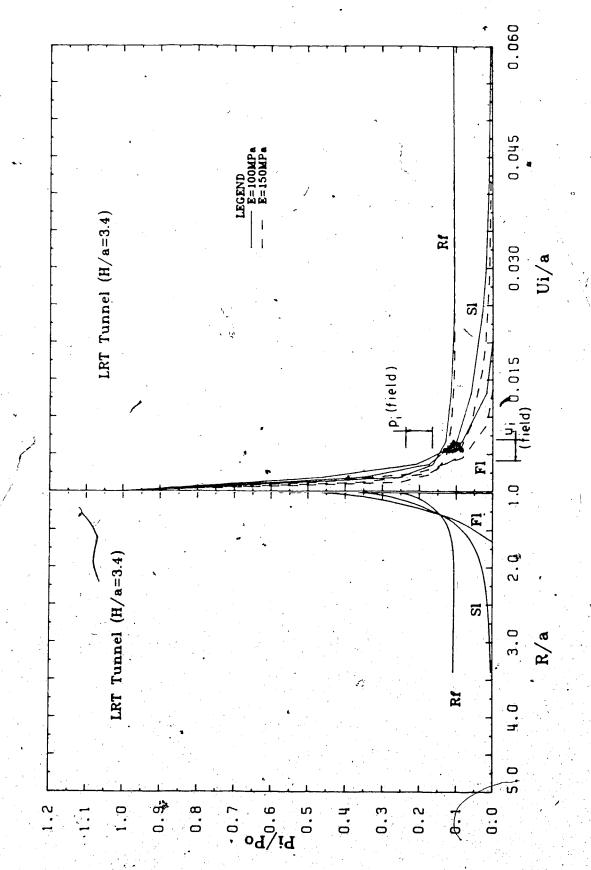
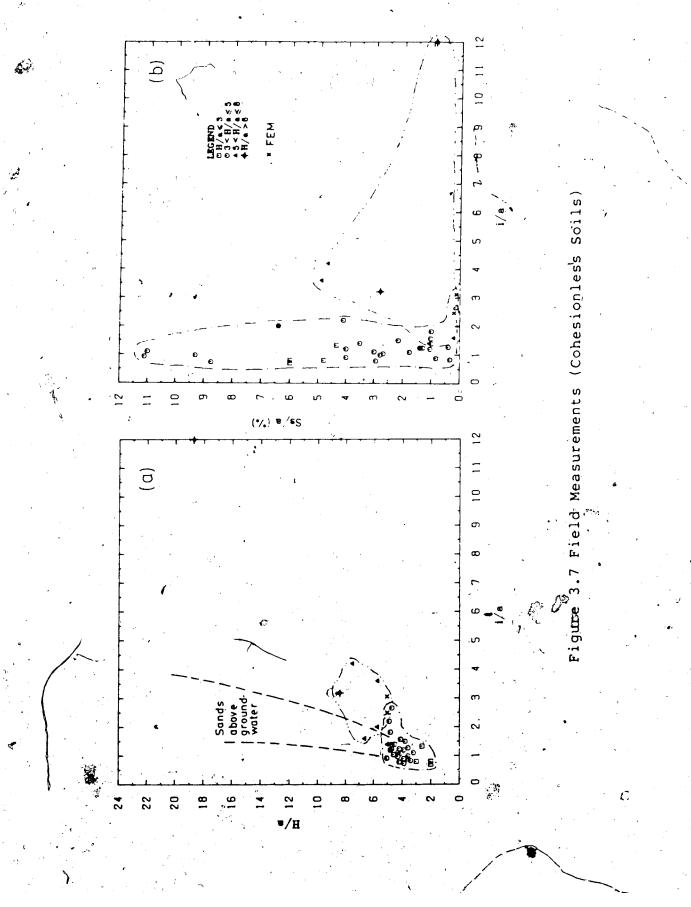


Figure 3.6 GCC and Support Pressure-Extent of Yield Zone



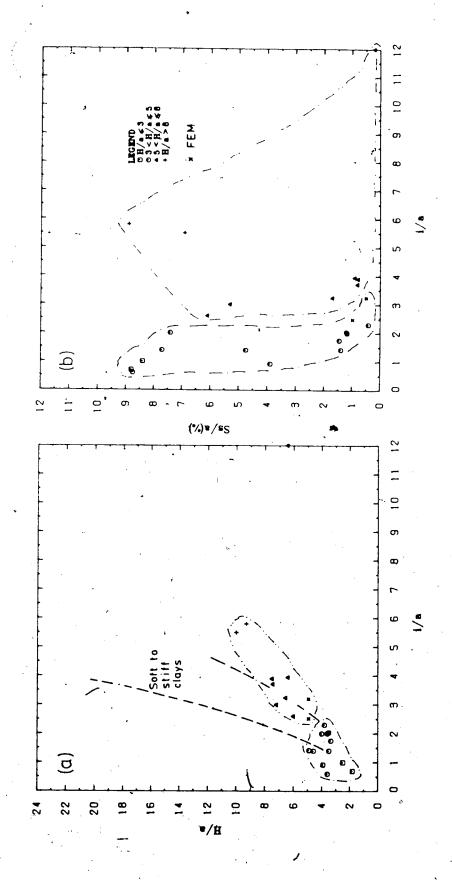
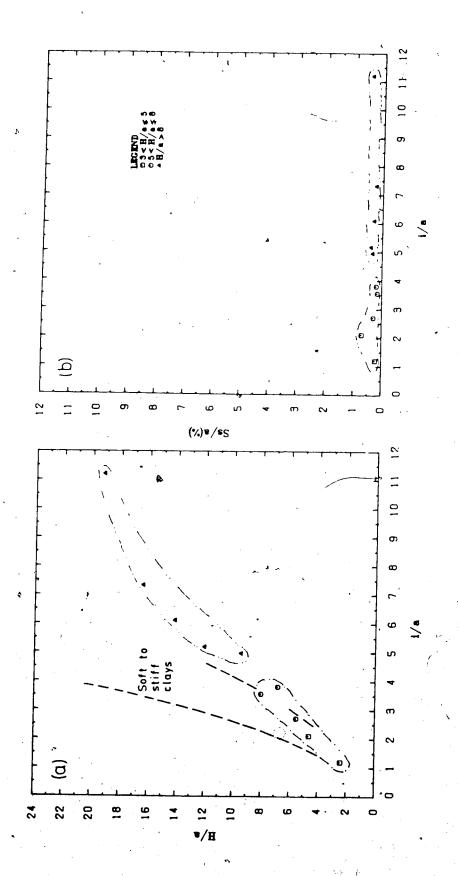


Figure 3.8 Field Measurements (Cohesive Soils, Cu= 0 to

100kPa)

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Figure 3.9 Field Measurements (Cohesive Soils, Cu>100kPa)

#### 4 . DESIGN OF SHAFT IN SOIL

#### 4.1 Introduction

Geotechnically, the design of a shaft in soil consists of two major steps: (1) to design the shaft lining to prevent instability of shaft wall, and (2) to estimate the soil movement associated with shaft construction. Although these two tasks are interrelated, they are usually dealt with separately.

Most of the currently available design approaches are based on limit equilibrium methods, e.g., Terzaghi (1943), Berezantzev (1958) and Prater (1977). This type of analysis considers only equilibrium of the forces acting on an assumed failure mass defined by hypothetical rupture surfaces. The formation pressures on the shaft lining are determined by satisfying force equilibrium, but these do not correspond with the actually expected pressure which depends on such factors as ground deformation, insitu stress, and ground strength-deformation properties. As the constitutive relationship of the ground is not considered explicitly, no closed-form solution is available for the determination of the displacements at the shaft wall and inside the soil mass. In these design methods, it is attempted to control and limit the displacements by choice of a suitable factor of safety. Furthermore, excessive yielding is prevented by the selection of an appropriate construction sequence. Besides for simplicity of analysis, there is no particular reason why one has to separate the determinations of the lining pressure from the ground deformation in shaft design. The convergence confinement method provides an analytical framework to predict the formation pressure and the soil deformation simultaneously. However, this method requires additional information about the ground and support, such as insitu stress, constitutive relationship of ground, support characteristics and construction details. Also, some restrictive assumptions must be made to obtain a closed-form solution for the ground convergence curve, e.g., plane strain condition. Hence, this method has, so far, only been applied to two-dimensional circular, horizontal openings like deep tunnels.

Since the analysis of a shaft is a three-dimensional problem, the direct application of the convergence confinement method is not adequate to describe the shaft behaviour. Hence, some techniques such as force equilibrium have to be included to account for the behaviour in the third dimension. By combination of the force equilibrium technique and the convergence confinement method (CCM), it is possible to derive ground convergence curves (pressure-displacement relationships) for the shaft wall, to rationally assess the required support pressure and to evaluate the limits of applicability of conventional limit equilibrium methods for shaft design.

In the following, conventional design methods for shafts are reviewed briefly and the proposed confinement convergence method for shaft design is then presented. Typical mechanisms of shaft behaviour including yield initiation, modes of propagation of the plastic zone and gravity effects due to vertical arching are studied in

detail. The second part presents examples generated by the finite element method and a comparison with results of the proposed CCM as well as the limit equilibrium methods.

The proposed technique has been applied to several case histories (Lade *et al.*, 1980; Muller-Kirchenbauer *et al.*, 1980; Britto and Kusakabe, 1984; and Kaiser and Wong, 1984) and good agreement has been achieved.

# 4.2 Review of conventional Shaft Design Techniques

### 4.2.1 Terzaghi s Method

Based on the state of stress in the vicinity of a drill hole, Terzaghi (1943) derived the earth pressure distribution for a circular shaft located above the water table in cohesionless soil.

 $(\sigma - \sigma)$  produces shear failure in vertical surfaces.

Terzaghi observed the downward and inward movements of the shaft wall, and argued that near the surface these movements of the shaft wall cause the tangential stress (intermediate stress) to increase until in the limit, the surface of plasticity becomes vertical instead of inclined (even at low  $K_{\hat{a}}$ ) (see Fig. D.1, Appendix D). This assumption permits the application of the equations describing the state of stress near the wall of a drill hole (hole-in-plate) for the shaft case. With a known state of stress and a linear Mohr-Coulomb failure critérion, the extent of the plastic zone can be determined. Terzaghi calculated the minimum support pressure to maintain the limit equilibrium state of such a plastic zone (McCreath, 1980). Terzaghi's approximation of a three-dimensional state may be viewed as a method of accounting for the necessary limit equilibrium condition of a potential failure block under the action of gravity.

#### 4.2.2 Berezantzev's Method

Berezantzev (1958) proposed the method of earth pressure calculation on the retaining walls of circular cylindrical form, based on the solution of axisymmetrical problems of the limit equilibrium theory (Fig. D.2). He assumed the soil medium surrounding the shaft yielded along inclined surfaces dipping toward the shaft.

In Berezantzev's method, the states of stress around the shaft are given by two differential equations of equilibrium. These equations are made determinant by

introducing the Mohr-Coulomb failure criterion and a simplified assumption of equalization of principal stresses. The stress conditions are determined by solving the differential equations using the numerical "step-by-step computation" technique (Sokolovsky, 1954).

#### 4.2.3 Prater's Method

Prater (1977) computed the earth pressure on shaft linings using a limit equilibrium analysis of an assumed truncated, hollow cone failure block. He included several parameters, such as the earth pressure coefficient K and the tangential forces induced by the inward movement of the cone, and investigated their influence on the calculated support pressure. He found that the tangential force (T) induced by the inward movement had an outward acting wedge component (F) that increases with depth and increasing earth pressure coefficient (see Fig. D.3). Consequently, Prater's method leads to a reduction in support pressure with great depth.

The active earth pressure for a shaft wall and the earth pressures as a function of the depth factor (H/a) for typical soil parameters predicted by these three methods, are shown in Fig. 4.1 for comparison. The distribution of earth pressure predicted by Terzaghi (1943) and Berezantzev (1958) are almost identical (for  $\phi$ =30°, c=0) and reach an asymptotic value at a depth of H/a>4. Prater predicts slightly higher pressures for this case at H/a=4 and zero pressure at H/a=9. It is intuitively obvious that the assumptions made by Prater are not valid at great depth.

There must be a minimum support pressure at great depth in yielding ground and, hence, Prater's solution cannot be valid below a critical depth unless unreasonably large movements are permitted (see also discussion on convergence confinement method).

#### 4.2.4 Convergence Confinement Method

The convergence confinement method is an analytical approach to predict the stresses and displacements in the soil around an underground opening and in the support elements. This method describes how ground and support interact with each other. The ground behaviour is represented by a ground convergence curve (GCC) and the lining by a support confinement curve (SCC). The former describes the ground convergence in terms of the internal pressure relief while the latter relates the confining pressure acting on the lining to its deformation. The equilibrium condition for the ground-support interaction is given by the interaction of the GCC and SCC, as illustrated and described in detail earlier in Fig. 1.1.

Two basic assumptions are normally made for derivation of the GCC and the SCC: (1) two-dimensional plane strain model; and (2) plastic yielding controlled by tangential-radial stress difference  $(\sigma_{t} - \sigma_{r})$ . These two conditions are generally applicable to tunnelling at depth but they do not necessarily satisfy conditions near a shaft.

Brown et al. (1983) presented a summary of currently available GCC formulations of different material models. The characteristics of the SCC for different liner or tunnel

support types were summarized by Hoek and Brown (1981) and do not need to be reviewed here. However, whether and under what conditions the shaft behaviour can be simulated by a "vertical tunnel" or by the two-dimensional, "hole-in-plate" theory needs to be investigated further.

The behaviour of a shaft is affected by gravitational forces and is a three-dimensional problem in which three stress components (\$\sigma\_{\text{t}}\$, \$\sigma\_{\text{a}}\$ and \$\sigma\_{\text{t}}\$) must be considered. The mode of yielding and its initiation and propagation, depend on the initial insitu stress state described by \$K\_{\text{c}}\$. The shaft response to stress relief during excavation is further complicated by the influence of the free stress boundary near the ground surface resulting in a linearly varying horizontal stress field with depth. In order to apply the two-dimensional "hole-in-plate" model to determine the relationship between support pressure and shaft wall displacement, it is necessary to understand the mechanisms of shaft behaviour so that some adjustments may be made to account for these aspects in the analysis.

The excavation of a shaft can be simulated by a stress relief ( $\Delta\sigma_{r}$  at the shaft wall causing the surrounding soil to deform both horizontally and vertically). Excessive stress relief may induce yielding near the opening and cause permanent plastic deformations. The magnitude of support pressure, wall displacement and extent of plastic zone are interrelated. The stress relief during shaft excavation causes stress redistribution near the opening, and thus induces (1) horizontal and (2) vertical archings as illustrated in Fig. 4.2. The horizontal arching is explained

by the development of hoop stresses in a horizontal plane, i.e., an increase in the tangential stresses. The vertical arching arises from the formation of a plastic zone of limited extent around the shaft with a tendency to move in the vertical direction. Vertical arching develops between the shaft support and the surrounding unyielded soil mass. Collapse is prevented by these two mechanisms that need to be considered to describe mechanisms of shaft behaviour.

## 4.2.4.1 Simulation of Horizontal Arching

The horizontal arching is best represented by the "hole-in-plate" model. The shaft may be viewed as a stack of plates with a hole subjected to insitu stress: (Fig. 4.2b). The differential wall displacement along the shaft is assumed to cause negligible shear developed between plates. Hence, each plate can be treated individually as a two-dimensional problem. The ground convergence curve and the extent of the yield zone are characterized by the material properties, the magnitude of stress relief or wall displacement permitted, and the initial state of stress. Solutions presented by Brown et al. (1983) are applicable to this condition but are restricted to two-dimensional plane strain  $(\sigma_1 - \sigma_1)$  yielding and initial uniform insitu stresses. Other modes of yielding (e.g., due to  $\sigma_1 - \sigma_2$ ) do result in different ground responses as will be discussed later.

4.2.4.2 Simulation of Vertical Arching (Gravity Effect)
Owing to the stress relief in the radial direction, the
soil particles around the shaft opening will undergo
vertical displacement under the effect of gravity \*simply

referred to as "gravity effects" in this chapter). The soil particles may remain stable and be kept in equilibrium position by the shear resistances derived from the interactions with the adjacent stationary soil mass and the shaft lining. This phenomenon is referred to as vertical arching (see Fig. 4.2a). If yielding around the shaft takes place and creates failure blocks that have a tendency to slide into the shaft opening under the gravity effect, an external support pressure may be required to prevent the instability. For example, in the case of yielding induced by the tangential-radial stress difference, the strength of the ground is reached along the vertical log-spiral surfaces only for horizontal movement. The soil will, however, still be able to resist forces due to gravity for movement in the vertical direction if the strength is not reached for movement in this direction. Hence only part of the gravity component will create additional stresses on a shaft support (besides support stresses in horizontal arching) if yielding in vertical and horizontal directions are exceeded simultaneously.

The support stresses required in resisting the gravity effect can be determined by considering the limit force equilibrium of the yield zone. This approach has been adopted in calculating horizontal stresses on silo walls (Kendal, 1980). Handy (1984) used the same approach with some adjustment in soil-wall interaction to estimate the lateral pressures behind retaining walls, and claimed satisfactory correlation with model test observations.

Fig. 4.2(d) shows forces acting on a horizontal differential soil element within the yield zone. The vertical stresses are assumed to be uniform at each depth. Summation of vertical forces in Fig. 4.2(d) give

 $d\sigma_{V} = \left[ \gamma - (2\pi\sigma_{V}/A) \left( K_{SS} R / \sin(a) + K_{WW} \mu_{a} \right) \right] dh$ 

where: A - the sectional area of the plastic zone

a - the inclined angle as shown in Fig. 4.2d (This angle can be taken as 90° on the conservative assumption)

K - K coefficient at soil to soil interface
 K - K coefficient at wall to soil interface
 μ - frictional coefficient at soil to soil interface
 μ - frictional coefficient at soil to wall

Integration of Eqn. 4.1 will give the vertical stress distribution along the shaft depth, and thus the horizontal stresses required to prevent instability due to the gravity effects. Detailed treatments on Eqn. 4.1 are given in Appendix D.1.

interface

In Eqn. 4.1, the geometry of the plastic zone and the soil-lining properties govern the stress distributions. The extent of the plastic zone which was usually assumed in most past works (e.g., Handy, 1985) can be predicted with reasonable accuracy (see section on Mechanisms of Shaft Behaviour). The soil-lining properties are discussed in

Appendix D.

Several numerical examples have been generated to illustrate the role of the gravity effect in shaft design. Fig. 4.3 shows dimensions of the shaft and extent of plastic zone. Results of analyses following equations of Appendix D.1 on these numerical examples for cohesionless and cohesive soils are plotted in Figs. 4.4 and 4.5 respectively, along with the soil properties.

Fig. 4.4 indicates that for cohesionless soil, the support pressure due to gravity effect increases as the extent of yielding increases. The distribution of the support pressures also depends on the configuration of the plastic zone. The support pressure reduces to zero at great depth for cone shaped plastic zone, and to a constant value for plastic zone of constant radius which was assumed by Terzaghi (1943). Near the surface, the pressure is close to the active earth pressure (K case), implying that the gravity effect dominates at shallow depths.

In the cohesive soil, the support pressure to prevent any instability due to gravity effect also increases for a larger yield zone. A distinct difference in vertical arching action exists between cohesionless and cohesive soils. For cohesive soils, the support pressure applied along the whole shaft depth does not enhance the stability because the shear strength of the cohesive soil is independent of the confining pressure. The support pressure has to be applied at the bottom to inhibit the collapse mechanism (Britto and Kusakabe, 1983).

## 4.3 Design Based on Convergence Confinement Method (CCM)

It has been shown that the behaviour of the shaft can be described by horizontal and vertical arching around the shaft. These two arching actions can be quantified by use of the convergence confinement method and the gravity effect. Hence, it is possible to derive a relationship between the support pressure and the displacement allowed during the construction. The CCM with inclusion of gravity effect for shaft analyses can be summarized in the following steps (Fig. 4.6):

- Identify the mode of yielding at the shaft wall. This mode, dependent on the initial in-situ stress (K), governs the extent of the plastic zone as well as the shape of convergence curve.
- Calculate the ground convergence curves along with the extent of plastic zone, using appropriate two-dimensional model for various depths (h<sub>i</sub>) (see Fig.4.6(a)).
- 3. For specific displacements (u<sub>s</sub>), establish pressure versus depth and plastic zone versus depth relationships from (2) (Fig. 4.6(b)).
- 4. With the configuration of the plastic zone around the shaft, determine the support pressure due to the gravity effect (p) with depth (Fig. 4.6(c)).
- 5. Two pressure distributions due to horizontal and vertical archings form an envelope of design pressure (Fig. 4.6(c) is the support pressure for given displacement u\_).
- 6. Adjust the design pressure envelope at the bottom of

the shaft due to face effect (Panet and Guenot, 1982).

It follows from the above calculation steps that a particular displacement corresponds to a unique support pressure distribution. Hence, it is possible to evaluate the limit equilibrium method of shaft design (e.g. Berezantzev, 1958) in terms of displacement.

#### 4.4 Mechanism of Shaft Behaviour

Prior to excavation, a soil element adjacent to the shaft wall is subjected to the initial stresses, as shown in Fig. 4.2b. The excavation of the shaft at any particular level may be simulated by progressively reducing the internal support pressure (radial stress,  $\sigma$ ). For the axisymmetrical case, as long as the soil material remains in the elastic range, the stress distributions in the plane of the section are given by Eqns. 4.2 to 4.4 (Lame's equations; Terzaghi, 1943).

$$\sigma_{v} = \gamma h = p_{o}$$
 4.2

$$\sigma_{r} = K_{0} p_{0} - [K_{0} p_{0} - p_{1}](a/r)^{2}$$
4.3

$$\sigma_{t} = K_{0} + [K_{0} - p_{i}](a/r)^{2}$$
4.4

It is important to realize that Eqns. 4.2 to 4.4 are derived based on the following assumptions: (i)  $\sigma_{\rm v}$ ,  $\sigma_{\rm r}$  and  $\sigma_{\rm t}$  are principal stresses, (ii) shear stresses along the shaft wall are small for a constant wall displacement and (iii) the bottom of the shaft is remote from the section.

As  $p_i$  is reduced, stress differences among  $\sigma_v$ ,  $\sigma_r$  and  $\sigma_t$  are generated at the shaft wall due to the increasing  $\sigma_t$  and decreasing  $\sigma_r$  (i.e., horizontal arching). If the stress difference exceeds the strength of the soil, three possible alternatives of stress combinations may initiate yielding. Plasticity could be generated by stresses in a vertical surface by the tangential-radial stress difference  $(\sigma_t - \sigma_v)$ . The onset of the plasticity (the mode of yield initiation) depends on the value of  $K_v$  and strength parameters of the soil. For simplicity, the following derivations are for cohesionless material only. The formulations for cohesive material are given in Appendix D.2.

Assuming a purely frictional material with a linear Mohr failure criterion, the maximum stress ratio which may be sustained is:

$$\sigma_1/\sigma_3 = N = \tan^2(\pi/4 + \phi/2)$$

Three yielding mode criteria can be calculated from Eqns. 4.2 to 4.4, and the support pressures (p<sub>i</sub>) corresponding to three possible modes of yield initiation are listed as follows:

For 
$$\sigma_t^{-\sigma_r}$$
:  $p_i = 2K p_o/(N+1)$  4.6

For 
$$\sigma_{v}^{-\sigma}$$
:  $p_{i} = p_{o}/m$ 

For

$$\sigma_{t} - \sigma_{v} : p = (2K_{o} - N)p_{o}$$

The largest value of p will, of course, govern the mode of yield initiation, which can be exptessed in terms of required K:

Mode A for 
$$\sigma_{t}^{-\sigma}$$
:  $(N+1)/2 > K_{o} > (N+1)/2N$  4.9

Mode B for 
$$\sigma_{v} - \sigma_{r}$$
: K <  $(N+1)/2N$  4.10

Mode C for 
$$\sigma - \sigma$$
: K > (N+1)/2 4.11

Mode A is the mode commonly evaluated for tunnels in yielding ground. For example, for a typical value N of 3 ( $\phi$  = 30°),

Mode A governs for  $0.67 < K_O < 2.0$  (Fig. 4.7a); Mode B governs for  $K_O < 0.67$  (Fig. 4.7b); and Mode C governs for  $K_O > 2.0$ .

Mode C, although possible, has been neglected in the following analysis because it is of less practical significance in soft ground. The boundary between the other two modes of yield initiation at the wall can be described by a critical K -value, K . Mode A is observed for K > K or For these two cases of Modes A and B, the relationships among support pressure, wall convergence and extent of plastic zone can be derived separately.

a) Mode A 
$$(\sigma_t - \sigma_r)$$
 at K > K cr:

Fig. 4.8 shows the sequential stages for Mode A as the internal support pressure is reduced. The vertical stress always acts as an intermediate stress. Initiation of yielding and propagation of the plastic zone (R tr) are controlled by the tangential-radial stress difference. The relationships between the support pressure and the radial convergence at the wall for yielding ground (i.e., ground convergence curve) have been well studied by other authors (e.g., Brown et al., 1983). For simplicity, the pressure-displacement relationships are calculated herein by application of the model proposed by Ladanyi (1974) assuming associated flow rule, constant volume change and plane strain condition. Other models could be incorporated, but will not be treated in detail.

During the first stage (Fig. 4.8), the vertical stress remains constant, radial stress decreases and tangential stress increases according to Eqn. 4.2 to 4.4. Yield initiation of Mode A occurs if the condition of

$$\sigma_{t}/\sigma_{r} = N \text{ or } p_{i} = 2K p_{o}/(N+1)$$
 4.12

is satisfied. The magnitude of wall convergence with p from Eqn. 4.12 is given:

$$u = a [K p - p] (1 + \nu) ]/E$$
 4.13

Further relaxation of the fictitious internal support pressure  $(p_i)$  causes the propagation of the plastic zone, and the tangential stress decreases to satisfy the failure

criteria. The extent of the plastic zone and the radial wall convergence (Ladanyi, 1974) are given by:

$$R_{tr} = a[2K_{0}p/(N+1)p_{1}]^{1/(N-1)}$$
4.14

$$u_i = a\{1-[(1-e_{av})/(1+A_v)]^{1/2}\}$$
 4.15

For a given  $p_i$ , the plastic radius,  $R_i$  can be calculated from Eqn. 4.14, and thus the corresponding wall displacement,  $u_i$  can be determined from Eqn. 4.15. Ground convergence graphs of  $p_i$  versus  $u_i$  and  $p_i$  versus  $R_i$  can be established, as schematically shown in Figs. 4.6.

As the tangential stress decreases during yielding and becomes equal to the vertical stress (Stage(2) of Fig. 4.8), p becomes equal to K p. Substitution of this p value into a o Eqns. 4.14 and 4.15 will yield R and u for this stress state, i.e.,  $\sigma_{\rm t} = \sigma_{\rm t} > \sigma_{\rm t}$  and  $\sigma_{\rm t} = \sigma_{\rm t} > \sigma_{\rm t}$  and  $\sigma_{\rm t} = \sigma_{\rm t} > \sigma_{\rm t}$  and  $\sigma_{\rm t} = \sigma_{\rm t} > \sigma_{\rm t}$ 

During Stage(3) at which the internal pressure continues to decrease, the tangential stress will decrease to satisfy the failure criteria and equilibrium state, whereas the vertical stress drops due to arching between the shaft support and the elastic ground. The value of  $R_{\rm tr}$  and  $u_{\rm i}$  at this stage can again be determined by Eqns. 4.14 and 4.15. An example is given later.

b) Mode B  $(\sigma_{v} - \sigma_{r})$  at K < K < c:

Fig. 4.9 shows the stress distributions in a horizontal section through a shaft at a particular depth in sequential stages of reducing the internal support pressure, i.e., simulating the shaft excavation.

At Stage (1), the yielding initiates at the wall as the vertical-radial stress difference violates the yield criterion. The stress distributions are given by Eqns. 4.2 to 4.4. Hence, the wall displacements are elastic and can be calculated using Lame's equation assuming elastic material behaviour.

$$u_{i} = a[(K_{0}p_{0}-p_{i})(1+\nu)]/E \text{ and } p_{i} = \sigma_{v}/N = K_{0}p_{0}$$
or  $u_{i} = a[(K_{0}-K_{0})p_{0}(1+\nu)]/E$ 
4.16

As p is further reduced,  $\sigma_{\rm t}$  still increases and  $\sigma_{\rm v}$  decreases due to vertical shear stress induced by vertical downward displacement (Stage (2) in Fig. 4.9). Yielding propagates outward from the wall and the radial and tangential stress distributions are still given by Eqns. 4. and 4.4. But these stresses at the wall (r=a) also satisfy the failure criterion, i.e.:

$$\sigma_{t}/\sigma_{r} = N$$

The extent of the plastic zone due to Mode B yielding at this stage is determined by equating the radial stresses at the elastic-plastic interface, i.e., the stress continuity. At the elastic-plastic boundary, the elastic

radial stress is given by Eqn. 4.3, and the radial stress in the plastic zone is given by:

$$\sigma = \sigma / N \text{ and } \sigma = p$$

$$V = 0$$

Equating Eqns. 4.3 and 4.18 yields

$$R_{vr} = a\sqrt{[(K_{o}p_{o}-p_{i})/(K_{o}-K_{a})p_{o}]}$$
 4.19

Once the extent of the yield zone has been determined, it is of interest to predict the related wall displacement allowed in Stage (2). One can make use of the same model as in Mode A yielding, i.e., Eqn. 4.15.

At the end of Stage (2), the tangential stress becomes equal to the vertical stress and the support (radial) stress is given by Eqn. 4.12, or

$$p_i = 2K_0 p_0 / (N+1)$$

4.20

Further relief in  $p_i$  as allowed in Stage (3) causes a decrease in  $\sigma_i$  and  $\sigma_i$  in order to satisfy the failure criteria given by Eqns. 4.17 and 4.18. At this stage, the extent of the plastic zone ( $R_i$  in Fig.4.9) is still governed by Mode B yielding, but the radial stress distribution will be different from that in Eqn. 4.3) because of the need to satisfy Eqn. 4.17. Hence, in order to determine the plastic zone in this stage, one needs to calculate the radial stress distribution.

Fig. 4.9(3) shows that inside the plastic zone (r<R ) tr where the states of stress are  $\sigma = \sigma > \sigma$ , the radial and tangential stress distributions are (Landanyi, 1974)

$$\sigma_{r} = p_{i}(r/a)^{N-1}$$

$$\sigma_{t} = Np_{i}(r/a)^{N-1}$$

$$4.22$$

and the plastic zone  $(R_{tr})$  is given by:

$$R_{tr} = a[2K_{o}p/(N+1)p_{i}]^{1/(N+1)}$$
4.23

Between the distance (R < r < R ), the radial stress distribution is governed by the stress equilibrium with tangential stress:

$$\sigma_{r} = K_{p} - (K_{p} - \sigma_{r}) (R_{tr}/r)^{2}$$
 4.24

where  $\sigma' = Kp - (N-1)/(N+1)Kp$  at r = R

Substituting  $r = R_{tr}$  and  $\sigma$ ' into Eqn. 4.24 yields the radial stress distribution in the plastic zone ( $R_{tr}$  <  $R_{vr}$ ). Continuity of radial stresses at the elastic zone boundary ( $\sigma = K_{tr}$ ) locates the extent of the plastic zone,  $r_{tr}$  as

$$R_{vr} = a\sqrt{\{K_{o}[(N-1)/(N+1)][2K_{o}p_{o}/((N+1)p_{1})]^{2/(N+1)}/(K_{o}-K_{a})\}}$$
4.25

The wall displacement, u can thus be found in steps similar to Stage (2). The ground convergence curve without inclusion of the gravity effect can be obtained. Examples are given later.

# 4.4.1 Effects of Variations of E and q along with Depth on Shape of Plastic Zone

In previous analyses, Young's modulus (E) and shear strength  $(q_u)$  were assumed to be constant with depth. This assumption, made for purpose of simplicity, may not be valid in practice. Equations governing the extent of the plastic zone for cohesionless and cohesive materials depend on the E and  $q_u$  parameters which may be of variation with depth in reality (i.e., due to confining and consolidated overburden pressures).

#### Cohesionless Soils

For E constant with depth, the extent of the yield zone along the depth can be determined as a function of depth using Eqn. 4.14 and Eqns. 4.19 and 4.25 for Modes A and B, respectively. For Mode A and a constant displacement imposed at the shaft wall, yielding occurs on the vertical (spiral) surfaces and its extent decreases rapidly with the depth. For Mode B, the yielding caused by the vertical-radial stress difference occurs along inclined (conical) surfaces following the Rankine's slip lines and its configuration looks like the truncated cone assumed by Prater (1977) (Fig. D.3).

For Soils with E linearly increasing with depth, the profile of the yield zone will look like those shown in Fig. 4.10. With a constant wall displacement imposed, the magnitude of stress relief with the depth is larger for this case than for E=constant. The induced stress difference in Mode A or B is larger and hence, the extent of the yield zone at depth. The geometry of the yield zone can be determined from Eqns. 4.13 to 4.25 with E=kh, where h is the depth from the surface.

#### Cohesive Soils

applicable to cohesive soils. Equations governing the shape of the yield zones for Modes A ans B are given in Appendix D.2 (Eqn. D.16 and Eqns. D.21 to D.27 respectively). Fig. 4.10 illustrates possible configurations of the plastic zones for different cases. The shape may be cone-shaped or collar-shaped depending on the strength-deformation property variation with depth. That the extent of the plastic zone diminishes at the bottom of the shaft is attributed to the face effect which results in more confinement that reduces displacemnets near the base.

Britto and Kusakabe (1982, 1983) investigated the mechanism of the collapse modes of unsupported axi-symmetric excavations in soft clays theoretically and experimentally. Their findings agree well with the shapes shown in Fig. 4.10.

### 4.5 Comparison of Proposed Solution with FE Analysis

Numerical examples generated by the finite element method (FEM) are used to compare results with those of the newly introduced CCM. As these examples are aimed at filustrating the mechanism of the shaft behaviour and are not designed to simulate any particular case history, typical soil conditions were selected as shown in Table 4.1. Analyses of shaft behaviour in cohesionless and cohesive materials are performed to distinguish two types of response to shaft construction. The chosen numerical values of K oresult in two conditions (K < K and K > K ) and permit two different modes of yielding, i.e., Mode A (tangential-radial stresses) and Mode B (vertical-radial stresses).

## 4.5.1 FE Analysis

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The finite element program SAFE, developed by Chan (1985), was used for this comparison. The soil near the shaft was discretized for the axisymmetric problem by two-dimensional elements. The configuration of the mesh is shown in Fig. D.6. Zero displacement boundaries are assumed at the three boundaries (AB, BC, and CD) for the initial step (i.e., switch on gravity). The boundary AB, representing the wall of a 2m diameter shaft, is allowed to move inward due to the shaft construction. The allowed displacement profile in the FE analyses is also shown in Fig. D.6. A constant displacement is assumed along the shaft depth except near the base.

Several problems were encountered during numerical -simulations. Mohr circles of stress state for cohesionless soil have small diameters the ground surface close to zero pressure (or elements at the ground surface). Small stress relies could cause yielding at these element. For the cohesive soil, tension might be induced in the top layer at which supports were not required. Thus to overcome this problem, a minimal surcharge equivalent to 3m and 4m depth of soil for cohesionless and cohesive materials, respectively, were applied on the ground surface. Some difficulty in obtaining convergence in equilibrium iteration was also experienced. This problem is commonly encountered in FE analysis involving non-linearity, plasticity and limit equilibrium state (Borst and Vermer, 1984; Griffiths and Koutsabeloulis, 1985). The incremental displacement in each step after formation of the plastic zone was u/a=0.01% and the number of iterations to obtain a convergence tolerance of 0.0001 varied from 4 to 10. Hence, it would be very costly to determine conditions at large displacements.

## 4.5.2 Results from Finite Element Analyses Stress Distribution

The radial, tangential and vertical stress contours near the shaft after the last displacement increment are plotted in Figs. 4.11 to 4.13 for three cases. Yielding has occurred, in all cases. In case SM1 (Fig. 4.11), the yielding is induced by Mode B, i.e., due to vertical-radial stress difference (the tangential stress is always the intermediate stress). In contrast, the cases SM2 and CM1 indicate that

Mode A governs the initiation of yielding and the vertical, stress is always the intermediate stress. However, common stress distribution patterns exist for three cases, except at the bottom boundary. Sections at various depths indicate that horizontal arching develops around the shaft openings. The radial stress decreases toward the shaft wall because of the stress relief at the support. This stress relief in radial direction causes the stress redistribution and, results in an increase in tangential stress. This increase in tangential stress occurs in elastic regions. Within the plastic zone, the tangential stress decreases toward the shaft wall. Arching in vertical planes that develop against the, elastic-plastic interfaces are observed in all cases. The plastic zone near the shaft wall tends to displace downward under the gravity effect, and is stabilised by the shear resistance derived from the adjacent elastic zone (the shaft wall being frictionless). The mechanism explains why the vertical stress decreases toward the wall within the plastic zone and experiences a small increase at the elastic-plastic boundary. The vertical arching due to the gravity effect is not very pronounced because the yielding allowed is not excessive.

The stress distributions at one horizontal section (depth h=3.789m) (are plotted in Figs. 4.14 to 4.16 to compare with the results predicted by the CCM solution. This particular section is chosen because it is remote from the influence of the surface and bottom boundaries. Stress distributions at two displacement stages are plotted: first, when yielding at the wall (r=a) is initiated and second,

when much yielding has occurred. A comparison of results in Figs. 4.14 to 4.16 indicates not only that the two methods agree well with repect to the mode of yield initiation at the wall, but also that stress distributions predicted by two methods are close.

Support Pressure-Displacement-Plastic Zone Radius Relationship

Figs. 4.17 to 4.19 present results of the radial pressure-wall displacement-plastic zone radius relationships obtained from the FE analyses (designated by discrete spoints) and the closed-form CCM (designated/by full lines). On the GCC plots (p, -u, ) are included zones corresponding to different stress stages de libed in Mechanisms of Shaft Behaviour (Section 4.4). For SM 1 of yield Mode B, Zone 1 represents the elastic response of the ground and Line E separates the elastic zone from the plastic. Further, displacements in excess of those given by Line E induces yielding by vertical-radial stress difference and this results in a non-linear response of the ground. Zone 2 corresponds to the stress Stage (2) of Fig. 4.9. During Stage (2), yielding  $(\sigma_-\sigma_-)$  has occurred. Further stress relief will cause  $\sigma_{+}$  decrease and become equal to  $\sigma_{+}$ . This state is represented by Line T, and Zone 3 corresponds to Stage (3) of Fig. Fig. 4.9.

Similar reasoning can be applied for cases SM2 and CM1, except that the mode of yielding is different, i.e., induced by tangential-radial stress difference, Mode A. Lines E and T in case CM1 (cohesive) are parallel, instead of

inclined as for cases SM1 and SM2 (cohesionless). This is because yielding in cohesive soil is induced by the constant stress difference given by the compressive strength, instead of stress ratio in cohesionless soil which is dependent of the confining pressure (overburden depth).

On the p - R plots, the plastic zone propagates outward with decreasing p . The results predicted by CCM agree well with those of FE analyses.

## Extent of Plastic Zone with Depth

The extent of the plastic zone with the depth in vertical sections at two displacement steps obtained from two methods (FEM and CCM) are plotted in Fig. 4.20. The results of FE analyses are those of the integration points, and hence show some variations. However, both methods yield consistent results. It is interesting to note that the configurations of the plastic zone around the shaft show the distinct features discussed earlier. In general, the three cases indicate that the extent of the plastic zone increases as displacement increases. The plastic zone in the cohesionless cases not only increases in the radial direction, but also with depth.

A comparison of behaviour in cohesionless and cohesive soils demonstrate that under a constant displacement (depth (except at the shaft bottom) the radius of the plastic zone decreases with depth in cohesionless soils and forms a cone whereas it is constant with depth in cohesive soils, forming a collar (cylinder above a cone). It is also observed from Cases SM1 and SM2 that at the same displacement a larger

plastic zone develop for Case SM1 (K =0.41) than SM2 (K =0.98). This implies that the pressure due to the gravity effect in Case SM1 (K < K ) is more dominant than in Case SM2 (K > K ). However, the support pressure due to horizontal arching is greater in Case SM2 because of higher insitu stress. The resultant support pressure of vertical and horizontal arching in both conditions (K < K ) and (K > K ) varies in each situation. For analysis of these or three cases, the soil parameters are assumed to be constant with depth. It can be expected that the shape of the plastic zone will be different if the soil parameters vary with depth. This phenomenon has been discussed in detail earlier.

#### Pressure Distribution at Wall

From the GCC (Figs. 4.17 to 4.19), the required support pressures for a given displacement are determined and plotted on Figs. 4.21 to 4.23 for two displacement levels, along with results from the FE analyses. Both methods give consistent results except at the bottom boundary. These rigures clearly show that the support pressure is a function of the displacement and initial insitu stress. The support pressure determined from the limit equilibrium methods proposed by Berezantzev (1958) is also included for comparison. Excessive displacement must be allowed in order to obtain these minimum pressures predicted by limit equilibrium methods.

Gravity effects due to the vertical arching within the plastic zone around the shaft are not dominant in these cases because small displacements (u/a=0.3 to 0.56%) are

imposed at the shaft walls. Further movements were not simulated because it is very costly to obtain excessive displacement and also the convergence in iteration processes became very unstable. Nevertheless, the CCM provides an excellent tool to predict the ground pressure at specified displacement levels.

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· grin · i · grin · i · i · i · i · i · i · i · i · i · i	, )	ф ( deд	φ c E (κν/m3) (κν/m3)	E (MPa)	(kN/m3)	>	Š	¥ CL	X Ø
SM.	Perfectly Elastic- Plastic, Associated, Flow Rule, Mohr-Coulomb	<b>6</b> 0 m	0 38 20 (constant with depth)	38 it with	20 depth)	0.290	0.409	0.647	0.295
SM2	as above	33	0 38	38	50	0.400		0.980 0.647	0.295
			(constant with depth)	¥ th	depth)				
CM 1	Perfectly Elastic- Plastic.	0	09	40	20	0.495	0.980	0.980   0.625-	,
•.	Von Mises		(constant with depth)	t with	depth)			0.864	

Table 4.1 Input Data for FE Analyses (Shaft)

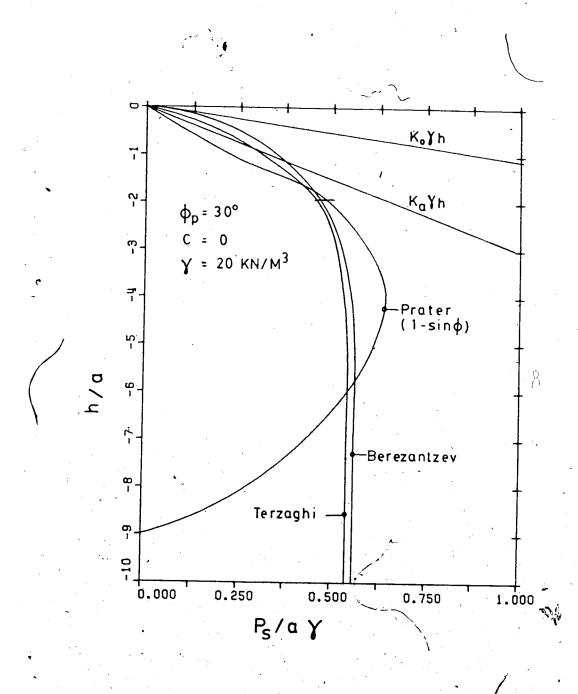


Figure 4.1 Earth Pressure as a function of Depth Ratio (H/a)

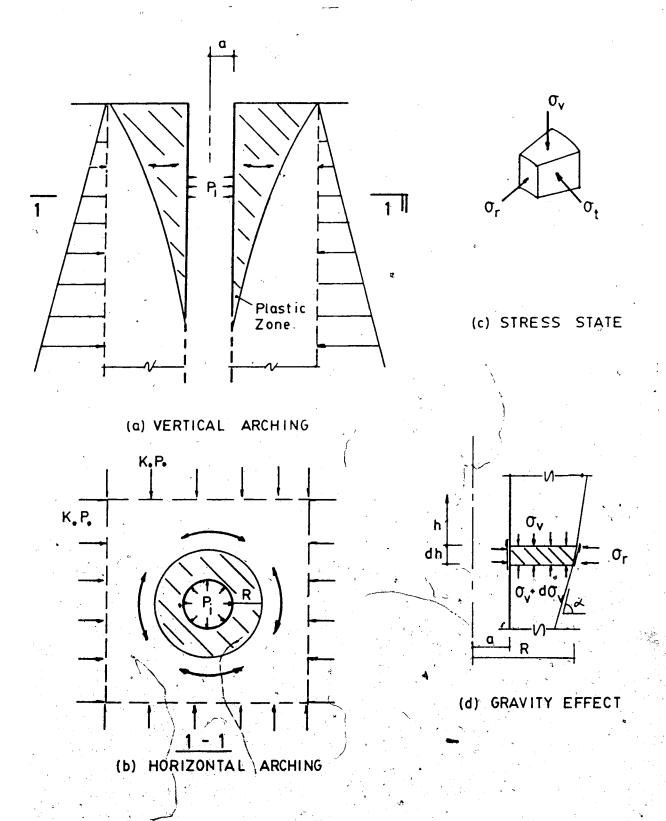
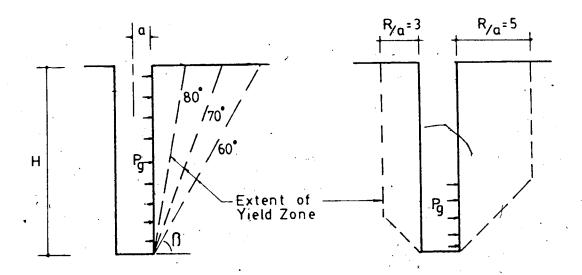


Figure 4.2 Mechanism of Shaft Behaviour



(a) COHESIONLESS

(b) COHESIVE

Figure 4.3 Gravity Effect due to Vertical Arching



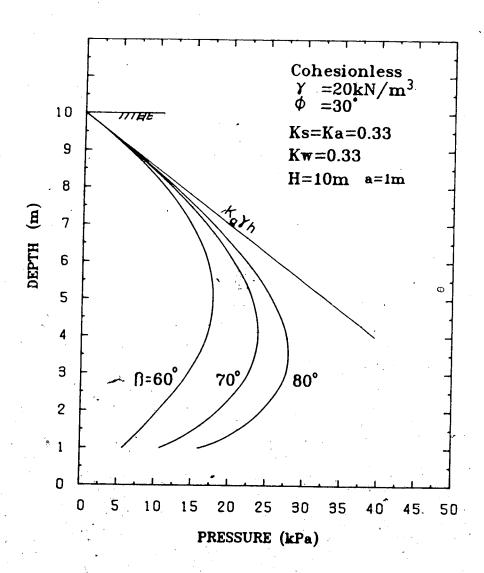
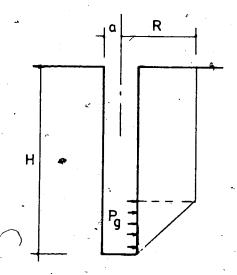


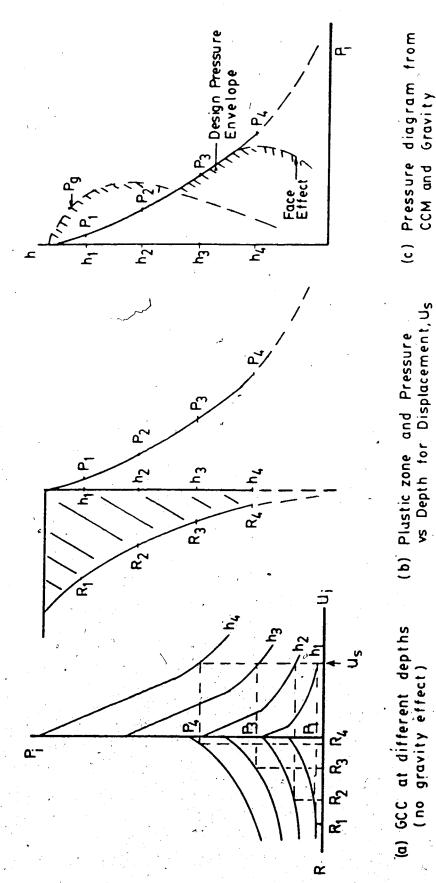
Figure 4.4 Support Pressure due to Gravity Effect (Cohesionless)



H = 10 m a = 1 m C<sub>u</sub>= 40 kPa (A

R(m) Pg(kPa)
3 17
5 183

Figure 4.5 Support Pressure due to Gravity Effect (Cohesive)

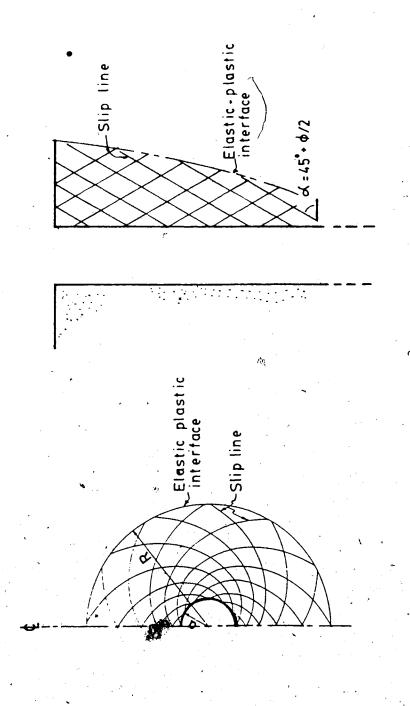


(c) Pressure diagram from CCM and Gravity

(no gravity effect)

Figure 4.6 Shaft Design Approach-CCM with inclusion of

Gravity Effect



i) MODE A . (b) MODE B  $(K_o > K_{cr}, \sigma_t - \sigma_r)$ 



Figure 4.7 Mode A and Mode B

					/
	Stress Distribution	Stress State	Equ	ations	for
	ground opening		R	σ	u
*	(1) Pi Ov Pi or	$\sigma_r$ $\sigma_v$ $\sigma_t$	-	4-2 -4-4	4 13
	(2)	$\sigma_r = \sigma_v \dot{\sigma}_t$	4 14	4·3-4·4 4·21-422	4 15
,	P <sub>I</sub> LR <sub>tr</sub>				
	Or Ot Riv	$\sigma_r$ $\sigma_v \sigma_t$			

Figure 4.8 Stress States (Mode A)

	1			ar e	
	Stress Distribution	Stress State	Equations for		
	around opening		R	σ	u
	(1) O <sub>V</sub>			4.2-4.4	4-16
	- Or	$\sigma_r \sigma_t \sigma_v$			4-10
	(2) O <sub>v</sub>	T .	·.		
4	P Or	$\sigma_r$ $\sigma_t$ $\sigma_v$	4 19	4:3 4:18 4:20	4 15
	+ LR <sub>vr</sub>				. #
	(3) <u>Ov</u>	Ť		4-21	
	Pi dr	σ <sub>t</sub> σ <sub>v</sub>	4-25	4-21	4 15

Figure 4.9 Stress States (Mode B)

J				
	VE	MODE B (K. < K.)	with depth  depth  depth  depth  depth	c
	COHESIVE	MODE A (K.>Kcr).	E. Que Constantswith  E. Que Lorease with  E. Const. with depth	
	ONLESS .	MODE B (K. Kr.)	depth.	
	COHESIO	MODE & (K.>K	E = Constant with de	

Figure 4.10 Configurations of Plastic Zones with De Cohesionless and Cohesive Soils

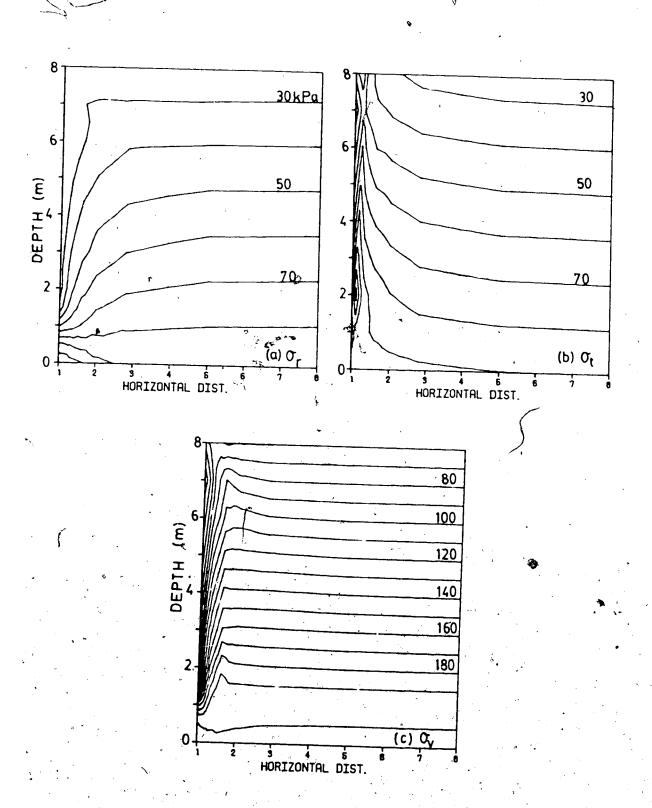


Figure 4.1 P Comparison of Stress Distribution (u/a=0.3%, SM1); (a)  $\sigma$  (b)  $\sigma$  (c)  $\sigma$ 

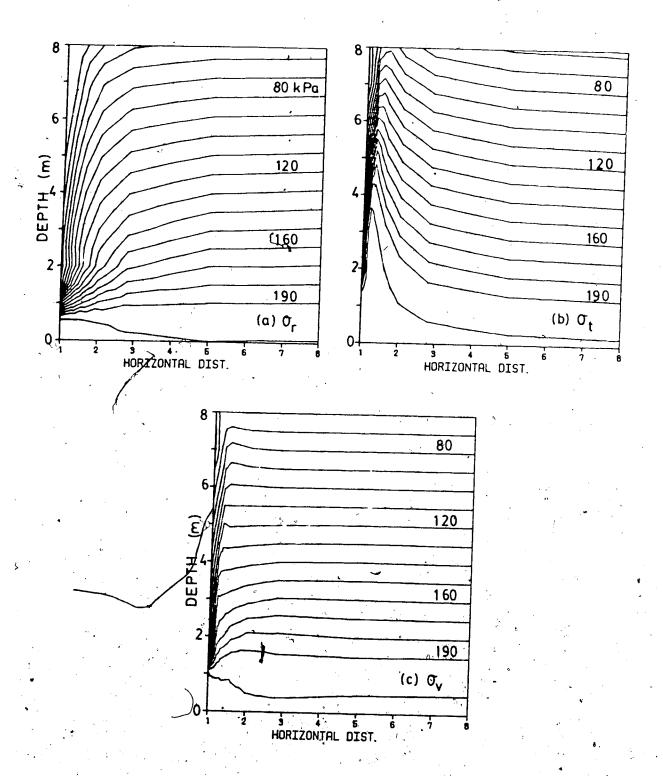


Figure 4.12 Comparison of Stress Distribution (u/a=0.5%, SM2); (a)  $\sigma$  (b)  $\sigma$  (c)  $\sigma$ 

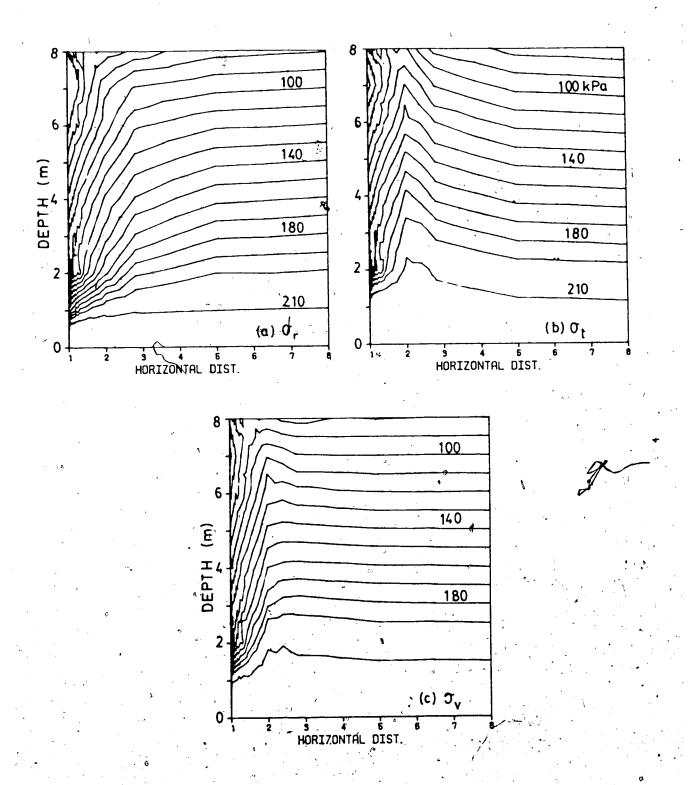


Figure 4.13 Comparison of Stress Distribution (u/a=0.56%, CM1); (a)  $\sigma$  (b)  $\sigma$  (c)  $\sigma$ 

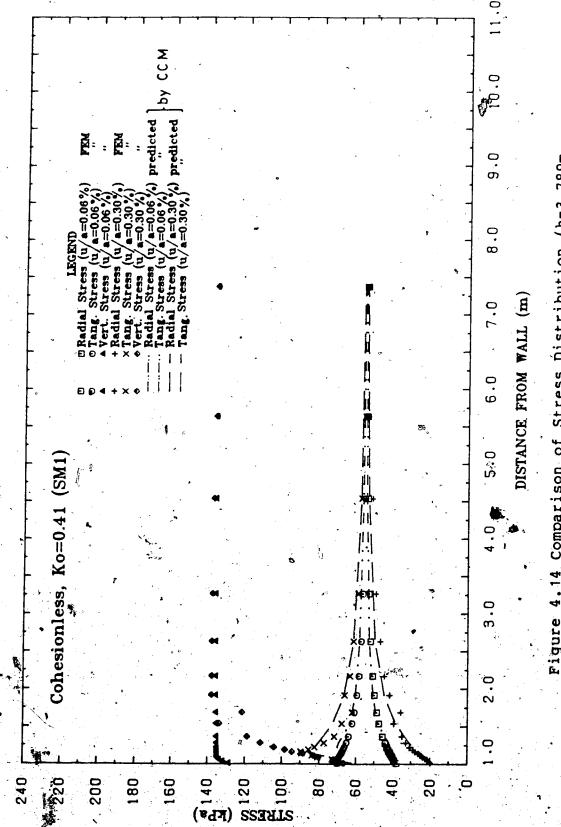


Figure 4.14 Comparison of Stress Distribution (h=3.789m,

SM1)

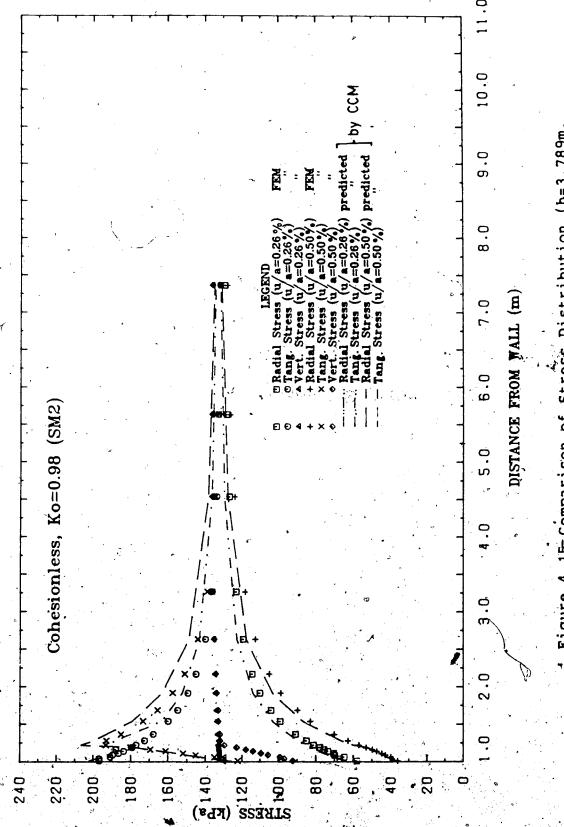
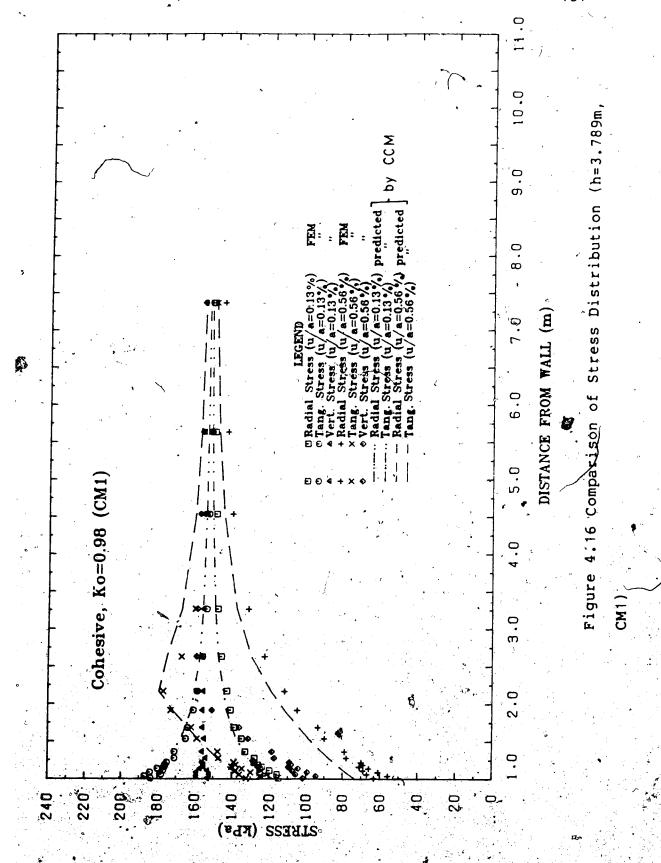


Figure 4.15-Comparison of Stress Distribution (h=3.789m,



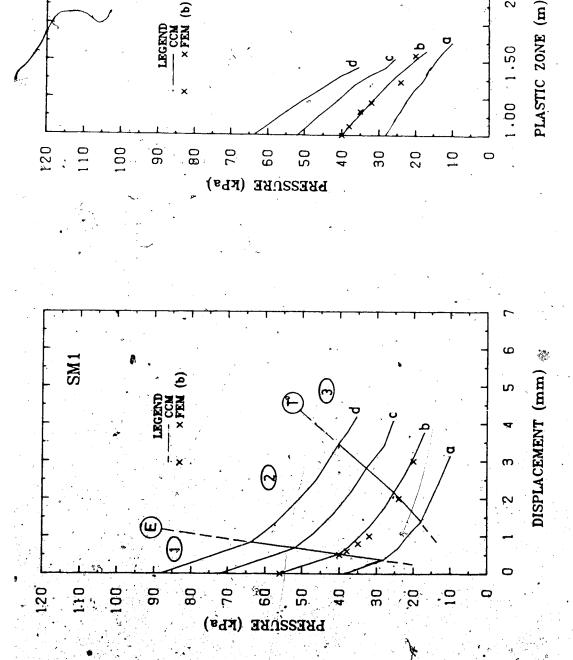


Figure 4.17 Comparison of Pressure-Displacement-Plastic Zone (SM1, Cohesionless, Ko=0.41); a) h=1.789m b) h=3.789m c)

h=5.789m d) h=7.789m

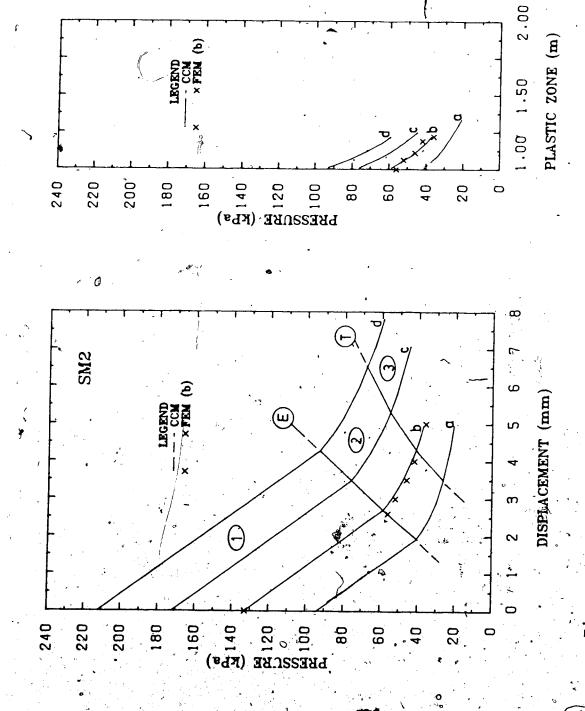


Figure 4.18 Comparison of Pressure-Displacement-Plastic Zone (SM2, Cohesionless, Ko=0.98); a) h=1.789m b) h=3.789m c) h=5.789m d) h=7.789m

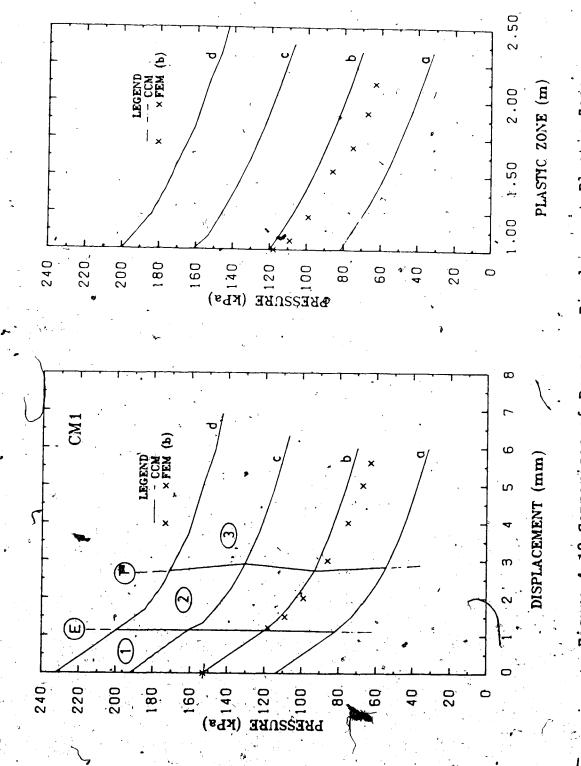


Figure 4.19 Comparison of Pressure-Displacement-Plastic Zone (CM1, Cohesive, Ko=0.98); a) h=1.789m b) h=3.789m c)

h=5.789m d) h=7.789m

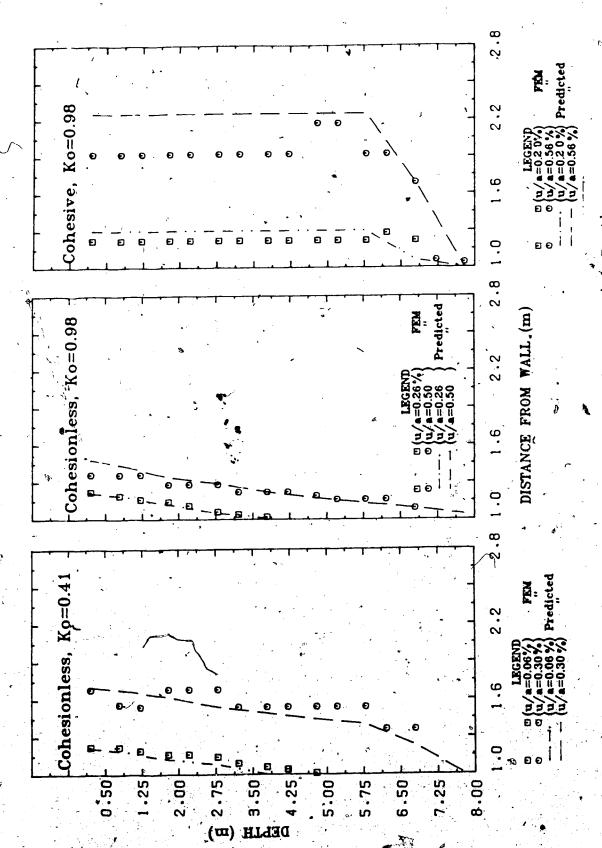


Figure 4.20 Comparison of Plastic Zone with Depth

(SM1, SM2, CM1)

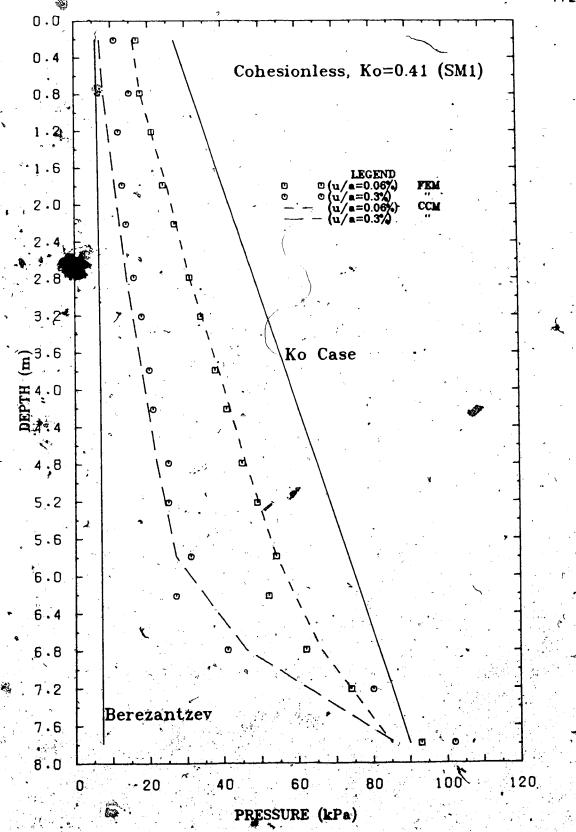


Figure 4.21 Comparison of Pressure with Depth (SM1)

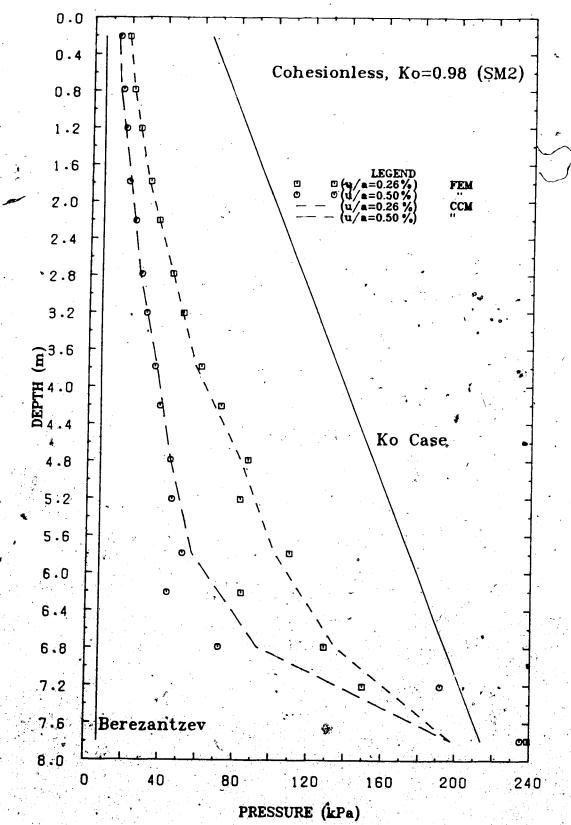


Figure 4.22 Comparison of Pressure with Depth (SM2)

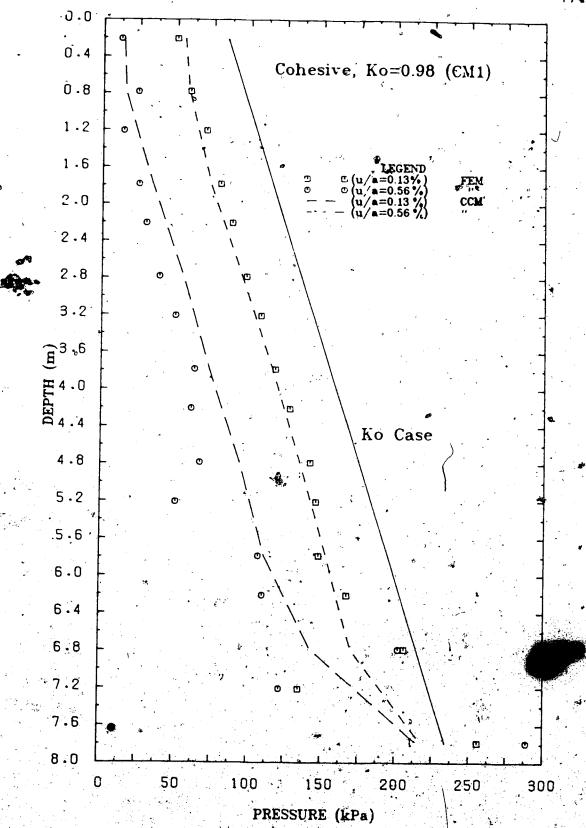


Figure 4.23 Comparison of Pressure with Depth (CM1)

# 5 . SHAFT - CASE STUDY (FIELD MEASUREMENTS)

## 5.1 E. L. Smith Plant - Terwillegar Shaft

## 5.1.1 Introduction

The construction of a shaft south of the E. L. Smith Plant (Terwillegar Shaft) provides an excellent opportunity to collect information for a rational evaluation of currently available design methods and the newly proposed CCM for shafts. Results of field monitoring, data interpretation and analyses of this shaft were presented in detail by Kaiser and Wong (1984). Only data and findings relevant to this thesis are highlighted hereis.

The behaviour of the Terwillegar Shaft excavated in sand and Edmonton Till was studied by an approach integrating field measurement with results from finite element analyses. For this purpose, it was necessary to determine the parameters for the stress-strain model of the soil mass by laboratory testing, to observe the displacement field of the soil near the shaft, to monitor the pressure development on the temporary support and to estimate the original insitu stress. A finite element program SAFE (Chan, 1985) was used to simulate the field conditions for a comparison with field measurements.

# 5.1.2 Project Description

Fig. 5.1 shows the stratigraphy at the location of the Terwillegar Shaft. The lacustrine deposits (above 6m) were found to consist of a brown silty clay with streaks of grey

clay intermixed. The glacial till consisted of a matrix of sand, silt and clay sizes with small peoples and pieces of coal intermixed. Numerous sand lenses, many of which were waterbearing, were encountered in the glacial till. The bedrock consisted mainly of interbedded clay shales and sandstones. Bentonitic layers were generally present in clay shale, however no pure bentonite seams were encountered during excavtion.

The Terwillegar Shaft with a finished diameter of 2.1m was sunk using conventional shaft sinking method. The rate of excavation was about 1.6m per day. Each stage of excavation was followed by the installation of temporary lining composed of corrugated and flanged steel plates which are fully assembled from inside the shaft to form support rings. The diameter of the excavated shaft changed from 3.2m (above depth -14.5m) to 2.4m (below depth -14.5m)

### 5.1.3 Field Instrumentation

Fig. 5.2 shows the layout of the field instrumentation. Due to the axisymmetry of the shaft, all instruments were positioned along the diametric axes. Furthermore, all instruments were located on one side of the shaft to enable easy access for recording during construction, to prevent interruption of the construction procedure and to minimize disturbance or damages to the instruments.

Fig. 5.3 depicts a transverse section (OA) giving the instrument elevations.

Vertical displacements near the surface and at depth were measured using surface settlement points and magnetic

multipoding extensometers. All readings were referred to a fixed reference point, a bench mark at about 60m from the shafts.

Two inclinometers (SI) were installed to detect the horizontal displacements due to shaft sinking.

Six pressure cells were installed on the steel liner to measure the earth pressure. These cells consisted of a flexible circular flat jack, constructed from two steel discs welded at their common periphery.

# 5.2 Analysis of Shaft Performance

The finite element method provides a powerful tool to analyze stress-deformation problems in geotechnical engineering where closed-form solutions are not applicable, but three pieces of information are required: (1) initial insitu stress, (2) material properties and (3) construction sequence causing a change in stress.

The finite element program SAFE used in this project provides options such as analysis for non-linear material models, excavation sequence simulation and liner installation.

The objective of the numerical simulation were:

- To evaluate the performance and effectiveness of the stress-strain models by comparison with field measurements; and
- To obtain numerically a realistic description of the magnititude and distribution of expected lateral pressures along the shaft lining.

Six FE analyses with different combinations of stress-strain models, support stiffnesses and construction sequences were performed. The input parameters for each analysis are listed in Table 5.1.

For the FE analysis, the soil near the shaft was discretized for the axisymmetric problem by two-dimensional elements. The configuration of the mesh is shown in Fig. 5.4. Zero displacement boundary conditions were assumed at three boundaries and zero pressure at the ground surface.

The coefficient of earth pressure at rest (K ) is assumed to be 0.8 (Medeiros, 1979). The initial insitu stress was applied to the finite elements by 'switch-on' gravity. Thus, the horizontal stress is governed by the relationship:  $\sigma_{\rm b} = (\nu/(1-\nu)) h \gamma$ .

FE analyses based on two different types of stress-strain models were performed: linear elasticity and non-linear elasticity (Duncan and Chan, 1970). The soil parameters were obtained from laboratory tests of collected block samples (Kaiser and Wong, 1984).

Two equivalent stiffness values of 30 GPa and 3 GPa were selected for the lining. The first value was calculated from the equivalent stiffness method assuming that the liner plates are fully activated and that no voids exist behind the lining. The second, reduced value, was assumed to approximate the effect of over-excavation and imperfect interaction.

Two cases simulating the excavation sequence and lining installation procedures in the field have been performed.

The first case followed exactly as those in the field and in

the second case, the lining installation was delayed by one excavation step. This delay may account for any over-excavation behind the lining and simulates the delayed activation of the lining.

# 5.2.1 Comparison of Predicted (FEM) and Observed Performance Surface Settlement

Surface settlements are presented in Fig. 5.5. The displacements measured are higher than those predicted by any one of the FE analyses.

The results of analyses LE2/30 and HY2/30 are closest to the measurements (needing a translation by 2 mm). The maximum displacement occurs near the shaft and the settlement decreases gradually away from the shaft. Larger surface settlements near the shaft observed along axes OB and OC are most likely due to the loss of ground in a sand lense, which only dominate the behaviour near the shaft. This lense was not properly simulated by FEM and hence, deviation is expected. It is interesting to note that analyses LE1/3, LE3/30D, HY1/3 and HY3/30D with reduced lining stiffness of delayed liner installation did not produce the greatest settlements, but tended to enlarge the zone of influence.

The results of analyses LE2/30 and HY2/30 are almost equal. This implies that non-linearity is of no significance for this problem because yielding is prohibited due to the relatively low stress levels as compared to soil strength (except in sand).

### Subsurface Settlement

Fig. 5.6a and b compares subsurface settlements at distances of 2.1 and 3.1m from the centre of the shaft (0.5m and 1.5m from the shaft wall). Settlements are generally underestimated by all models. At depths above -6.0m, the analyses LE2/30 and HY2/30 yield closer prediction, and at depths below -6.0m, the analysis HY1/3 gives better correlation with field measurements. Heaving was recorded at a depth -17.5m and a radial distance of 2.1m. This response was not predicted by analysis HY1/3, but by the analyses LE2/30 and HY2/30.

# Horizontar Displacement

Horizontal (radial) displacements at depths monitored by the inclinometers and predicted by the FE analyses are plotted in Fig. 5.7 for comparison.

Except the displacements in the top layer (depth -1.5m) and the displacements due to loss of ground in the sand lenses (at -6.0m), the field measurements fall near the lower bound of the ranges predicted by the FE analyses.

The maximum displacement recorded in the field is 5.0mm (at -10.5m). The calculated tangential strain is 0.24%, and this value is well within the linear portion of the stress-strain curve of Edmonton Till. Hence, it is not surprising that the linear model LE1/3 and LE2/30 (or the non-linear model HY2/30) give the best prediction.

Displacements predicted by the non-linear analyses
HY1/3 and HY3/30D are much higher than those predicted by
the linear analyses and those observed in the field,

especially at greater depth. The delayed lining installation or the reduction in lining stiffness allows excessive soil deformation, and thus induces a decrease in deformation modulus if the soil behaves non-linear. This effect is more significant at greater depths because higher deviatoric stress causes higher deformation. However, this effect is not reflected in the linear model because it does not exhibit any reduction in stiffness with respect to stress levels.

## Pressure on Lining

Results obtained from the FE analyses are presented in Fig. 5.8, together with field measurements.

Little pressure was recorded by the cells at elevations -12.7m and -16.0m. These low pressures indicate that the contact between the soil mass and the cell surface was poor. Soils pressures may have been transmitted to the lining by arching of the soil over the pressure cells. The pressure recorded by these cells are most likely not representative.

Cells at elevation -18.5m give better correlation with results predicted by the FE analyses. The large variation of pressure (two different directions) recorded by these cells may be attributed to various factors that cannot be quantified rationally.

5.3 Comparison of Measurements with Predictions from other Methods

# 5.3.1 Earth Pressure

The pressure as a function of depth calculated by the, methods proposed by Terzaghi (1943), Berezantzev (1958) and Prater (1977) is shown in Fig. 5.9 for a set of assumed soil parameters. These soil parameters are equal to those shown in Fig. 5.10 except the sohesion (c) is assumed to be to zero. The assumption simplifies the calculation, permits a comparison with techniques that are only applicable for cohesionless materials (e.g., Terzaghi, 1943) and should lead to an overestimation of the pressure.

The convergence curves for various depth were constructed using the CCM described in Section 4.2.4 and are shown in Fig. 5.10 together with the assumed soil parameters. Using the measured displacements (u<sub>m</sub>), the formation pressure on the shaft was determined from Fig. 5.10 and plotted on Fig. 5.9. Futhermore, the pressures for assumed displacements of 1/3u<sub>f</sub>, 0.5u<sub>f</sub>, and 0.75u<sub>f</sub> are listed in Table 5.2 and plotted in Fig. 5.9.

The pressure distribution along the shaft, analysed by the FEM model LE2/30 (see Fig. 5.8) is also plotted on Fig. 5.9.

# 5.3.1 Discussion

Fig. 5.9 shows the pressure disbributions predicted by five different methods. The following features are observed:

1. The five methods predict relatively consistent pressure distributions to depths of up to -6m. At greater depths, the limit equilibrium methods yield lower pressure than the CCM and FEM except if very large wall

displacements are permitted (e.g.,  $u = 0.75 u_f$ ). From Fig. 5.9, it follows that displacements at depths below

-6m must be in excess of 0.5 to  $0.75 u_f$  if limit equilibrium methods are adopted for design. Compared to the results of the CCM and FEM, displacements at depths of -6 to -10m are about 0.5u and at depths greater than -10m are well below 0.33u , i.e., excellent ground control.

This clearly shows a dependency between the support pressure and the wall displacement. Large pressures are exerted on the lining when displacements are limited. The support pressure calculated from the limit equilibrium methods is a minimum pressure required to maintain the stability and thus induces large wall displacements.

By plotting the pressure distributions predicted by the limit equilibrium methods on the convergence curves of Fig. 5.10, it is possible to estimate the ultimate wall displacement when limit equilibrium conditions are reached. This is illustrated on one example shown in Fig. 5.11 (same assumptions as for Fig. 5.10). The CCM and FEM give close correspondence with field measurements. The displacements corresponding to the methods proposed by Terzaghi (1943) and Berezantzev's (1958) are 3.5-4.0 times those measured or estimated by the GCC and FEM. Because no pressure is predicted by Prater (1977) for this depth, the displacements are excessive, i.e., about 8 times greater than those of the CCM and FEM.

However, from the above analysis it follows that the method proposed by Prater (1977) is only applicable for depth of less than 8m (h/a<5) on this project. The methods proposed by Terzaghi (1943) and Berezantzev (1958) are acceptable for this project if relative poor ground control is allowed, i.e., the displacements are in excess of about 16mm. The CCM provides fairly similar results as predicted by the limit equilibrium method for h<10 m. At h>10 m, the CCM predicts higher support pressure if displacements are restricted. The measured pressures are generally in support of these findings even though they cannot be considered adequate as proof for the proposed approach. Low pressures were recorded where large shaft wall movements were permitted and higher pressures in areas where movements were more restricted.

 	-	<del></del>					,		· .	
Reference	Duncan	Chang (1970)			(1963)	VKUAC Tests	Katser and Wong (1984)	SPT (field tests)		
Model '8' Hyperbolic	E1=15MPa,V=0.3	# 307 c=0 Rf=0.9 Kur=3.0		E 1=20MPa, V=0 3	# = 25" C= 10.5KP	E1=36MRa, ν=0.44 φ = 22.5 c=31.5KPa	Rf=0.92 Kur=2.5	E i = 120MPa √ = 0.44	E-1= 1GPa .v.	
Model 'A' Linear Elastic	E=15MPa	n.0=∧	. ,	E=20MPa	V=0.3	E=22MPa V=0.44	E=35MPa E=35MPa	E=120MPg V=0.44	E = 1GP a V = 0.44	
	Sand	, i		Çlay	•	<u>-</u>	10 - 10 - 10 - 10 - 10 - 10 - 10 - 10 -	Trans	Clay shale	,
Construct ton Sequence Case	NO. 1	No. 1	2 0 2	No. 1	No. 1	No. 2	luence Henri (Valena end Lana 2002	Tation		
Support stiff (GPa)	<b>m</b>	30	30	'n	30	30	nce (Kalen	ining Installation		
 Strain Model	⋖	⋖	¥	80	₩.	, <b>m</b>	*Construction sequence No. 1 - Cite Conditon	layed 1fn		
 Analysis \Stress- No. Strain Model	LE1/3	LE2/30	1E3/30D	HY 1/3	HY2/30	HY3/30D	onstructi	.2 - Del		•

Table 5.1 Input Data for FE Analyses

E	u / a	r / E O	ઝ	Pressure Pt (KPa)	•		. •
3.2	.001	0018	Measured P1/Po Pt	0.33(Uf) P1/Po P1 F1 52 25	0.5(Uf) P1/P0 P1	0.75(Uf)	
6.4	.0036	. 6100	. 33 . 42	R	35 45	91 27	
9.6	. 0071	.0029	. 48 > 93	41 78		- •	
12.8	.0127	.0028	. 44 112	27 69		•	
16.0	. 0220	.0020	.60 ,192	. 20 64	10 32		
19.2	0000	8008	93 35,8				~

Table 5.2 Convergence Curye Data

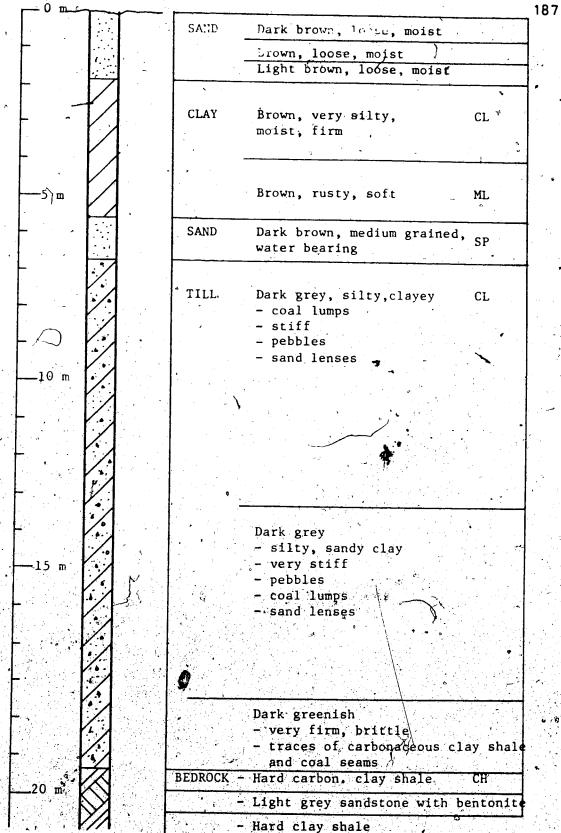


Figure 5.1 Subsurface Stratigraphy (at Terwillegar Shaft)

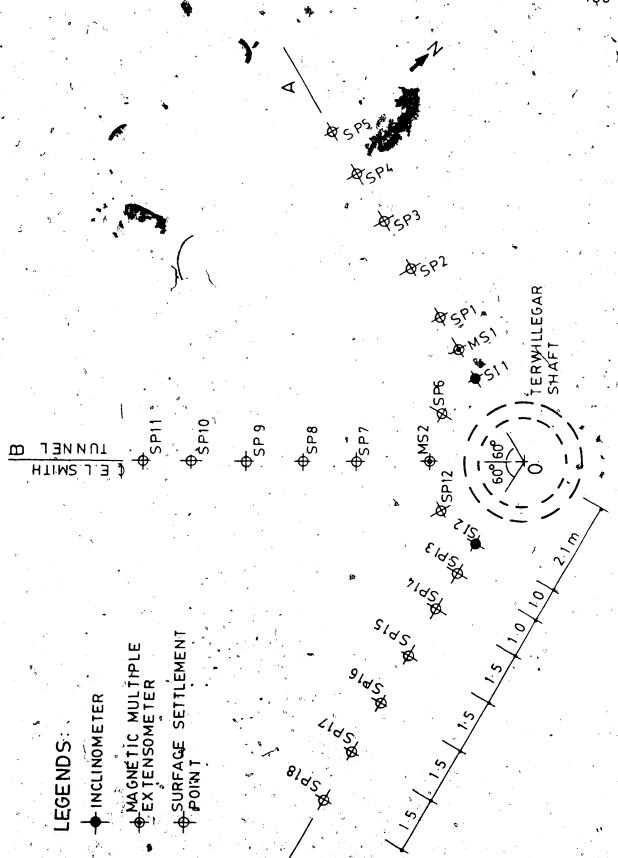


Figure 5.2 Layout of Field Instrumentation)

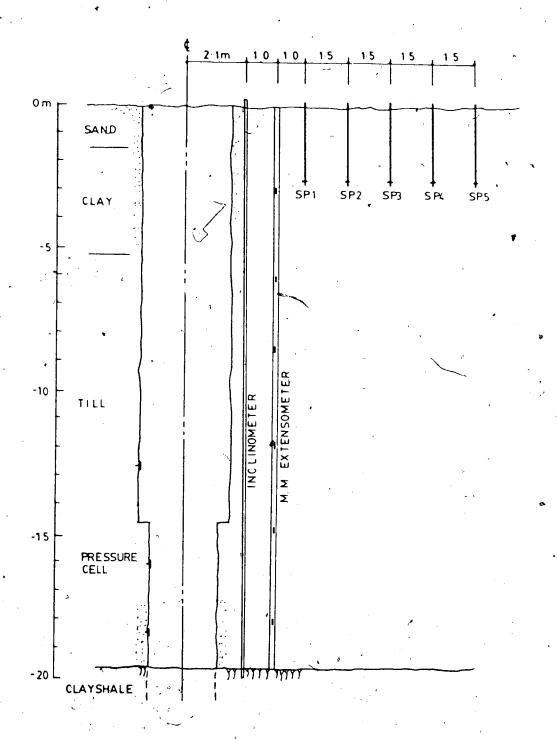


Figure 5.3 Transverse Section showing Shaft Instrumentation (along Axis OA)

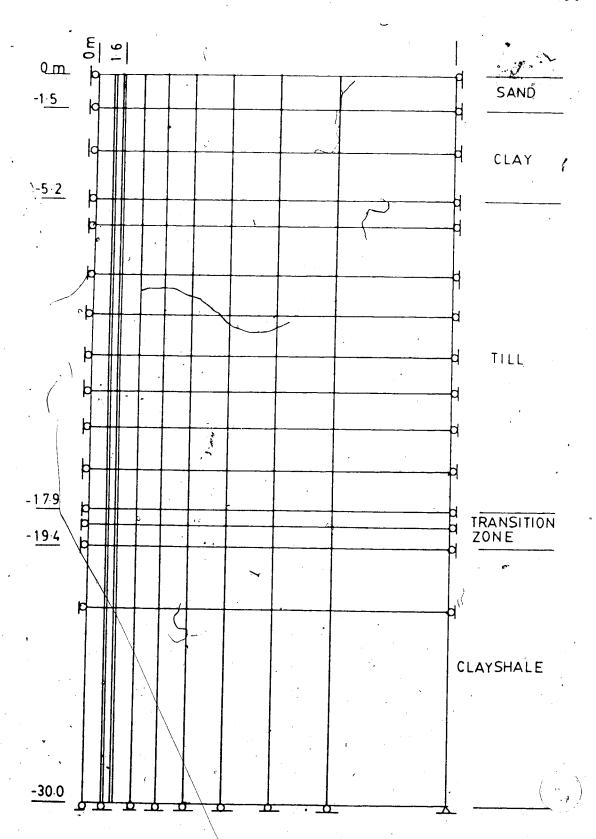


Figure 5.4 Two-Dimensional Mesh for the FE Analyses of the Terwillegar Shaft

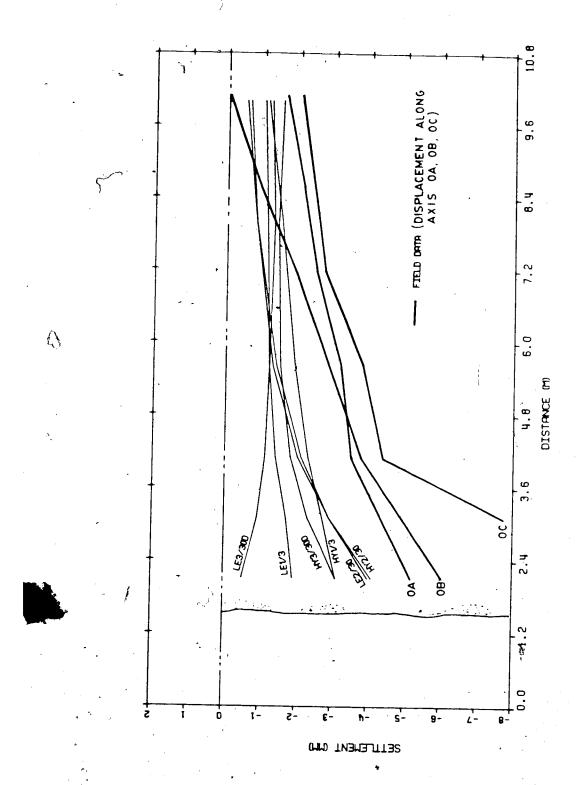


Figure 5.5 Comparison of Surface Settlements

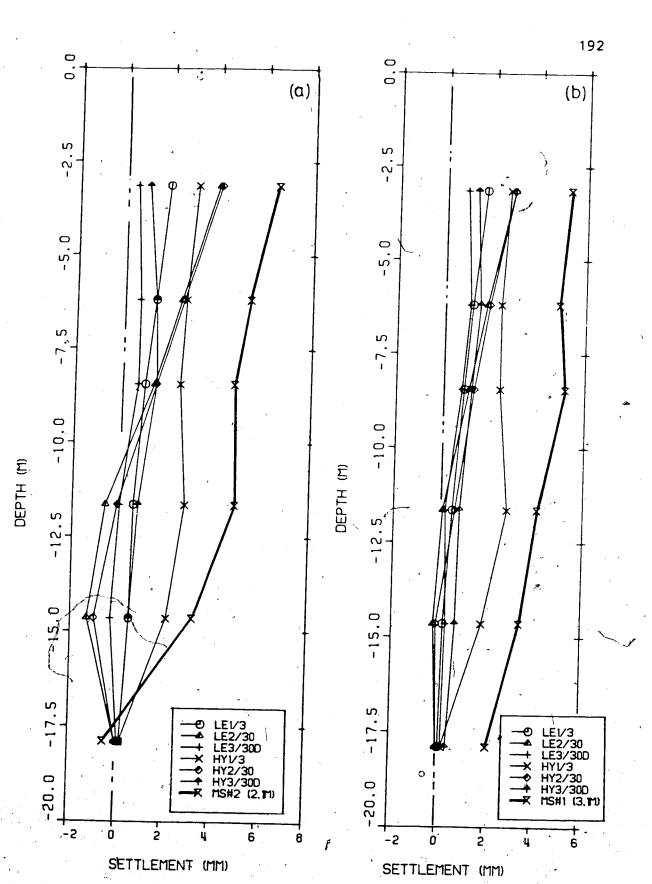


Figure 5.6 Comparison of Subsurface Settlements

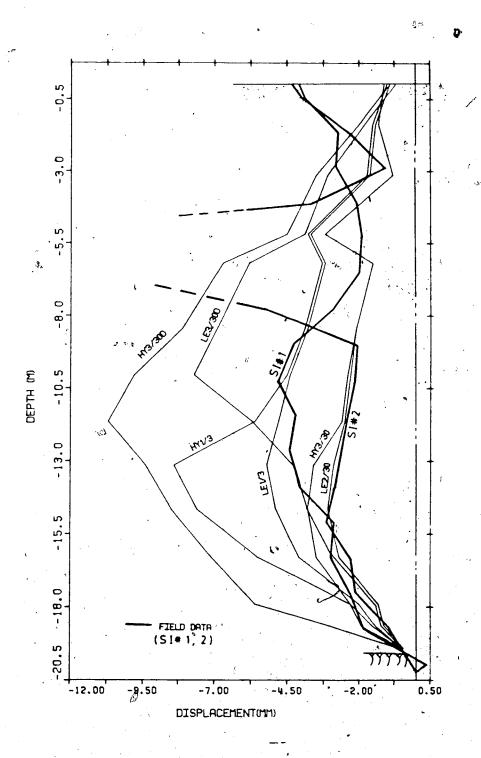


Figure 5.7 Comparison of Horizontal (Radial) Displacements



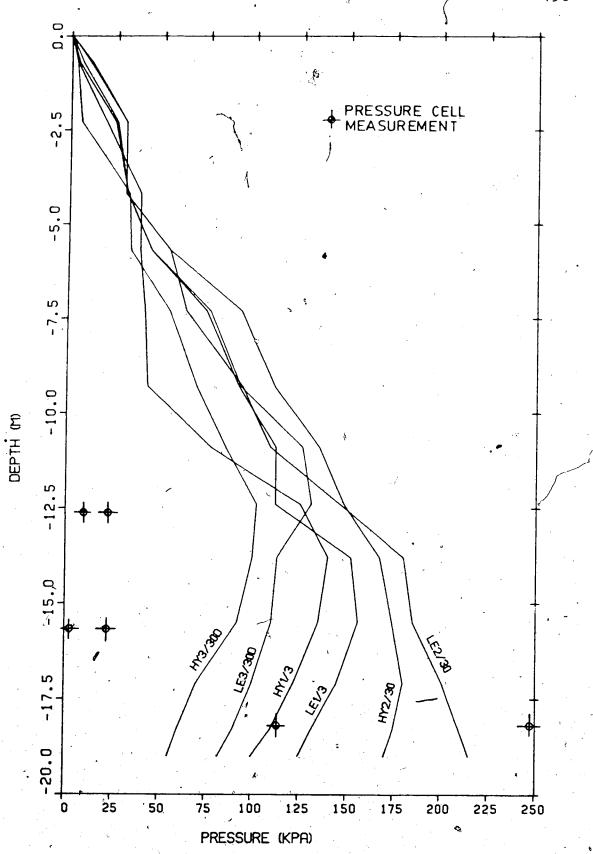


Figure 5.8 Comparison of Soil Pressures

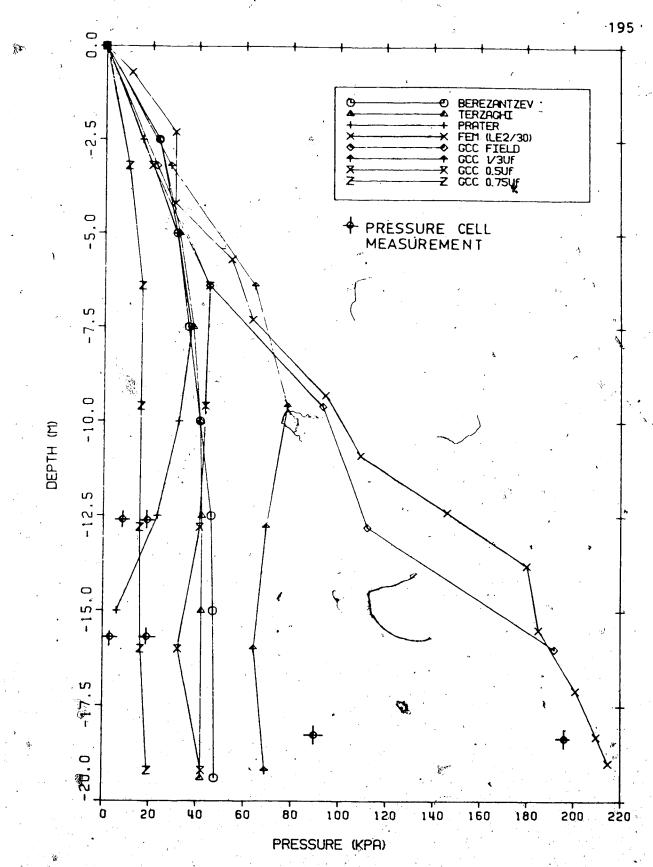


Figure 5.9 Earth Pressures versus Depth

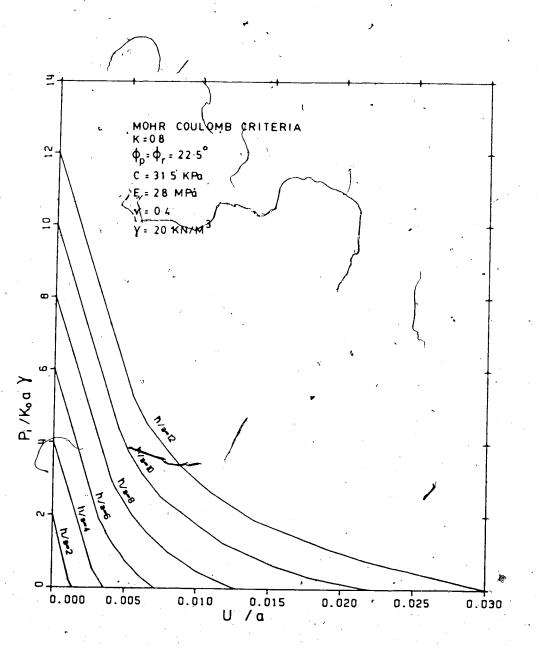


Figure 5.10 Convergence Curves

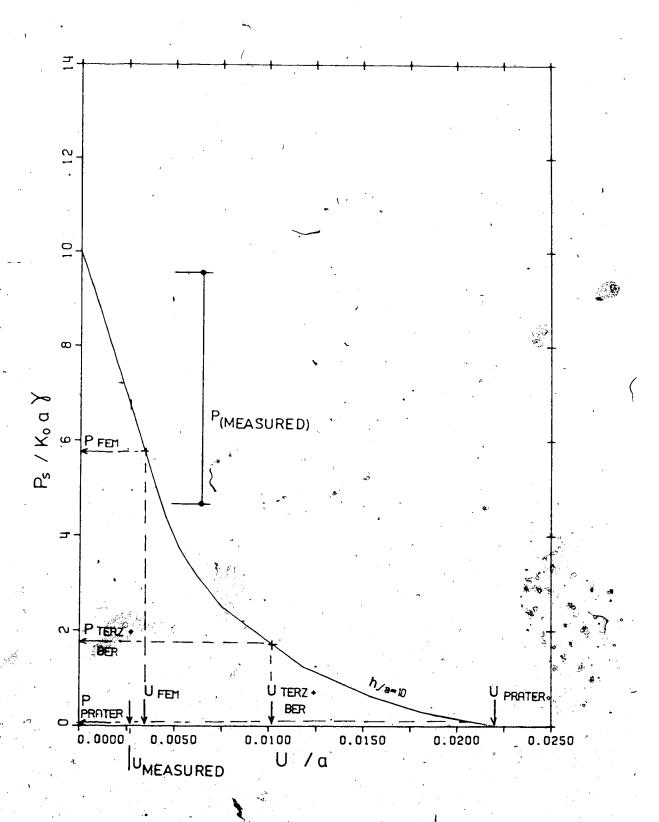


Figure 5.11 Comparison of Shaft Design Methods

# 6 . SHAFT - CASE STUDIES (MODEL TESTS)

## 6.1 Berlin Model Tests (Mode A Yielding)

### 6.1.1 Introduction

Muller-Kirchenbauer et al. (1980) carried out a series of model tests to measure the soil pressure exerted by dry sand on a cylindrical caisson (shaft lining). They observed that the pressures measured differed from those predicted by established theories (e.g., Berezantzev, 1985). Based on their results, they extended Terzaghi's method (1936) for the calculation of the active soil pressure on a plane wall to the axi-symmetrical problem by considering vertical and horizontal arching near the shaft. These two arching effects were described by two lateral pressure parameters  $(\lambda, \lambda)$ which were obtained by fitting the measurements by polynominal functions (note: Terzaghi (1936) assumed  $\lambda_{\star}$  and λ, to be linearly distributed with depth). The active earth pressures on the shaft lining were then calculated by the force equilibrium method on horizontal slices of an assumed failure cone around the shaft.

In the following, a brief description of the testing procedures and certain test results, revelant to this thesis, are given. Their findings will be compared with predictions by the proposed convergence confinement method.

# 6.1.2 Testing Apparatus and Procedure

The shaft model of 100mm in diameter and 650mm in height was fabricated from steel hollow cylindrical

sections. The bottom section was equipped with a sharp-edged shoe to allow easy driving. A recess of several millimeters (ranging from 0 to 5mm) was provided behind the shoe. This recess corresponds to the soil displacement allowed during the excavation. The dimensions of the shaft model are shown in Fig. E.1.

The model was sunk into a test container filled with sands (D=0.2 to 1.0 mm, n = 44%, n = 32.5%) of constant density index. The soil was excavated from inside the shaft. The forces on the shaft lining were measured at different depths. Two types of models (Model I and II) were employed to monitor the pressure. In Model I (Fig. E.2(a)), the axial forces along the shaft height (i.e., frictional forces on the shaft wall) were measured, and the active pressures on the rings were calculated indirectly by assuming a constant wall friction. In Model II (Fig. E.2(b)), three sets of calibrated strain gauges or fiented at 120° are installed in each steel segment to monitor the pressure directly.

Model shafts with different recesses (s=0, 1.0, 1.5, 3.0, 3.5, 5.0mm) were excavated in sand with various density indices ( $I_D$ =20%, 55%, 80%, 90% and 95%) corresponding to unit weights of 16.2, 16.8, 17.6 and 18.2 kN/m³ were measured. No laboratory test results were presented on the strength-deformation properties of the sand.

Earth pressure cells were placed in the soil to measure the radial stresses outside the shaft before and during excavation. The vertical settlement profiles at the surface were also monitored.

#### 6.1.3 Test Results

An example of the active earth pressures (ID=80%) on the shaft lining using Model I is plotted against the depth of the shaft in Fig. 6.1. The larger pressure distribution corresponds to the model test with recess (s=0mm). The pattern of the pressure diagram indicates that some difficulties in measuring pressures existed. Fig. 6.2 shows the effect of recess (s) on the total active force acting around the circumference of the shaft at different depth ratios (h/a) for four density indices. The active force decreases significantly with small increase in recess (s) and reaches a minimum. After a critical recess that depends the density index) the active force starts to increase slightly.

Results of one test using Model II are given in Fig. 6.3. The sand sample had a density index of 80% and the recess was 1.5mm. The active pressures plotted were obtained by averaging readings from three sensors. A further discussion of this figure will follow later.

Fig. 6.4 shows some typical radial stresses measured in the sand outside the shaft before and after excavation. It can be seen that even before excavation, the initial radial stress varied significantly within the sand mass. Assuming the initial vertical stress is equal to the overburden pressure, the calculated K value varies from 1.27 to 1.45. The radial stress decreases due to stress-relief in shaft excavation. Typical settlement profiles near the shaft for I =55% and 95% and different reccesses are plotted in Fig. 6.5. For

I =55% the magnitude of the settlements in general increases with increasing s and the maximum settlement occurs near the shaft wall. For I =90% surface heaving instead of settlement is observed for small s=0.5mm. This response is due to the dilation of dense sand at small straining.

## 6.1.4 Interpretation of Model Test Results by CCM

Two systems (Models I and II) were employed to measure the active earth pressures on the shaft lining using different approaches. The authors reported that tests on Model II gave better and more representative results than those on Model I. For this reason, the results of Model I are treated qualitatively and those of Model II quantitatively.

Model I

The sinking of the shaft consists of pushing the shoe into the soil and removing the soil inside the shaft. The recess provided a gap between the shaft lining and the vertical excavated face. Hence, the soil can displace inward by a variable amount (s) until it interacts with the lining. For the soil lining interaction, the recess can be viewed as the displacement component allowed during sinking in addition to the displacement ahead of the face (assumed to be negligible). Since the active pressure on the shaft lining depends on this displacement, it should be related to the selected recess. Larger recess will cause smaller active pressures. This phenomenon is confirmed by the results shown in Fig. 6.1 where the pressures for s=0mm are about 2.5 to 3

times larger than for s=5 mm. It is interesting to note that the pressures for zero wall displacement is significantly lower than the pressure at rest. This implies that some stress-relief and displacement u took place ahead of the bottom of the shaft.

The pressure-displacement relationship in underground opening is best described by the ground convergence curve. As long as the ground does not yield (elastic), the relationship between displacement and pressure is linear as shown in the inset of Fig. 6.2. Non-linear response will take place due to general plasticity or localized yielding. The shapes of the curves shown in Fig. 6.2 cam be compared with those shown in the inset. These curves actually represent the ground convergence curve for a wall displacement in excess of the displacement u (ahead of th face). The curves of Fig. 6.2 are reproduced from Muller-Kirthenbauer et al. (1980) and are plotted in terms of total force instead of pressure. These curves are not accurate because they are obtained by joining discrete measurements and the initial slope may actually be steeper. The true ground convergence curve can be compained by shifting the measured data horizontally with u . It can be seen that taking displacement as s (neglecting u ) will not induce significant error in pressure prediction when s is large (s >1.0mm) because the slope of the convergence curve is gentler in this stage than the initial slope.

It is most important to note that the total force increases in most cases for large recesses. This can be explained by the gravity effect (see later) due to vertical

arching after a failure mechanism has been established at a critical amount of wall movement (between s=1 and 3 mm). This gravity effect may become dominant as excessive displacements are allowed.

The negative slopes of the initial portion of these curves (Fig. 6.2) increases as the density index decreases. This response is expected because sand of higher density index has a higher Young's modulus. Thus, an equal amount of displacement imposed on sand of high density index will cause relatively more stress-relief than for sand of low density index. Furthermore, the compaction to achieve a dense sand will produce higher initial stress, i.e., the GCC will originate at a high initial stress as can generally be observed from Fig. 6.2.

### Model ÌI

The model test with 80% density index and 1.5mm recess was chosen for a quantitative interpretation and comparison with predictions from the CCM because this test gave typical and representative features of the pressure distribution of all model tests. The measurements were presented earlier in Fig. 6.3 and show: two maxima at H/a= 3 and 10.

In order to calculate the pressure using the CCM and to compare them with the measured results, some necessary parameters had to be assumed. These include the insitu stress state, the stress-strain law and the strength of the soil. No measurements of insitu stresses were reported by Muller-Kirchenbauer et al. (1980) for this model test, but some results from model tests of

 $I_{\rm p}=90\%$  were given as shown in Fig. 6.4. The initial insitu stress coefficient obtained from these measurements were found to vary between 1.27 and 1.45. It is reasonable to assume that the model test ( $I_n=80\%$ ) has a K close to or in excess of 1.0. (The test results with  $I_p=90\%$  was not used here because the pressure distribution was not presented by Muller-Kirchenbauer et al. (1980)). Furthermore, no test results were available to study the stress-strain relationship of the sand. However, the stress-strain curves should be hyperbolic with an initial almost linear portion. From the settlement measurements shown in Fig. 6.5 the sand of high density index (I = 95%) exhibited heave at small straining and must have dilated. This complicates the strain calculation. For this and other reasons the model test with high density index was discarded for data interpretation. Because of the limited information available, the following simplified assumptions were made to obtain a closed-form solution for the calculation of the GCC. The sand is assumed to behave as an elastic perfectly plastic, frictional material that obeys the associated flow rule during plastic straining. The angle of friction of the sand and Poisson's ratio were assumed to be 37° (also used by Muller-Kirchenbauer et al., 1980) and 0.4, respectively and the Young's modulus was asumed to be 44MPa.

The mode of yielding at the shaft wall is due to the tangential-radial stress difference because K is greater than K (=0.62). Hence, horizontal arching can be described cr quantitatively by the GCC of a 2-D plane strain hole-in-plate. The GCC at various depth levels were

calculated, and the pressures were determined by imposing a wall displacement of 1.5mm (or u/a=3.0%) and plotted in Fig. 6.3. The extent of the plastic zone (Fig. 6.6) was also obtained for each elevation and used to calculate the component due to gravity or vertical arching. Assuming a wall friction angle of  $2/3\phi$  (downward shear) as suggested by Muller-Kirchenbauer et al. (1980). Two sets of coefficients on soil to soil and soil to lining wall interaction (K and  ${\tt K}_{\_}$ ) were used and the resulting pressure distributions are shown in Fig. 6.3. These distributions reduced to zero at about h/a=12.5. The pressures exerted on the lining are greater for K = 0.6 and K = 0.5. It was shown in Appendix D that for Mode A yielding (tangential-radial) K and K should be closer to K than K because the vertical stress always remains the intermediate stress. Hence, the pressures due to the gravity effect given by the larger values seem to be more appropriate in this model test. Combination of the pressure envelopes given by the vertical arching (gravity effect) and the horizontal arching yields the pressure distribution on the shaft lining for u/a=3.0%. The resultant envelope compares well with the measured one, except at the bottom of the shaft where the shoe provided a shielding effect attracting the pressure from surrounding soil and reduced the pressure on the lining just above (h/a=13).

From the above analysis, it becomes obvious why the pressure distribution on the shaft lining has two maxima. The upper maximum is caused by the gravity effects which dominate only near the surface due to the shape of the plastic zone and the relatively large wall displacements.

1.

The lower maximum is induced by the steadily increasing pressure with depth and the elevated pressure due to restraint of movement by the shoe (s=0mm). This reduction is only temporary and larger pressures would develop as the shaft advances further.

## 6.2 Cambridge Centrifuge Model Tests (Mode B Yielding

#### ⇒6.2.1 Introduction

A series of centrifuge tests was performed in the Cambridge University Geotechnical Centrifuge to study the behaviour of deep, vertical shafts in dry sand with particular emphasis on the horizontal pressures acting on the shaft lining. A detailed description and the results of a series of these small scale shaft tests in Leighton Buzzard sand are given by Lade et al. (1981). Schofield (1980) presented a detailed description of the Cambridge University Geotechnical Centrifuge and some of the projects performed with it. Here, we shall give a brief description of the testing procedure and results of certain tests which will be needed for comparison with predictions of the CCM.

# 6.2.2 Testing Apparatus and Procedures

A sketch of the model package along with the dimensions and locations of instruments is given in Fig. E.3. The shaft, its lining made of polyetheylene Melinex, was installed within the sand before excavation to simulate artifical ground freezing techniques in construction of deep shafts (physical properties of shaft lining listed in Table

E.1).

Strain gauges bonded to the surface of the shaft tube were employed to measure the radial strain in the tube. Pressure cells were placed in the soil to monitor the vertical, radial, and tangential stresses as shown in Fig. E.1. The vertical settlements across the tub and the vertical movement of the bottom of the shaft were recorded by LVDTs.

Since it was not possible to actually excavate soil in the centrifuge test during flight, the soil in the shaft was substituted with a fluid which could be removed in stages to model the excavation process of the shaft. The ZnCl<sub>2</sub>-solution of density 1.55 g/cm<sup>2</sup> and a parafin oil with a density of 0.765g/cm<sup>2</sup> were used to model the vertical stresses at the shaft bottom and the horizontal stresses inside the shaft, respectively.

#### 6.2.3 Test Results

The centrifuge tests were performed using dry fine Leighton Buzzard Sand 120/200, whose properties are given in Table E.2. Fig. E.4 shows the stress-strain relations obtained from triaxial compression tests over the appropriate range of confining pressures.

Results of three successful tests PL2, PL5 and PL6 will be summarized and discussed. The details of each of the three tests are listed in Table 6.1.

Fig. 6.7 shows a comparison of the radial strains measured in the shaft liner for each test plotted versus non-dimensional depth h/a. The vertical axis on this figure

indicates zero strain corresponding to the unstressed . Melinex liner without earth and fluid pressure. Ideally, the shaft lining should not expand during centrifuge acceleration but even when paraffin oil with much lower density than soil was used did some expansion occur as observed in Fig. 6.7. However, for the three tests discussed here, the initial horizontal earth pressure near the shaft were only slightly higher than the at rest pressure (i.e.,  $K = 1-\sin\phi$ )

Fig. 6.7 indicates that as excavation of the shafts proceeds, the flexible lining moves inward and the formation pressures on the shaft lining increase to the bottom level as shown in Fig. 6.8, a plot of normalized earth pressures versus normalized depth. In all cases do the earth pressure increase near the bottom of the shaft. The earth pressures calculated from Berezantzev's formula are shown for comparison. It appears from this comparison that the earth pressures measured in a flexible shaft are higher than those predicted.

The radial and tangential earth pressures measured with earth pressure cells at level 3 after excavation for Test PL5 are shown in Fig. 6.9. It is clear from the measured pressure distribution that arching develops around the shaft causing the radial stresses to decrease in response to the small radial movement of the shaft wall during excavation.

# 6.2.4 Prediction of Model Test Behaviour by the CCM

One of the essential input parameters for the CCM analysis is the stress-strain relationship of the soil. The

triaxial compression tests on Leighton Buzzard Sand (Fig. E.4) indicate that the stress-strain relations are hyperbolic and dependent on the confining pressure levels. The sand is slightly dilatant at small strains and contractant at large strains.

To simplify the formulation of the ground convergence curve, it is necessary to make the following assumptions on the stress-strain law of sand:

1. The hyperbolic stress-strain relations are approximated by the elastic perfectly plastic models. Because of the non-linearity of the stress-strain curve the Young's moduli E are obtained by averaging the slope on the portion of the curve of the expected strain range. The displacements measured in the model tests were about 0.2%, hence, it seems reasonable to use E at 0.2%. However, it is also realized that there exists an initial strain (0.2 to 0.4%) if the stress difference in K condition is imposed in the isotropic compression triaxial tests (Kaiser and Wong, 1984). To account for this aspect, the E values at 0.5% strain are the appropriate parameters for the elastic perfectly plastic models.

The E-values also depend on the confining pressure or the depth. The confining pressures are isotropic in the triaxial tests whereas they are anisotropic in the model tests. To account for this difference, the first stress variance is used (i.e.,  $I=(\sigma_1+2\sigma_3)/3$ ) for the triaxial tests and  $I=(1+2K_0)\sigma_1/3$  in the model tests).

2. The stress path involved in the triaxial tests

(isotropic consolidated passive compression) is different from that (anisotropic consolidated active compression) involved in the model tests. To make use of the results from the triaxial test it is required to assume the stress path did not exert great influence on the stress-strain-relations.

With the above assumptions the E-values as a function of the confining pressures for 0.2% and 0.5% strains are calculated and plotted in Fig.6.10. The modulus increases with the confining pressure or overburden depth, which can be approximated by linear equations as shown.

Mode of Initial Yielding and Extent of Yielding Zone

For  $\phi=38.3^{\circ}$ , K (=0.38) is less than K (=0.62) and or initial yielding should be induced by the vertical-radial stress difference or Mode B.,

The displacement u to determine how much radial y displacement can be tolerated before yielding around the shaft wall takes place is:

$$u_i = a[(K_0 - K_a)p_0(1+\nu)]/E$$
 6.1,

Table 6.2 lists the calculated displacement u for test y no. PL2, PL5 and PL6, along with intermediate steps involved in calculations. A comparison of the normalized displacement (u /a) with the measured radial strains in the model test (Fig. 6.7) reveals that the soil in the upper zones (PL2, PL5: h/a<1/2, and PL6: h/a<1) is in the elastic state and the soil below has been strained beyond the yield point.

With the measured strain of 0.2% relationships between the extent of the plastic zone and the shaft depth for the model tests are calculated using Eqn. 4.23 or 4.25 and are plotted in Fig. 6.11. The shapes and extent of the predicted plastic zone are cylindrical with depth, which are comparable to the one shown in Fig. 4.40.

## Soil Pressure Distribution along the Shaft Depth

The soil pressures acting on the shaft lining are calculated using Eqns. 4.15 with the assumed strain of 0.2% for the tests, and plotted in Fig. 6.8, along with the measured results for comparison. The observed pressures correspond well with the predicted values. The pressures near the surface lie within the envelopes given K and K whereas below a depth (h/a=1.0) the pressures are much lower than the given by K, i.e., yielding and related stress redistribution take place. The soil pressures on the shaft linings are, however, higher that those predicted from Berezantzev's formula because Berezantzev's method only provides minimum support pressures for conditions of large or excessive deformation.

Higher pressures developed near the shaft bottom. As the vertical confining pressure exerted by the soil above (fluid in model tests) is removed during excavation, the soil beneath the excavation bottom yields. The horizontal pressures are redistributed and attracted to the shaft lining of high stiffness. This result in high pressure at the bottom of the shaft. If the material at the face is strong enough to prevent yielding, the face effect would

reduce, instead of increase the pressures.

## Stress Distribution

From the reduced tangential stress near the shaft wall (Fig. 6.9), it is evident that arching develops in horizontal planes around the shaft.

Lade et al. (1981) calculated the stress distribution in the horizontal plane and compared it with the measured pressures. They found that the general pattern of the stress distribution predicted by the model proposed by Klein and Gerthold (1979) agreed reasonably well with the measured stress distribution but that the extent of the plastic zone was underestimated. This underestimation is due to the fact that Lade et al. (1981) assumed Mode A yielding (tangential-vertial stress). The actual mode is, however, Mode B which is governed by stress difference (vertical-radial) as discussed earlier. The extent of the plastic zone induced by Mode B yielding is given by Eqn. 4.23 or 4.25.

Using the measured pressure on the Melinex tube, one can calculate R. The calculated R of 1.69a or 97mm and the related tangential stress distribution agree better (Fig. 6.9) with the measured and confirms that the CCM with the appropriate mode of yielding is capable of predicting the ground behaviour with sufficient accuracy.

Gravity Effect due to vertical arching inside the yield cone (Fig. 6.9) are not dominant in these model tests because of the relatively small straining.

Table 6.1 Centrifuge Model Test Data

				<b>*</b> -	•			
Test No	Mode4 Dimensions			Prototype -			Soll	Fluid
	d (cm)	H (cm)	N	d (m)	(m)		Density (kN/m3)	, used
PL2	, <del>8</del> .3	24.0	39 , 6⊲	3.3	9.5	5 76	15,50	ZnC12
PL5	<b>/</b> 1.5	<b>45.0</b>	111.8	12.9	50.3	7.80	15.31	2nC12
PLET	7.8	54.0	112.9	8.8	61.0	13.86	15,43	Paraf-
				- *		.~/	•	
d d H - d	adius of Hameter Jepth of Jravity I	of sha f shaft	if t	Ifune				

Table 6.2 Calculation of Displacement  $u_v$ 

				·		
TEST NO.	h/a	h (m)	Po (KPa)	I (KPa)	E (MPa)	u/a (%)
PL2 (H=9.5m, a=1.65m)	<ul><li>♀ 1</li><li>2</li><li>3</li><li>4</li><li>5</li></ul>	1.65 3.30 4.95 6.60 8.25	50.2	30.4 45.6	19.0 23.4 27.7 32.1 36.5	0.032 0.053 0.067 0.077 0.084
PL5 (H=50.3m, a=6.45m)	1 2 3 4 5 6 7	12.90 19.35 25.80 32.25 38.70	292.2 1 389.6 2 487.0 2 584.4 3	118.1 177.1 236.1 285.2	31.6 48.6 65.6 82.6 99.6 116.6 133.6	0.076 0.098 0.109 0.116 0.120 0.123
PL6 (H=61.0m, . a=4.4m)	1 3 5 7 9 11 13	13.2 22.0 30.8 39.6	332.2 2 465.1 2 598.0 3 730.8	20.8 201.3 281.9 362.4	49.4 72.6	0.062 0.099 0.112 0.119 0.123 0.126

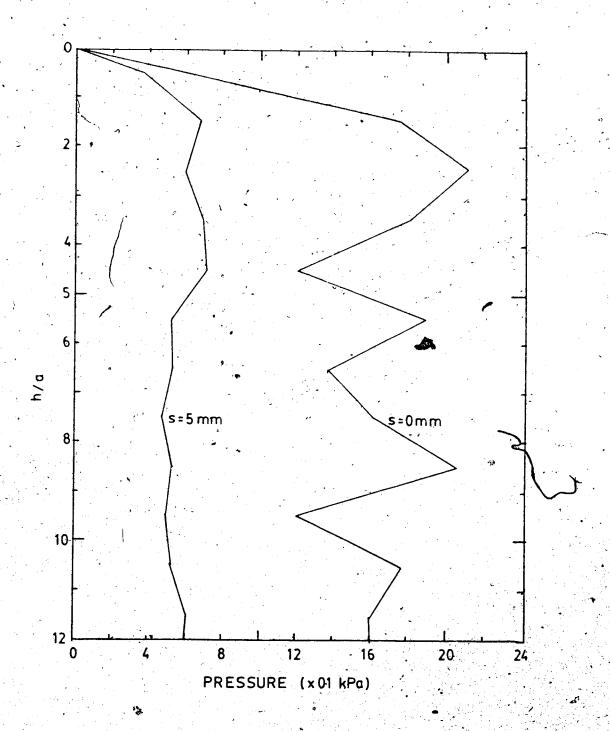
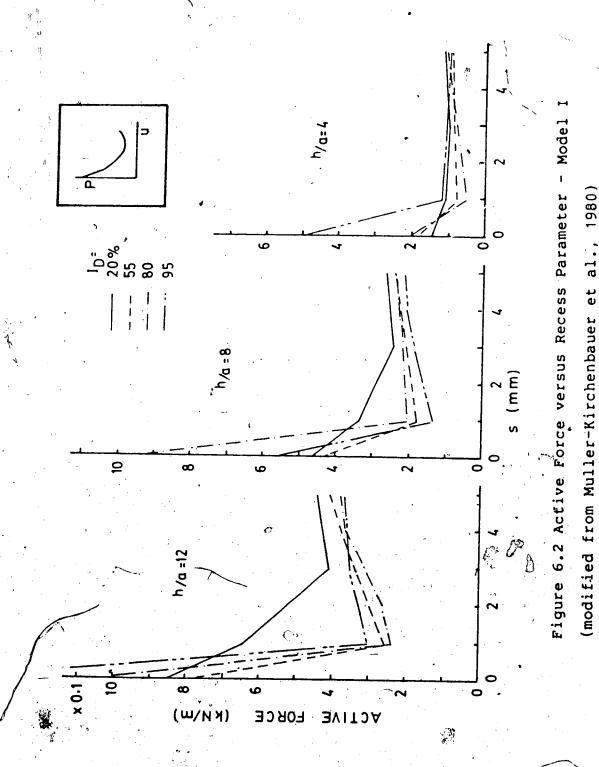


Figure 6.1 Pressure versus Depth & Model I (modified from Muller-Kirchenbaudr et al., 1980)



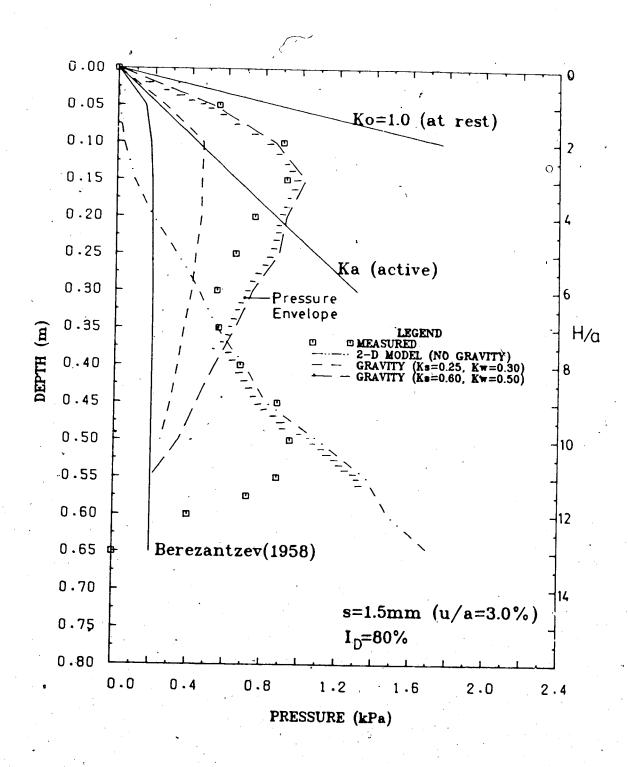


Figure 6.3 Comparison of Active Pressures - Model II

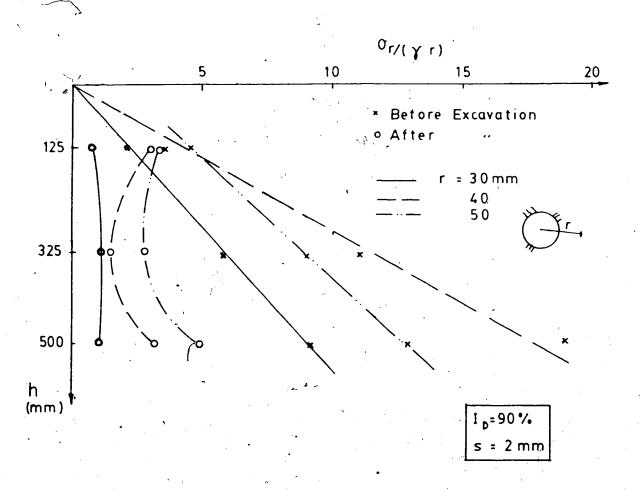


Figure 6.4 Measurement of Stresses Before and After Excavation (modified from Muller-Kirchenbauer et al., 1980)

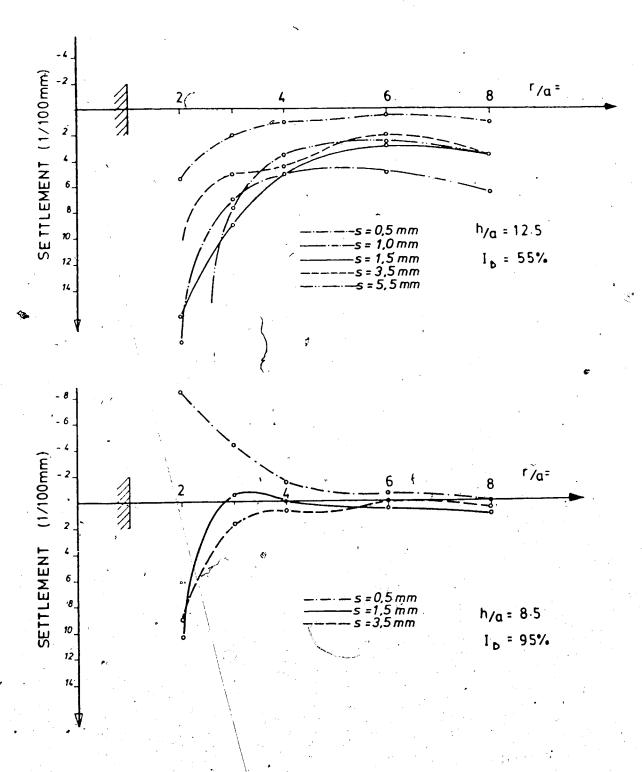


Figure 6.5 Measurement of Surface Settlement (modified from Muller-Kirchenbauer et al., 1980)

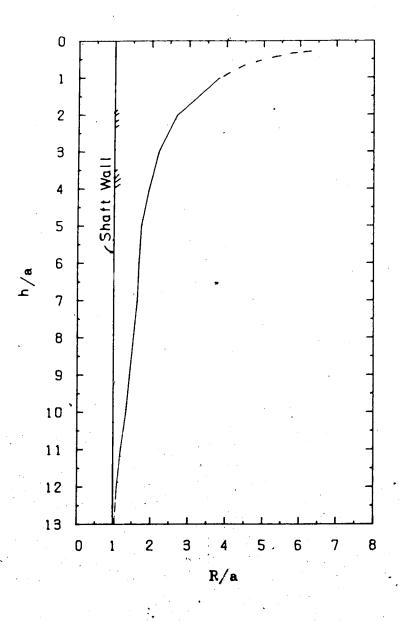


Figure 6.6 Predicted Plastic Zone with Depth

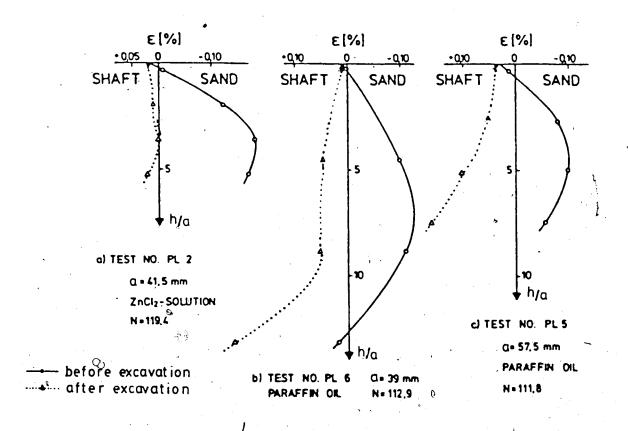


Figure 6.7 Radial Strains in Melinex Shafts versus Depth \*modified from Lade et al., 1981)

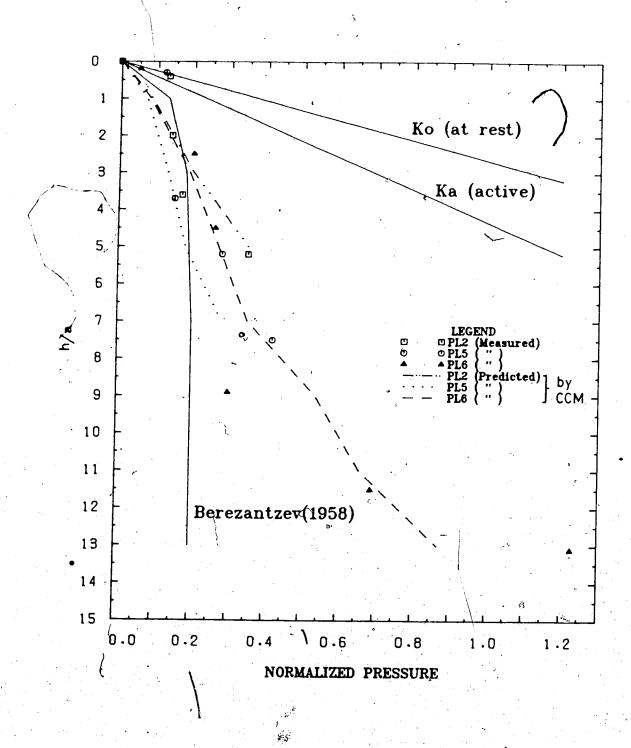


Figure 6.8 Comparison of Pressures with Depth

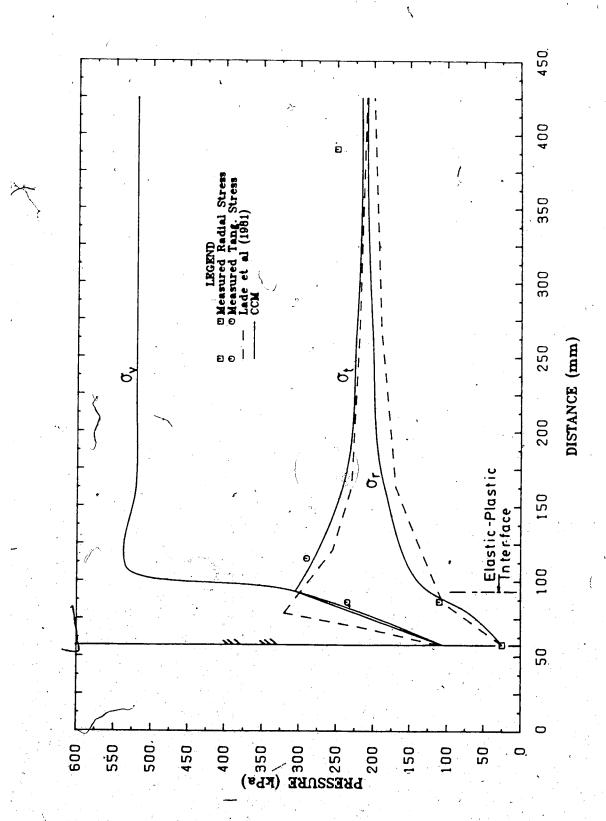


Figure 6.9 Comparison of Stress Distribution

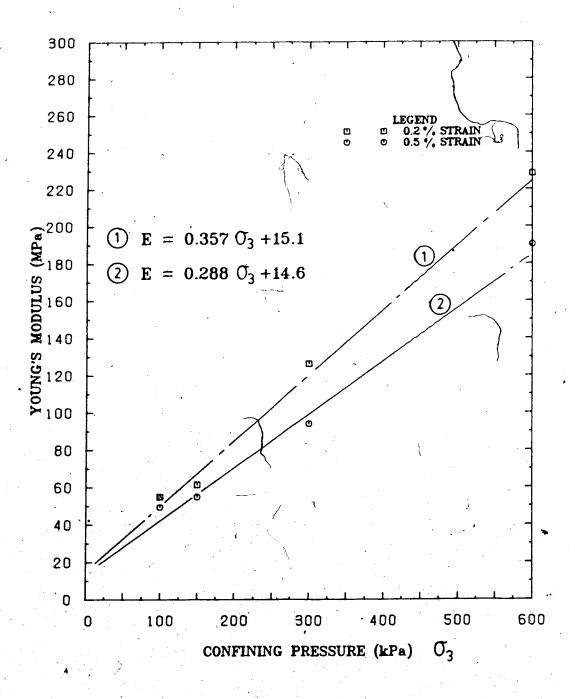


Figure 6.10 Young's Modulus versus Confining Pressure

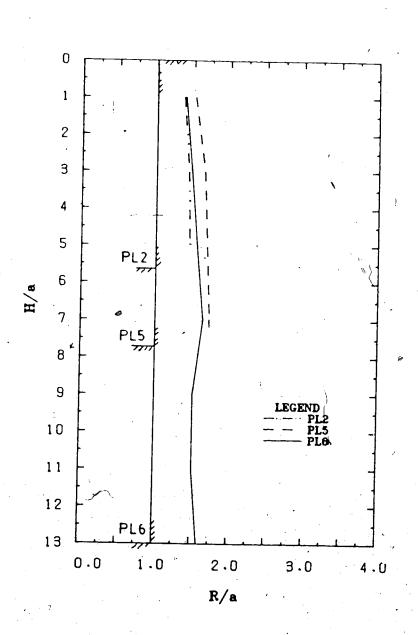


Figure 6.11 Predicted Plastic Zone

#### 7 . SUMMARY AND CONCLUSIONS

## 7.1 Introduction

This thesis examines the behaviour of tunnels and shafts in soils including elastic response, yield and collapse states. These ground responses can be interpreted by the use of the ground convergence curve and analytical techniques are proposed to predict the ground behaviour. These techniques are evaluated by comparison of the predicted with the observed behaviour in model tests and field measurements or by comparison with results from numerical simulations.

The following sections of this chapter summarize the conclusions, practical implications and recommendations for further studies.

## 7.2 Conclusions on Tunnel

#### 7 2 1 Tunnel Behaviour

The objectives of this study are (i) to investigate the mechanisms of ground behaviour near shallow and deep tunnels, (ii) to develop an approach to assess behavioral modes of both shallow and deep tunnels, and (iii) to verify this approach by comparison with results of behavioral mode, support pressure and ground displacement from finite element analyses, field measurements and model tests.

The tunnel is viewed as a two-dimensional plane strain "hole-in-plate" model and the excavation process can be simulated by progressively proportional unloading of the

initial insitu stress at the tunnel periphery. All ground behavioral stages from elastic unloading to collapse due to stress relief are studied by use of continuum mechanics, finite element method and limit state theory. It was found that the mode of behaviour of tunnel and surrounding ground is dictated by the mode of initiation and propagation of yield zone. Four different modes (Modes I to IV) are identified and they are governed by the following parameters: (1) the depth ratio (H/a) relating the free surface boundary, overburden and geometry of the opening, (2) the external boundary stress  $(K_0)$ , (3) the internal boundary stress  $(p_i, \tau_i)$  and (4) the ground strength properties. Mode I with localized yield zone propagation from the shoulders of the tunnel to the free surface and Mode II with continuous yield zone propagation near the opening are only considered in detail in this thesis because they are commonly encountered in soft ground tunnelling at K<sub>0</sub><1.0.

For a given tunnel configuration a critical K value (K ) exists that separates Mode I from Mode II. Mode I will occur for K < K and Mode II for K > K (e.g., K = 0.7 for o cr a tunnel with H/a=5,  $\phi=30^{\circ}$  and c=0). The magnitude of K cr decreases with increasing H/a and  $\phi$ . The existences of Modes I and II have been confirmed by the finite element analyses and model tests (Atkinson et al., 1974; Atkinson and Potts, 1977; Cording et al., 1976). The signifiance of identifying these modes of yielding is two-fold. Firstly, it can verify the mode of yielding assumed in any available numerical models predicting the pressure-deformation around the

opening. The mode of yielding in a proposed model must be consistent with that observed for satisfactory prediction. Secondly, a proper understanding of the mode of yielding provides excellent insight into the tunnel behaviour with aspects of displacement fields including surface settlement profile.

The ground responses of different stages in Modes I and II can be identified in sequence of progressive stress-relief as follows:

- elastic unloading
- 2. yield initiation
- 3. yield zone propagation (yield zone may or may not intersect the free surface)
- 4. kinematic collapse at the  $ro\phi f$
- 5. propagation of roof collapse to the free surface (with states of stress at the roof disturbed)
- From Stage (1) to Stage (4) the support pressure (at roof) decreases at the expense of increased displacements. This support pressure reduces to an optimum value and then may increase beyond Stage (4). At Stage (5) a kinematic collapse mechanism occurs at the roof. The soil block above the tunnel will undergo a sudden collapse and the related settlement will in general be excessive and unacceptable in practice. Hence, one of the practical implications in locating Stage (4) by the use of the ground convergence curve is that the tunnel design engineer can estimate the marginal safety against the collapse state (Stage (4)) by comparing the displacement permitted during construction with the displacement initiating the occurrence of Stage

can be determined by the use of the finate element method until severe problems are confronted during the numerical iteration process. The ground responses of Mode II may be approximated by the "hole-in-plate" model under isotropic stress field with the effect of the surface boundary (supported by studies on two case histories of field measurements) provided that all governing parameters satisfy conditions of Mode II. The responses in the remaining stages (i.e., 4 and 5) are determined using limit state theory.

This method only provides the ultimate support pressures but the related displacements are currently unresolved. However, extrapolation of the ground convergence curve from Stage (2) or (3) to (4) or (5) can permit one to obtain an approximate support pressure-displacement relationship beyond Stages (4).

The surface settlement profile, defined by the maximum surface settlement (S) and the point of inflection (i) is a function of the support pressure (p) or crown displacement (S), depth ratio (H/a) and soil properties. In general, the distance from the centreline to the point of inflection decreases with increasing S at a faster rate for shallow tunnels than for deep tunnels. Modes I and II display several distinctly different and unique features in their settlement profiles because they represent different modes of yielding. If the crown settlement S is small, the extent of the yield sone is small. The vertical settlement profiles of Modes I and II are initially very similar, but the

magnitude of settlement is larger in Mode I than in Mode II. Small displacement occurs at the elastic zone and large plastic straining within the yfeld zone. For excessive S (i.e., the yield zone reaches the surface), Modes I and II exhibit distinct differences in vertical settlement profile above the crown. In Mode I, two localized shear planes develop and the soil block displaces toward the opening as a rigid body. It remains elastic so that the differential displacement between the crown and the surface is small. Hence, the ratio of the surface to crown displacement (S /S ) tends to be close to unity. In Mode II, a plastic zone surrounded by the elastic ground develops around the opening. The elastic displacement is relatively small and most of the straining will occur inside the plastic zone. Thus the settlement ratio S\_/S\_ must be much less than unity. At an equal support stress, the surface settlement in Mode II (K = 1.0) is much smaller than that in Mode I (K =0.5): This difference is attributed to the fact that tangential arching near the opening is higher for K = 1.0 and the resistance against the downward movement of the soil block above the crown is larger. This results in smaller ground settlements.

## 7.2.2 Practical Implications

Based on all of the results presented herein, several observations of practical importance can be made:

1. From the ground convergence concept, the support pressure decreases with increased displacements. This

support pressure reduces to an optimum value (specified by displacement u ) and increases thereafter.

Reduction of lining pressure can be ensured if the displacement permitted during construction (u ) is smaller than u . This displacement u not only corresponds to the optimum support pressure, but also denotes the point of initiation of roof collapse and a state of change in arching action above the crown (from convex upward to inverted downward). This transition in arching signifies an abrupt increase in settlement above the roof. Hence, a comparison between u and con u provides information on the marginal safety against the development of the roof collapse. For good ground control u must be much less that u con

- 2. The initial insitu stress (K<sub>0</sub>) not only governs the mode of yielding but also influences the displacements of soil around the tunnel. There exists a critical K<sub>0</sub> value (K<sub>0</sub>) to differentiate between Mode I (K<sub>0</sub> K<sub>0</sub>) and Mode II (K<sub>0</sub> K<sub>0</sub>) of yielding. The surface settlement in Mode II is much smaller than that in Mode I at an equal support pressure. In addition, Mode I has a narrower settlement trough than Mode II. It implies that the potential for damage due to surface settlement becomes more critical for Mode I than for Mode II. For normally consolidated soils where K<sub>0</sub>=1-sinø, the mode of yielding is generally Mode I. Tunnelling in low K<sub>0</sub> is much more risky as intuitively expected.
- 3. Boundary stresses along the periphery of a tunnel opening exert great influence on the tunnel behaviour

in terms of yield zone development and displacement field. Studies demonstrate that the application of uniform pressure tends to inhibit the propagation of localized yield zones to the ground surface and reduces the surface settlement. This explains why tunnel construction methods such as air-pressurized tunnelling and earth pressure balance shields, provide better ground control. However, this also shows clearly that the approach of applying a constant fictitious support pressure at the tunnel circumference to simulate excavation process does not provide a reasonable solution of ground responses. Ground convergence curves predicted by constant pressure relief are inaccurate. The effects of K and H/a must be evaluated by a correct simulation procedure.

- 4. The German Tunnelling guidelines (Duddeck, 1980 and 1982) for 'shallow' (H/a<5) and 'deep' (H/a>7) tunnels are generally justified both in terms of ground pressure and structural models. These guidelines are reasonably conservative for normally consolidated soils provided that u is always less than u con
- Terzaghi (1943) derived a solution for the roof loading above a tunnel in cohesionless soil. It is shown that Terzaghi's case is only valid for one special case of tunnel behaviour (for a specific crown displacement the horizontal stress to vertical stress coefficient K of soil above the crown is equal to unity). The support pressure may not be on the safe side if u is greater than (or much less than)

u. ic

#### 7.3 Recommendations for Further Studies (Tunnel)

Although the tunnel behaviour has been quite extensively studied in this thesis, there are still many unresolved areas. For example, the boundary between Mode I-2 and Mode I-3 is not well defined. Also, the relationship between the surface and crown settlements was not quantitatively treated and further studies are required to arrive at the exact solution.

Finite element method has proven to be an effective tool in understanding the tunnel behaviour. Techniques in predicting the responses near collapse are definite assets to understand the tunnel behaviour at the limit. Improved techniques in soil behaviour modelling are also necessary for predicting the real and actual tunnel behaviour.

#### 7.4 Conclusions on Shaft

## 7.4.1 Shaft Behaviour

The behaviour of a shaft is a three-dimensional problem in which three stress components ( $\sigma_v$ ,  $\sigma_t$  and  $\sigma_t$ ) must be considered. The excavation process simulated by relief in radial stress component causes both horizontal and vertical archings around the shaft.

Horizontal arching develops due to hoop stresses in a horizontal plane (an increase in tangental stress), and can be quantified in terms of stress-relief (or support pressure) and wall displacement by the use of the

convergence confinement method. This support pressure-wall displacement relationship is dependent on the insitu stress and soil strength-deformation properties.

Vertical arching arises from the formation of a plastic zone of limited extent around the shaft with a tendency to move in vertical direction under the effect of gravity. The support pressure due to vertical arching (or gravity effect) can be determined using the theory of arching. The magnitude and distribution of this support pressure (due to the gravity effect) depends on the extent of the plastic zone, the soil and interface properties. The extent of the plastic zone is governed by the mode of yield initiation (which is dictated by  $K_{\alpha}$ ), the soil strength parameters and wall displacement permitted during construction. An envelope of support pressures resulting from horizontal and vertical archings yields the formation pressure around the shaft for a given wall displacement. In other words, the convergence confinement method with consideration of the gravity effect, which may be viewed as a technique of accounting for the three-dimensional conditions near a shaft, can be used as an analytical approach to predict the true support pressure and the wall displacement around the vertical opening.

From studies on mechanisms of shaft behaviour, it was confirmed that there modes of yield initiation (Modes A, B and C) are possible, and that they are primarily governed by K. Mode A  $(\sigma - \sigma)$  and Mode B  $(\sigma - \sigma)$  are commonly encountered in the field. A critical value K may be used to differentiate between these two modes. The support pressure-wall displacement relationships have been derived

for Modes A and B. They are completely different in terms of yield initiation, yield zone propagation, distribution and magnitude of support pressure, and wall displacement.

Variation of strength-deformation parameters (e.g., q, q, u)
E) with depth also exerts a significant influence on the shape and extent of the plastic zone and, thus, on the distribution and magnitude of the support pressure.
Furthermore, it was also confirmed that the shaft exhibits distinctly different behavioral modes in cohesionless and cohesive soils.

Numerical examples generated by the finite element method (FEM) were used to compare results obtained from the FEM and the newly introduced CCM. These examples included analyses of shaft behaviour of Modes A and B yielding, and responses in purely cohesionless and cohesive materials. Good agreement in terms of mode of yield initiation, extent of yield zone, stress distribution and support stress-wall displacement relationship was found for the two techniques. This suggests that the proposed confinement convergence method provides a valid analytical approach to predict the behaviour of shafts.

The results of three case histories (including one field study and two model tests) have been described and analyzed. The field measurement (Terwillegar Shaft, Edmonton) and the Berlin model tests (Muller-Kirchenbauer et al., 1980) provided cases histories for Mode A yielding whereas the Cambridge Centrifugal model tests (Lade et al., 1981) simulated Mode B of yielding. It was shown that the mode of yielding dominates the behaviour of shaft with

respect to the extent of the yield zone, wall displacement and formation pressure around the shaft. The CCM, which includes important governing factors such as mode of yield initiation, insitu stress state, strength-deformation properties of soils and wall displacement permitted during construction (i.e., effect of construction method), provides a much improved framework of interpretation for the three case histories than conventional techniques.

## 7.4.2 Practical Implications

Based on the studies of shaft behaviour presented herein, the following implications of practical importance can be summarized:

With the formation pressure-wall displacement 1. relationship calculated from the CCM, it is possible not only to rationally assess the required support pressure for a given shaft, but also to evaluate the limit of applicability of the conventional limit equilibrium methods for shaft design. The studies of three case histories showed that the support pressures determined from the limit equilibrium methods proposed by Terzaghi (1943), Berezantzev (1958), and Prater (1977) are much lower than those actually observed in the field or model tests. These conventional methods do not distinguish properly among the various yield mechanisms, and neglect the significance of the shape of the yield zone. They are based on limit equilibrium analyses of one specific assumed failure mode. No strength-deformation relationships are included in

their formulation and full strength mobilization is assumed. Hence, these methods only provide lower limits of support pressure or the pressure required to prevent collapse. These limits are only reached after large ground movements have occurred either due to poor ground control during construction or in cases where large movements have accumulated for other reasons.

2. It was shown that the mode of initiation and propagation of yield zone governs the ground response near an opening. Hence, for the prediction of shaft behaviour it is essential to verify the mode of yielding assumed in a numerical model and to ensure that it is consistent (with the actual mode observed in the field or model test. The convergence confinement method is the only effective technique to evaluate. field observations. Well positioned radial extensometers should be employed to determine the extent of yield zone. Unfortunately, in most field monitoring projects and model tests displacements or stresses are only measured at few locations. Kaiser et al. (1985) demonstrated that data interpretation based 🔩 on one-directional convergence measurements may lead to misinterpretation because many displacement fields and internal stress redistribution mechanisms could cause the same wall movements. Case histories such as Terwillegar shaft and Berlin model test (modified from Muller-Kirchenbauer et al., 1980) have the same limitations. Cambridge model tests (Lade et al., 1981) provides better data on stress distribution but still

insufficient data on displacement field near the shaft opening.

### 7.5 Recommendations for Further Studies (Shaft)

The convergence confinement method has been shown to be a practical analytical method to predict actually expected, support pressures and wall displacements. The results determined from the CCM compare well with field measurements in the experimental shaft study and the results obtained from finite element analyses. However, one limitation of the CCM is the assumption of the two-dimensional plane strain condition. This condition prevails at depth, but not near the ground surface where vertical displacements dominate. Hence, some studies on how the two-dimensional plane strain condition affects the results should be encouraged. Further studies should concentrate on effects of boundary conditions, water table and construction sequence on the shape of the convergence curves.

The ground control involved in the case histories studied were relatively good, and little yielding did occur. In future projects, it might be valuable to allow large ground deformations to observe the collapse mode of a poorTy supported or unsupported shaft.



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## APPENDIX

#### Notations

## Symbols '

The following symbols are used in this thesis:

a = radius of tunnel or shaft

H = depth of tunnel or shaft

h = depth measured from ground surface

y = elevation measured from bottom of shaft

A = area of cross section

R = radius of plastic zone

R = radius of plastic zone induced by ?
tr tangential-radial stress difference (Mode A)

R = radius of plastic zone induced by vertical-radial

vr stress difference (Mode B)

r = radial distance

dh = differential soil element thicknessO

u = radial displacement (convergence) at wall (r=a)

u = radial displacement allowed during construction con

u = radial displacement at support pressure p ic

u = radial displacement at support pressure p ys ys

u = radial displacement at support pressure p fc

 $\sigma$  = radial stress

 $\sigma$  = vertical stress

 $\sigma_{\perp}$  = tangential stress

 $\sigma_1, \sigma_3$  = major and minor principal stress

 $\sigma_{\downarrow}$  = normal stress in x-direction

 $\sigma_{\perp}$  = normal stress in y-direction

```
= shear stress along the periphery of the opening
\tau_{xy}
         = shear stress in xy-plane
         = horizontal stress
σ
h
         = internal support pressure (radial stress) at wall
P;
         = vertical initial insitu stress
         = stress due to gravity effect
\mathbf{p}_{\text{ic}}
         = support stress for initiation of roof collape
           support stress for yield zones propagating to the
P
ys
           support stress at the final limit equilibrium
p<sub>ro</sub>
         = support stress at the time of installation
         = differential (incremental) vertical stress
ďσ
         = soil unit weight
γ
         = density index (relative density) of sand
ם <u>.</u>
         = Poisson's ratio
ν
         = Young's modulus
         = average plastic dilation
         = parameter calculating plastic deformation
        = angle of internal friction of soil
        = \tan^2 (45^\circ + \phi/2)
        = unconfined compression strength of cohesive soil
        = shear strength of cohesive soil
        = wall friction coefficient
        = soil friction coefficient
        = ratio of horizontal to vertical stress
        = K at soil-soil boundary
```

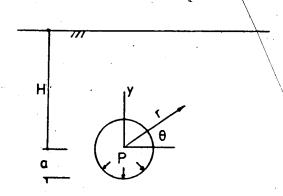
K at soil-wall boundary

```
= K at rest '
К
а
        = K at active state ( = tan^2(45^{\circ}-\phi/2)
Kcr
        = critical K-value to distinguish Mode I (or_A) from
        Mode II (or B) yielding
        = coefficient of tangential arching above the roof
        of the tunnel
S
        = a surface settlement constant
S
S
        = maximum surface settlement
        = vertical displacement at the crown
S
        = point of inflection at the surface settlement
        trough
```

#### APPENDIX B

### Tunnel Ahalyses

B.1 Schmidt's Equations of Stress Distribution around a Tunnel



Radial Stress:

$$\sigma_r = -pq - AH(1-q) + (1-q)(A+B/2-(m\gamma +B/2)q)rsin\theta + BH(1-q)(1-3q)cos2\theta-B/2(1-q)(1+q-4q^2)rsin3\theta$$
B.1

Tangential Stress:

$$\sigma_{t} = pq - AH(1+q) + (A-B/2-m\gamma q - (m\gamma + B/2)q^{2})rsin\theta - BH(1+3q^{2})cos2\theta + B/2(1-q^{2}+4q^{3})rsin3\theta$$
B.2

where:  $q = (a/r)^2$ 

$$A = (1+1/(u-1))^{1/2}$$

$$B = (1-1/(u-1))\gamma/2$$

 $u = 1/\nu$ ,  $\nu$  is Poison's ratio

 $u_1^{\ell} = (u+1)(u-1)/u$   $u_2 = (u+1)(u-2)/u$  $m = -u_2/4u_1$ 

(the negative normal stress is compressive)

Computer programs (TUN1 and TUN2, Appendix F) are written for the above calculations.

B.6

#### B.2 Formulation of Mode II

5

For the bifurcation Mode II, i.e., continuous domain (Fig. B.1a), the stresses within the plastic zone must satisfy the equilibrium equations (it is assumed that the unit weight of soil  $\gamma$  and the direction of gravity is in direction at an angle  $\theta$  to x-axis):

$$d\sigma_{x}/dx + d\tau_{xy}/dy = \gamma \cos\theta$$
 B.3

$$\frac{d\tau}{xy}/dx + \frac{d\sigma}{y}/dy = -\gamma \sin\theta$$
B.4

and the Mohr-Coulomb failure criterion (for cohesionless soil)

$$(\sigma - \sigma)^{2} + 4\tau = (\sigma + \sigma) \sin^{2} \phi$$
 B.5

Substitution of the failure criterion (Eqn. B.5) into the equilibrium equations yields a pair of quasilinear hyperbolic equations (Booker and Davis ,1977). Associated with a set of hyperbolic equations are two families of lines known as a- and  $\beta$ - characteristics. The equations of these characteristics (Fig. B.1b) are:

a lines: 
$$dy/dx=tan(\theta-\mu)$$

 $\beta$  lines: dy/dx=tan( $\theta+\mu$ )

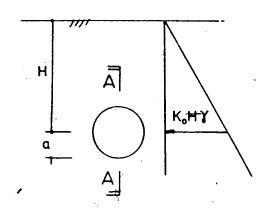
where: 
$$\mu = \pi/4 - \phi/2$$
 (Fig.B.2)

Integration of stress equations along the  $\alpha$ - and  $\beta$ lines with known boundary stress at the periphery of the tunnel can yield the stress distribution within the plastic zone. Booker and Davis (1977) described methods of numerical integration for these hyperbolic stress equations for three classes of problem. The tunnel problem (Fig. B.1a) belongs to the initial value , or Cauchy problem. The boundary stresses along the periphery of the tunnel opening are known, i.e., P1, P22, and etc. The values at successive points, e.g., P<sub>12</sub> are initially estimated from those at P<sub>11</sub> and P22 and solved iteratively by expressing the derivatives in hyperbolic and characteristic equations in the finite differences. The solution is assumed to converge when the sum of the absolute difference among the values in each iteration does not exceed a specified tolerance. Non-convergence may be encountered, implying that the hyperbolic and characteristic equations are not satisfied at that point. Similar procedures can be applied to find values at Points P13, P23 and etc. A computer program (TUN3, Appendix F) entailing all numerical integration processes described above is given herewith (Appendix F).

Within the plastic zone R (Fig. B.1a) where the hyperbolic and characteristic equations are satisfied, the solution will converge. Outside the plastic zone divergenece exists. Thus, using this criterion one can determine under what parameters (H/a,  $\phi$ , p and K) Mode II becomes dominant in a tunnel. Limitations of using the convergence criterion to locate Line i-l are recognized. The postulate that Mode II develops into Mode II-1 is made based on

observed results of model tests, and needs further confirmation. Also, Line i-l separating Mode II from Mode II-1 is obtained from trial cases. The exact solution of this boundary which must satisfy all necessary conditions such as equilibrium, stress continuity across elastic-plastic interface, and boundary conditions (e.g., Detournay; 1983) is not available in this stage.

# B.3 Calculation of Coefficient K



Along A-A, the resultant of horizontal forces are:

initially:  $K_{O}\gamma H(2a)$ 

at  $\eta\%$  support pressure:  $\eta\gamma H(2a)$ 

Assuming even stress redistribution at the roof and floor, the total horizontal force at the roof is given by:

$$1/2K_{O}\gamma(H-a)^{2} + (1-\eta)\gamma Ha$$

which can be expressed in terms of

$$1/2(\rho + \Delta K_r)\gamma (H-a)^2$$

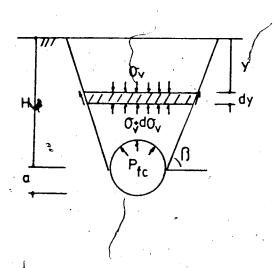
Therefore

$$\Delta K_r = [2K_0(H/a)(1-\eta)]/(H/a-1)^2$$
 $K_r = K_0 + \Delta K_r$ 

B.8

For results see Fig. B.2.

# B.4 Calculation of Tunnel Support Pressure p



Summation of vertical forces for the slice yields

$$-d\sigma_{v} = (1/A_{r})[\gamma dy(A_{r}) - F_{s}]$$
B.9

where:  $A_r = 2a + 2(H-y)/\tan \beta$  $F_s = K_s \sigma_t \tan \phi(dy)$ 

Integration of Eqn. B.9 yields the vertical stress distributions above the roof, i.e., p . A computer program (TUN4, Appendix F) is written for this purpose.

B.5 Analysis of 2-D Plane Strain Hole-in-Plate Problem (Ladanyi, 1974)

Material Characteristics (Linear Coulomb):

Peak (Elastic):

s = c /tan(φ)
p p
p
N = tan²(45°+φ/2)
p
M = (1+(N-1)K p/σ)/(N+1)

 $s_{r} = c_{r}/\tan(\phi_{r})$   $N_{r} = \tan^{2}(45^{\circ}+\phi_{r}/2)/(15^{\circ}+\phi_{r}/2)$ 

Residual (Plastic):

Extent of Plastic Zone:

 $R = a[(K_{0}p + s - M_{0}\sigma)/(p + s)] p$  Radial Displacement at Wall:B.10

$$u_{i} = a\{1-[(1-e_{av})/(1+A_{v})]^{1/2}\}$$
where:  $A_{i} = (2(u_{i}/R) - e_{av})(R/a)^{2}$ 

$$e_{i} = [2(u_{i}/R)(R/a)^{2}]/[((R/a)^{2}-1)(1-1/B)]$$

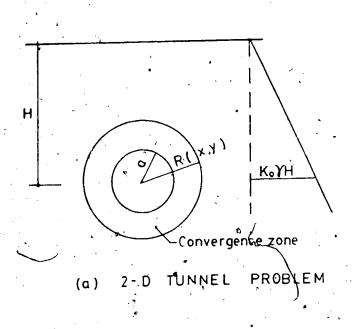
$$e_{i} = 2D[\ln(R/a)] \text{ if } R/a < 1.732$$

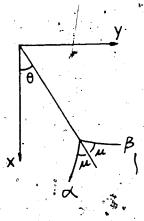
$$e_{i} = 2D[\ln(R/a)] \text{ if } R/a > 1.732$$

$$e_{i} = -\sin\phi_{i}$$

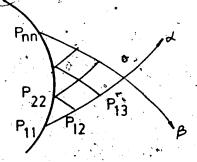
Convergence Curves:

Combining Eqns. B.10 and B.11, p can be expressed in terms of u. A computer program (SFT1, Appendix F) is written to calculate this relationship.





(b) DEFINITION OF



(c) CHARACTERISTICS FOR CAUCHY PROBLEM

Figure B.A. Cauchy Problem

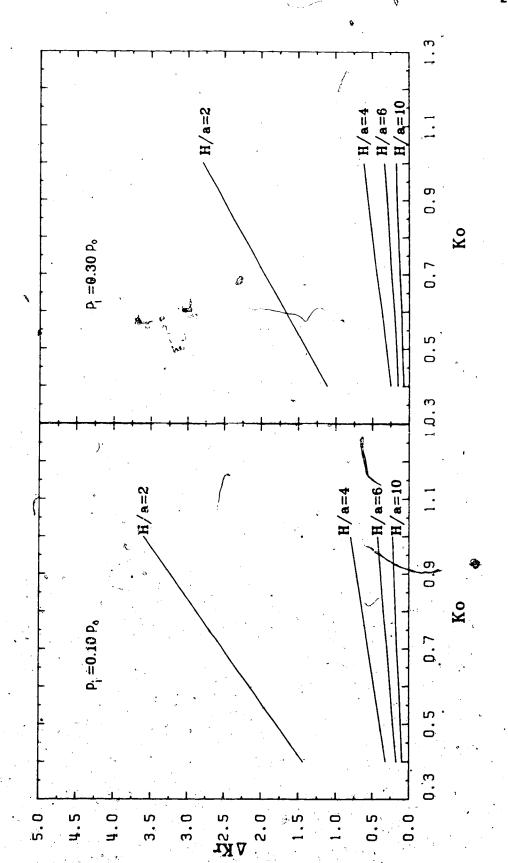


Figure B.2 K\_-values





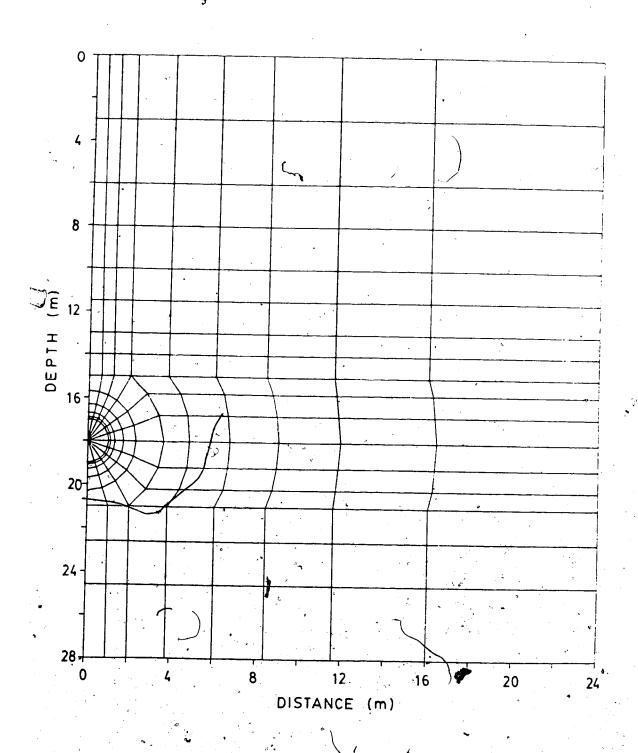


Figure B.3 Finite Element Mesh for Tunnel Analysis

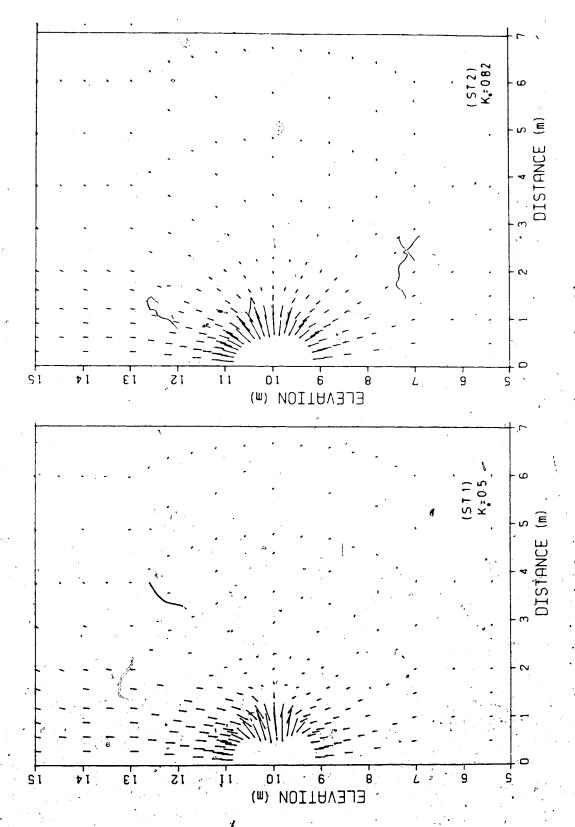
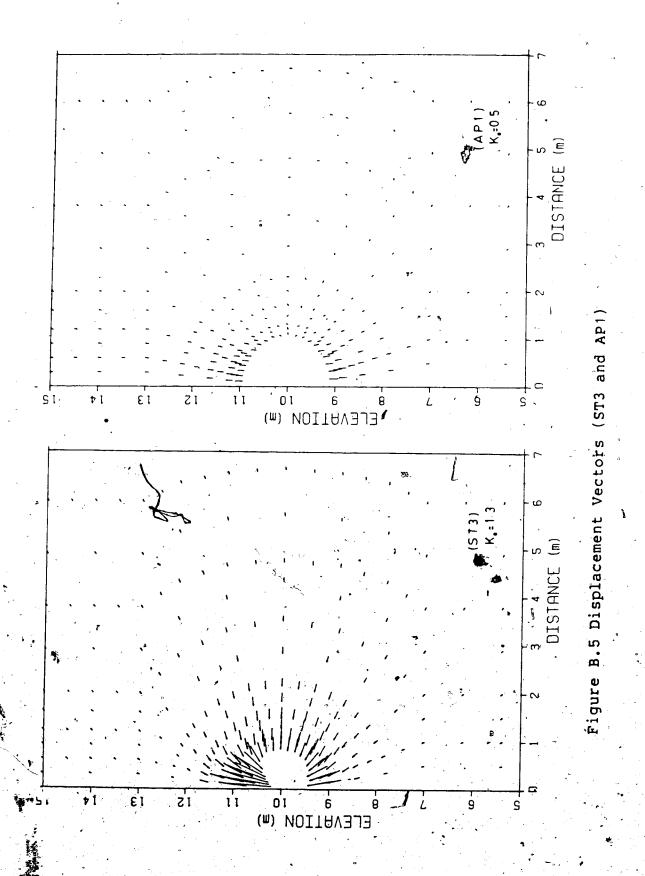
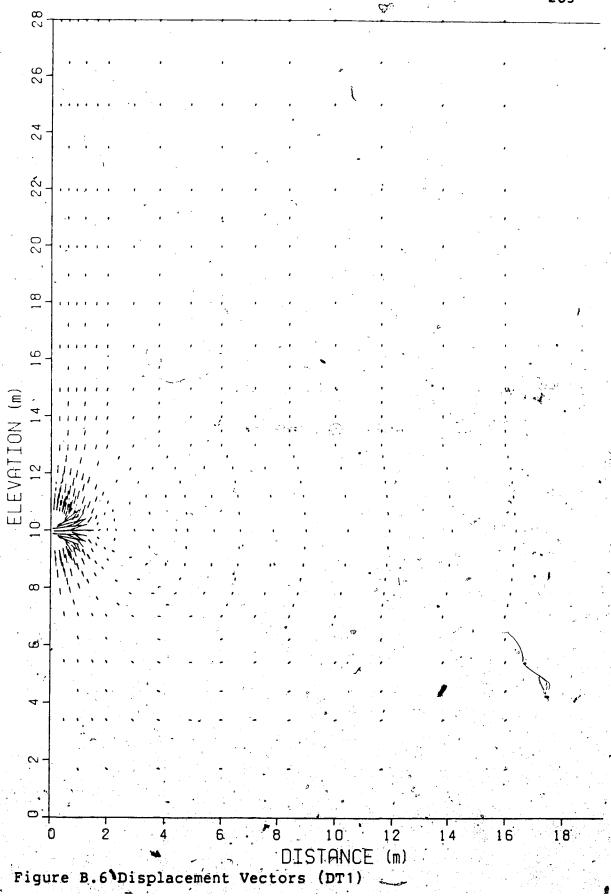
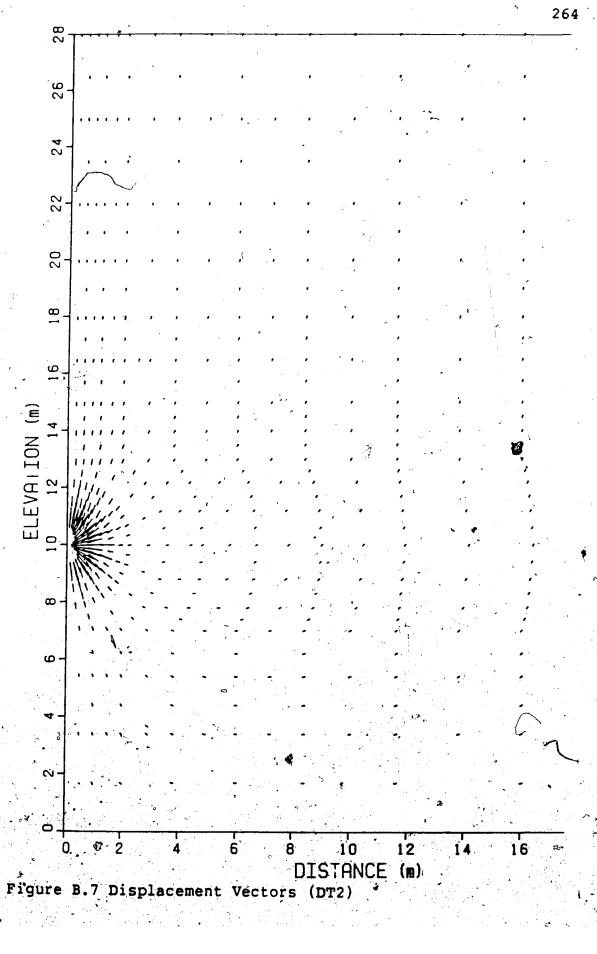


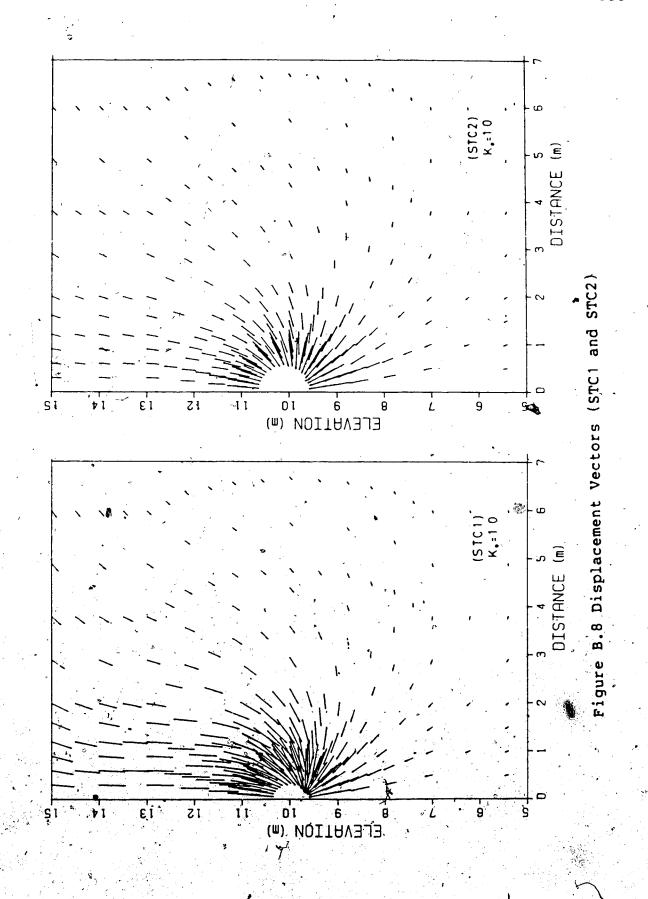
Figure B.4 Displacement Vectors (ST1 and ST2)











## APPENDIX C

Tunnel Case Histories

Table C.1 Tunnel Case Histories - Cohesionless Soils

Cases	H/a	Ss/a (%)	1/a	Reference	No .
London TEL New Cross	4 . 8	1.0	1.8	(1)	(G)
Washington DC Metro	4.5	3.5	1.47	<b>*(1)</b>	(1)
•	6.6	0.2	1.6	1 (1):	(L) ·
Frankfurt Shield	₹3.8	2.2	1.5	(1)	(M)
Brussels Metro	3.2	3.0	1.1	(1)	(т)
Mission Line BART	4.1	1.0	1.6	, (°1)	(1)
Toronto Subway	4.6	2.7	1.0	(1)	(2)
•	3.9	8.7	0.7	(i) ,	(15a)
•	5 . 1	11.1	Q.9	(ဌာ	(156)
Ayrshire JD Scheme	4.3	9 <b>3</b>	1.0	(2)	5. #
Warrington Sewer	4.7	1.4	1 . 2	(2)	_ •
,	4.7	1.0	1 . 4,	(2)	-
,	4.2	2.9.	0.8	(2)	-
- • ·	4.2	1.3	1.2	(2)	~
	-4 , 1	1.B	1.1	(2)	
•	. 4.7	Q.2	2.7	.(2)	-
WNTDC Sewer	2.6	4.3	- 1.3	(2).	
	5.0	1.1	1.4	(2)	-
H	3.6	O.B.	· 0.9	(2)	-
	3.76	2.8	1.0	(2)	-
j	3.6	0.4	1.3	(2)	<i>}</i> -
North West A Sewer	. 8.4	2.8	3.2	(2)	\
Northumbrian Sewer	7.5	4.7	4.2	(2)	
LRT Edmonton	3.4	0.3	0.8	(3)	· -
ED Tunnel Edmonton	18.8	0.9	12.5	(4)	- ·
Model Test (Cambridge)	2.0	6.0	0.7	(5)	
	3.0	4.1	0.8	(5)	• •
	4.0	3.5	1.2	(5),	·•·
Model Test (Illinois)	4.9	4.61	2.2	(6)	
•	5.7	4.9	3.6	(6)	<b>-</b>
	5.7	6.	2.0	(6)	
Ref.: (1) Attewell (19	) 1781 (2	) 2/1-1	•	the second second	

<sup>(1)</sup> Attewell (1978), (2) O'Reilly and New (1981) (3) Branco (1981), (4) El-Nahhas (1979) (5) Atkinson and Potts (1977), (6) Conding et a

Table C.2 Tunnel Case Histories - Cohesive Soils

l

Cases	H/a	\$s/a (%)	) j/a	Cu (kPa)	Referen	ice No:
'Model Tests (Cambridge	2.5	8.4	1.c	26,	(5)	-
•	1.8	8.8	0 -7	26	<b>\</b> (5)	-
Washington Fletro	4.6	4.8	* 1.4	75	7 (1)	(J)
, <b>6</b>	3.6	B.8	0.6	75	(1)	(K)
Model Tests (Cambridge)	3.5	7.7	1.4	<b>2</b> 6	(5)	-
•	, 4, 0	7 . 4	2.0	26	(5)	-
Toronto Subway	4.9	1.4	1,4	67	(1)	(12)
Chicago D-5	3.9	1.3	0.9	67	(1)	(13)
Sewerage, Belfast	316	1.2	2.0	10	(2)	
BART San Francisco	3.3	≒∠1,5	1.8	10 .	(2)	_
Sutton Sewerage.	3.8	0.4	2.3	90	(2)	
Bristol Sewerage	3.5	1.2	2.0	/ 18	(2)	-
IWA Seweage Scheme	7.5	0.8	3.7、	731	(1)	(-B)
lodel Tests (Camporidge)	7.2	5.3	3 7	26	(5)	•
	6.0	6.1	3 <sup>*</sup> 0	26 .	′(5)	_
yoto, Tokyo Subway	6.4	0.3	15.5		(1)	(9)
ART San Francisco	6,5	1.7	.3.2	77,	(1)	(10)
WA Sewerage	7.5	О.В	3.9	73	(2)	
uttón Sewerage	6.4	0.9	4.0	90	(2)	_
ockton Council	10.0	6.9	5.5	50	(2)	•
• ,	9.3	8.9	5.8	50	(2)	. <u>_</u>
ashington Metro	4.6		2.1	155-550		(N)
• • • • • • • • • • • • • • • • • • •	.2:4	0:2	1.2	72-275	(1)	(R)
ınn ing <u>T</u> unne i	5.5 -	0.3			(2)	. <u>.</u>
ondon Tránsport	6.8	0.2		230	(2)	
ford Trunck Sewer	•	0.2				
	14 1	0.3		270	(1)	(A)
		• •	<b>~</b> ,		(1)	(E)
		0.2	•	230	(1)+	(E) .
	12.0	0.4		1	(1)	<b>13</b>
		1.		354	(0)	(11)
CCOII SAMAII	19.2	0.4.1	1.2	180	(2)	-

<sup>(5)</sup> Atkinson and Potts (1977), (6) Conding et al. (1976)

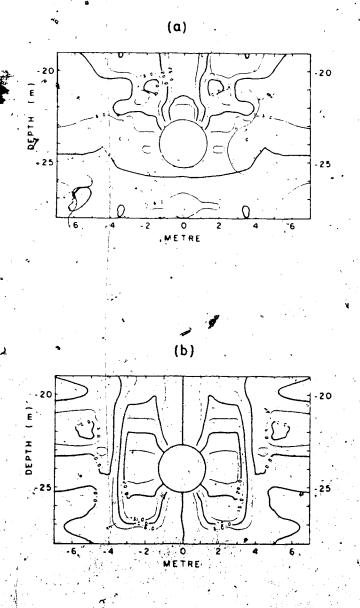


Figure C.1 Contours of Vertical and Horizontal Displacements around EXP Tunnel (modified from Eisenstein et al., 1981)

#### APPENDIX D

## Shaft Analyses

### D.1 Vertical Arching

Vertical soil arching sometimes can be referred to as "bin effect in silo" quantified by use of a horizontal differential element whose supports derive from the shear resistance on the interfaces of the stationary wall or soil mass (Fig. 4.2d). Summation of the vertical forces can be used to calculate the horizontal pressure required to maintain equilibrium of the soil mass in vertical directions. Because the shear resistances on the interfaces are derived from the frictional resistance and the shear strength of the cohesionless and cohesive materials respectively, the vertical arching action for these two types of materials is treated separately.

## Cohesionless Soils

Terzaghi (1936) adopted the vertical arching action developed behind the timbering of cuts and derived a relationship between the allowed wall yielding and the lateral pressure distributions behind. However, he had to assume some linear function on the coefficient of horizontal/vertical pressure (K) along the depth. Recently, Handy (1985) expanded his investigation on this coefficient (K) and found that K is a function of friction angle ( $\delta$ ) of the wall and soil internal friction angle ( $\phi$ ) of the

soil. The lateral pressure distribution calculated from the classic arching theory with adjusted K compared well with published data of some model and field tests. It is of practical importance that Terzaghi (1936) and Handy (1985), both assumed the plastic yielding behind the wall followed the Rankine's slip lines. However, this assumption is not applicable to the shaft case.

Vertical arching near a shaft can be quantitatively

described by the same concepts as found in the two papers

mentioned above, although several modifications are

necessary for their generalization. These will be discussed
below for the shaft behaviour.

Summation of vertical forces of a horizontal differential ring element in Fig. 4.2d gives

$$d\sigma_{V} = \left[\gamma - (2\pi\sigma_{V})/A \left(K \cdot \mu_{S} R/\sin\alpha\right) + K \mu_{W} a\right) dh \qquad D.1$$

In Eqn. D.1, the extent of the plastic zone, R can be determined from Eqn. 4.14, 4.19 or 4.25 depending on the material type and the mode of yielding. Typical values of K and K can be found in the paper by Handy (1985) for typical soil properties. These coefficients, K and K, define the horizontal to vertical stress ratios at soil and wall, depend on the stress distribution around the shaft which in turn are functions of displacement, material types, and modes of yielding. Fig. D.5 shows stress distributions around the shaft diagrammatically, along with ranges of K and K values for yielding Modes A and B. It is important to realize that the vertical stresses calculated from Eqn. D.1

are assumed to be constant across the plastic zone. This assumption should be taken into account when estimating the K and K values. In case of the yielding Mode A (K > K ) or cr where the vertical stress is always the intermediate stress, the K value at the elastic-plastic boundary is greater than K and less than K, but close to K. The K value at the wall depends on the displacement allowed ,i.e., K > K for small displacement and K < K for excessive yielding (tangential stress equal to vertical stress). For Mode B (K < K ) where vertical stress is always the principal or cr stress, the K value is slightly greater than K and the K is less than K. It can be seen that the K and K values vary along the depth of the shaft.

With all necessary parameters in Eqn. D.1 determined, the vertical stress distribution along the depth can be obtained by numerical integration of  $d\sigma$  over the whole depth (see the computer program SFT2, Appendix F), and the horizontal stress due to gravity effect is given by

D. 2

Cohesive Soils

The shear resistance of cohesive materials is independent of the confining pressure. Hence, applying horizontal pressure on the shaft wall does not enhance the strength of the material in the plastic zone. It is necessary to identify the possible collapse mechanism within the plastic zone and inhibit the mechanism by external forces.

Fig. D.4 shows one of the possible collapse mechanisms around the shaft. The collar within the plastic zone will cave in vertically and exert a resultant weight (W') on the conical mass beneath. To inhibit the horizontal caving in of the conical mass, horizontal pressure has to be applied. The required total force is given by considering the force equilibrium along the inclined failure surface of the conical mass

 $(W'+Wc)\sin a - P(\cos a) - F = 0$ D.3

## D.2 Mechanism of Shaft Behaviour in Cohesive Soil

This section contains equations governing the mechanism of shall behaviour in cohesive soil. Equations are designated in the same sequence as those for cohesionless soil except a prefix - D. Thus, the explanations in cohesionless soil section is applicable to this appendix. The notations and symbols are given in Appendix A.

The stresses in an elastic thick walled hollow cylinder are given by:

$$\sigma_{\mathbf{r}} = \chi \mathbf{h} = \mathbf{p}_{\mathbf{0}}$$

$$\sigma_{\mathbf{r}} = K_{\mathbf{0}} \mathbf{p}_{\mathbf{0}} - [K_{\mathbf{0}} \mathbf{p}_{\mathbf{0}} - \mathbf{p}_{\mathbf{i}}] (\mathbf{a}/\mathbf{r})^{2}$$
D.5

$$\sigma_{t} = K p + [K p - p] (a/r)^{2}$$
D.6

For a cohesive material the strength is assumed to be constant:

$$\sigma_1 - \sigma_3 = q$$

Hence, the support pressures for three possible modes of yield initiation are:

For 
$$\sigma_i - \sigma_i$$
:  $p = K p - q/2$ 

For 
$$\sigma_{\mathbf{v}} = \sigma_{\mathbf{i}} = \sigma_{\mathbf{i}}$$

For 
$$\sigma_t - \sigma_v$$
:  $p = (2K - 1)p - q_v$ 

D.10

The largest value of p will govern the mode of yield initiation, which can be expressed in terms of required K:

Mode A for 
$$\sigma_t - \sigma$$
:  $q_u/(2p_0) + 1 > K_0 > 1 - q_u/(2p_0)$ 
D.11

Mode B for 
$$\sigma_{v} - \sigma_{v}$$
:  $K < 1 - q_{u}/(2p_{o})$  D.12

Mode C for 
$$\sigma_t - \sigma_v$$
:  $K > q_u/(2p_o)+1$  D.13

Mode C has been neglected in the following analysis because it is of less practical significance. A critical K-value, K is used to distinguish Mode A from Mode B, i.e., Mode A and Mode B observed for K > K and K < K or cr respectively.

a) Mode A  $(\sigma_{t} - \sigma_{r})$  at K > K cr

The yielding initiates at the condition of:

$$\sigma_{t} - \sigma_{r} = 2K_{o}p_{o} - 2p_{i} = q_{u}$$

D.14

The magnitude of displacement corresponding to the yield initiation is:

$$u_{i} = [a(K_{o}p_{o}-p_{i})(1+\nu)]/E$$

D.15

Further decrease in the internal support pressure causes propagation of the plastic zone and increase in wall convergence which are as follows:

$$R_{tr} = a\{exp[(K_{o}p_{o}-p_{i})/q_{u}-1/2]\}$$
 D.16

$$u_i = a\{1-[1/(1+A_v)]^{1/2}\}$$
 D.1

b) Mode B  $(\sigma_r - \sigma_r)$  at K < K or cr

Yielding initiates at the condition of Eqn. D. 12 and the displacement at this stage is given by:

$$p = \sigma - q = K p$$

or 
$$u_i = a[(K_0 - K_0)p_0(1+\nu)]/E$$
 D. 18

As p is progressively reduced, the vertical stresses will also decrease due to the vertical shear stress induced by vertical downward displacement, and the tangential stress will increase until it becomes equal to the vertical stress at the wall. At this state, the radial stress in the plastic zone is given by:

$$\sigma = \sigma - q$$
 and  $\sigma = p$ 

Equating this stress with that in the elastic zone given by Eqn. D.5 yields the extent of plastic zone

$$R_{vr} = a\sqrt{\{(K_{o}p_{o}-p_{i})/[(K_{o}-1)p_{o}+q_{i}]\}}$$
 D.21

The corresponding wall convergence at this stage can be found by application of the model proposed by Ladanyi (1974),

$$u_i = a\{1-[1/(1+\hat{A})]^{1/2}\}$$
D.22

and the support pressure is

$$p_{i} = (K_{0}p_{0}-q_{u}/2)$$

Further relief in support pressure induces the propagation of the plastic zone. The mode of yielding is still controlled by Mode B. But near the shaft wall, there exists a zone at which the tangential stress becomes equal to the vertical stress. Within this zone, the radial and tangential stresses governed by the failure/criteria are given as Ladanyi (1974),

$$\sigma_{r} = p_{i} + q_{i} (\ln(r/a))$$
D.24

$$\sigma_{t} = p + q (1 + \ln(r/a))$$
D.25

and the extent of this zone is

$$R_{tr} = a \{ exp[K_{o}p - p_{i})/q_{u} - 1/2] \}$$
 D.26

To calculate the extent of the plastic zone due to Mode B, one needs to know the radial stress distribution. Between the distance (R < r < R > vr shown in Fig. 4.9, the condition of equilibrium with tangential stress gives radial stress as

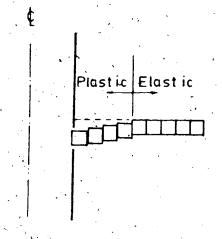
$$\sigma_{r} = K p - [(K p - \sigma_{r})(R_{r}/r)^{2}]$$
D.26

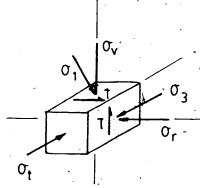
where:  $\sigma_r' = 2K p - q$  at r = R tr

Substituting r = R and  $\sigma$  into Eqn. D.26 yields the radial stress distribution in the plastic zone (R tr < r < R). Continuity of radial stresses at the elastic zone boundary ( $\sigma$  = K  $\sigma$ ) locates the extent of the plastic zone, R as:

$$R_{vr} = \frac{vr}{a\sqrt{\{(K_{o}p_{o}-q_{u})exp[2((K_{o}p_{o}-p_{i})/q_{u}-1/2)]/[(1-K_{o})p_{o}+q_{u}]\}}}$$

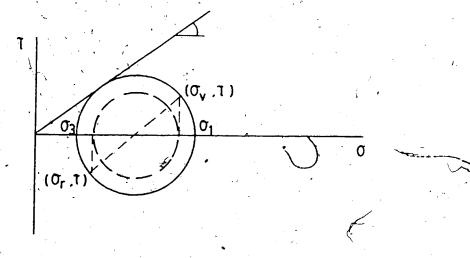
The wall displacement can also be found similar to Mode A.





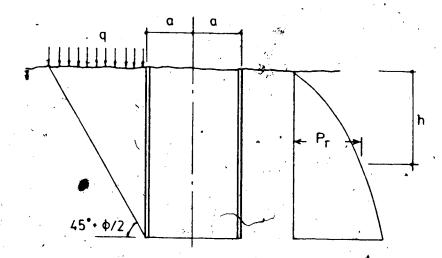
(a) Schematic Movement in Plastic Zone

(b) Stresses on Element in Plastic Zone



(c) Mohr Diagram

Figure D.1 Terzaghi's Method (1943)



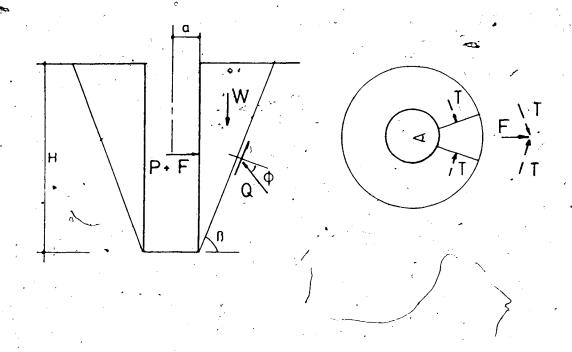
## Berezantzev (1958)

Radial stresses (Pr) acting on a vertical cylindrical sunface is given by:

$$P_r = ah \frac{\sqrt{K}}{g-1} (1 - (\frac{a}{r})^{g-1}) + q(\frac{a}{r})' + c \cot \phi ((\frac{a}{r})^g K'-1)$$

where: 
$$K = \tan^2(45^\circ - \Phi/2)$$
  
 $r = a + h/K$   
 $g = 2 \tan \Phi \tan (45^\circ + \Phi/2)$ 

Figure D.2 Berezantzev's Method (1958)



radius of shaft

depth of shaft

angle between the horizontal and the failure surface of the cone

angle of internal friction of soil

weight of the sliding cone

reaction acting on the sliding cone earth pressure acting on the shaft lining

tangential force which has a radial component Facting in outward direction, i.e.

 $F = 2 \tan(\phi/2)$ 

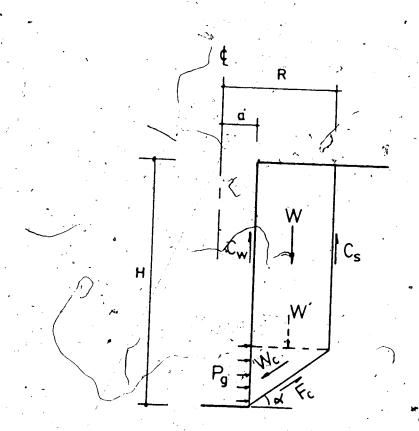


Figure D.4 Gravity Effect in Cohesive Soils

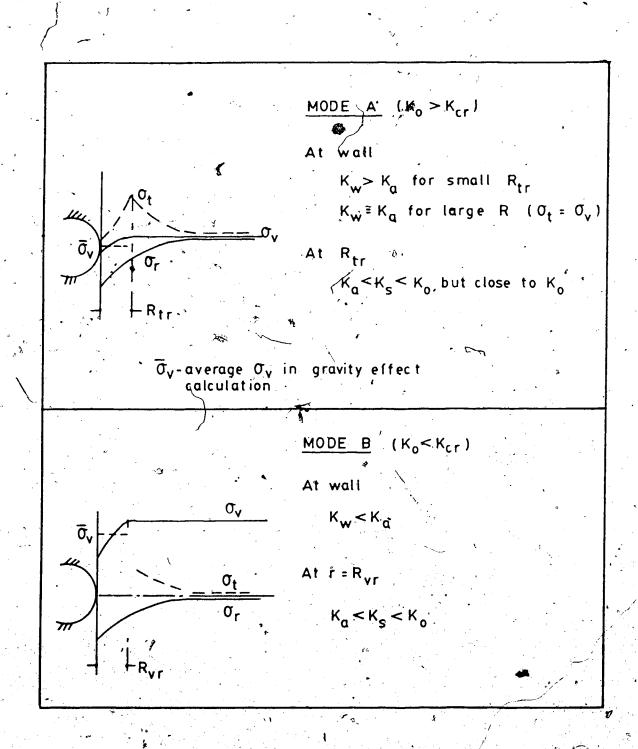


Figure D.5 Coefficients of Vertical-Radial Stress (R, K)

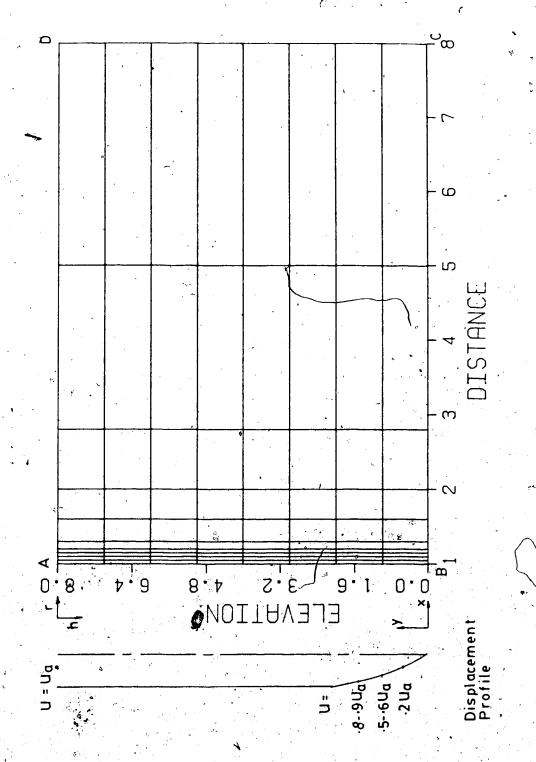


Figure D. Finite Element Mesh for Shaft Analysis

APPENDIX E

Shaft Case Histories

Table E.1 Properties of Melinex (Shaft Lining)

Tensile Strength
Yield Strength
Young's Modulus
(at 1% strain)
Elongation at
Yield Point
Coefficient of Thermal
Expansion

17.65 kN/sq.cm:
9.81 kN/sq.cm.
430.0 kN/sq.cm.
5 %
17.10 x10 -6 cm/cm-C°

# Table E.2 Properties of Leigthon Buzzard Sand 120/220

Specific Gravity 2.66
Minimum Density 12.95
Maximum Density 15.86
Average Grain Size 0.13
Coefficient of Uniformity 1.40
(=D60/D10) 2.66 12.95 kN/m3 15.86 kN/m3 \*0.13 mm

Density in Test Density Index Frictional Angle

1.535-1.550 kN/m3 85%-90% 38.3 deg

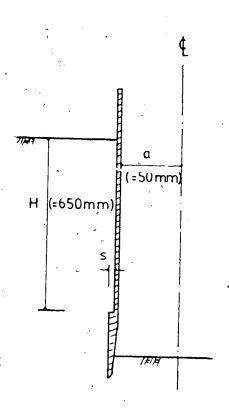


Figure E.1 Shaft Details (modified from Muller-Kirchenbauer et al., 1980)

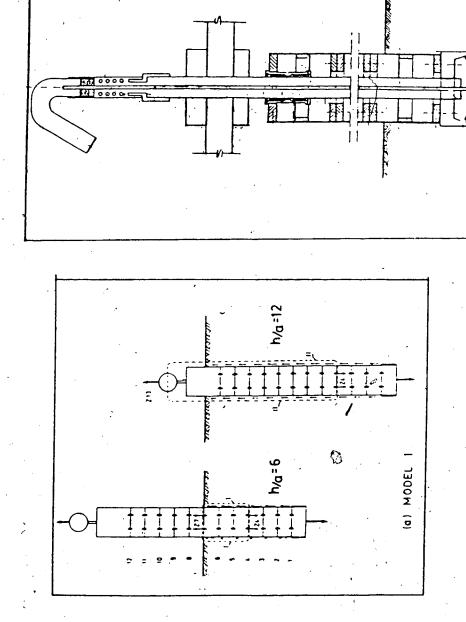


Figure E.2 Shaft Models-Types I and II (modified from

Muller-Kirchenbauer et al., 1980)

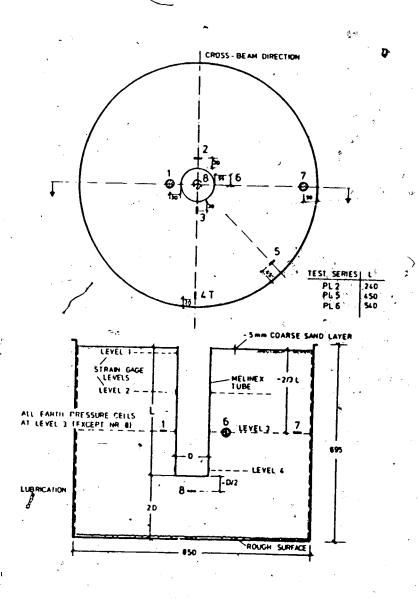


Figure E.3 Model Test Setup (modified from Lade et al., 1981)

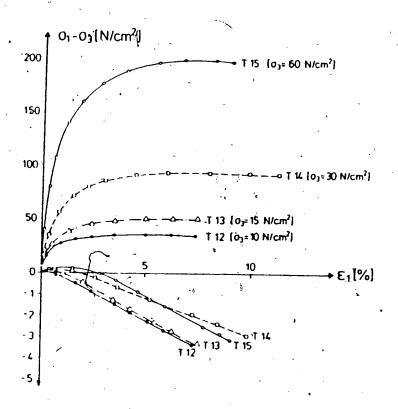


Figure E.4 Stress-Strain and Volume Change Curves from Triaxial Compression Tests on Leighton Buzzard Sand (modified from Lade et al., 1981)

## APPENDIX F

Computer Program

## Tunnel Analysis

TUN1 - Calculate support pressure for yield initiation at the roof (Line c-d, Fig. 2.4)

TUN2 - Calculate support pressure for yield initiation at the springline (Line h-k, Fig. 2.4)

TUN3 - Determine Regime 'J' of Mode II (Line i-1, Fig. 2.4)

TUN4 - Calculate support pressure for final equilibrium state ( $p_{fc}$ )

TUN5 - Calculate support pressure for p (LIne e-f, Fig. 2.4)

## Shaft Analysis

SFT1 - Calculate Ground Convergence Curve of soil model (Ladanyi, 1974)

SFT2 - Calculate support pressure for vertical arching

```
C*
  3
        C+
  4
  5
        С
        C CALCULATION OF SUPPORT PRESSURE(IN TERM OF YH) FOR
              YIELD INITIATION AT ROOF 'IN TUNNEL
  8
  9
        С
 10
        С
           USING:
 11
           ELASTIC STRESS DISTRIBUTION AROUND A CIRCULAR TUNNEL .
12
        С
 13
            UNDER INTERNAL PRESSURE USING SCHMIDT METHOD
 14
 15
        С
 16
        С
           INPUT FILE:
              1. KO,PHI,DEP
        С
17
              2. EOF
        C
 18
19
        С
           OUTPUT FILE:
20
              SUPPORT STRESS IN NEGATIVE
21
       ٠C
22
        С
23
        Ċ
              READ(5,511) XKOO, PHIO, DEPO
24
        511
              FORMAT(3F10.0)
25
26
        С
27
              DXK0=0.1 /
              DDEP=1.
28
29
              DPHI=5.
        С
30
           00 11 I=1,5
31
32
              PHI=PHIO+DPHI*(I-1)
              XPHI=PHI*3.14159/180.
33
              XNP = (TAN(3.14159/4.+0.5*XPHI))**2.
34
              WRITE(6.601) PHI, XNP
              FORMAT(/, ' ANGLE =', F6.2,'
36
      . 601
                                                   ',F6.3)
37
       С
38
              DO 22 J=1.8
              XKO=XKOO+DXKO*(J-1)
39
              XMIL = XKO/(1.+XKO)
40
41
       С
42
              XN=1.0/XMIL
              D1=(1,+1,/(XN-1.))*0.5
43
44
              D2=(1.-1./(XN-1.))+0.5
              XMIL1=(XN+1)*(XN-1.)/XN
45
              XMIL2=(XN+1)*(XN-2)/XN
46
47
              XM = -XMIL2/(4 + XMIL1)
48
       C.
              WRITE(6,602) XKO, XMIL
49
50
       602
              FORMAT(' KO = ', F5.3, ' POISSON RATIO =', F6.2)
51
       C
              WRITE(6,603)
52
53
       603
              FORMAT(' H/A ','
54
              DO '33 L+1.9
              DEP = DEPO +DDEP*(L-1)
55
56
       C
57
       С
            AT ROOF
       C٠
58
59
              FETA=3.14159/2.
60
              Q=1.0
```

```
′ح`ار
                  XST* -D1*(1.+Q)+
& (D1-0.5*D2-XM*Q-(XM+0.5*D2)*Q*Q)*(1./DEP)*SIN(FETA)-
     62
     63
     64
                      D2+(1.+3+Q+Q)+COS(2+(FETA))+
    65
                    0.5*D2*(1.-Q*Q+4*Q*Q*Q)*(1./DEP)*SIN(3*FETA)
    66
            С
    67
                   ALPHA = XST/(XNP+1.)
    68
            С
                   WRITE(6,612) DEP .ALPHA FORMAT(F5.2,F10.3)
    69
    70
            612
    71
            33
                   CONTINUE
    72
            22
                   CONTINUE
    73
                   CONTINUE
    74
                   STOP
    75
                   END
End of file
     2
            C 4
     3
     4
                  ****** TUN2*****
            С
     6
               CALCULATION OF SUPPORT PRESSURE(IN TERM OF YH) FOR
            С
     7
                 YIELD INITIATION AT SPRINGLINE IN TUNNEL
     9
            C+
    10
            С
            C
               ELASTIC STRESS DISTRIBUTION AROUND A CIRCULAR TUNNEL
    12
            C
    13
           C
               UNDER INTERNAL PRESSURE USING SCHMIDT METHOD
    15
           С
    16
           С
               INPUT FILE:
    17
           C
                  1. KO, PHI, DEP
    18
           C
                  2 EOF
    19
           С
   20
           C
              OUTPUT FILE:
   21
           C
                  SUPPORT STRESS IN NEGATIVE
   22
           C
   23
           С
   24
                  READ(5,511) XKOO, PHIO, DEPO
   25
           511
                  FORMAT (3F 10.0)
   26
           С
   27
                  DXK0=0.1
   28
                  DDEP=1.
   29
                 DPHI=5.
   30
           C
   31
                  DD 11 I=1.5
                 PHI=PHIO+DPHI+(I-1)
   32
   33
                 XPHI=PHI=3.14159/180.
   34
                 XNP =(TAN(3.14159/4.+0.5+XPHI))++2.
                 WRITE(6,601) PHI, XNP
FORMAT(/; ANGLE =',F6.2.'
   35
   36
          601
                                                         '.F6.3)
  37
          C
  38
                 DO 22 J=1.8
  39
                 XKO=XKOO+DXKO+(J-1)
  40
                 XMIL = XKO/(1.+XKO)
  41
  42
                 XN=1.0/XMIL
                 D1=(1.+1./(XN-1.))*0.5
  43
  44
                 D2=(1.-1./(XN-1.))+0.5
```

```
XMIL1=(XN+1)*(XN-1.)/XN
     45
                   XMIL2=(XN+1)*(XN-2)/XN
     46
     47
                   XM=-XMIL2/(4*XMIL1)
     4R
            С
                   WRITE(6,602) XKO,XMIL
FORMAT(' KO = ',F5.3, ' POISSON RATIO =',F6.2)
     49
     50
             602
     51
            С
     52
                   WRITE(6,603)
                   FORMAT(' H/A ','
     53
            603
                                          PI/YH './)
     54
                   DO 33 L=1.9
     55
                   DEP = DEPO + DDEP * (L-1)
     56
            С
     57
            С
                 AT SPRINGLINE
     58
            С
                                                                            ( ·
     59
                   FETA=0.0
     60
                  Q=1.0 .
     61
            C
                           -D1+(1,+Q)+
    62
    63
                    (D1-0.5*D2-XM*Q-(XM+0.5*D2)*Q*Q)*(1./DEP)*SIN(FETA)-
    64,
                     D2*(1.+3*Q*Q)*COS(2*(FETA))+
    65
                    0.5*D2*(1.-Q*Q+4*Q*Q*Q)*(1./DEP)*SIN(3*FETA)
    66
            C
    67
                  ALPHA = XST/(XNP+1,)
    68
           C
    69
                 "WRITE(6,612) DEP
                                     , ALPHA
    70
           612
                  FORMAT(F5.2, F10.3)
    71
           33
                  CONTINUE
    72
           22
                  CONTINUE
    73
                  CONTINUE
    74
                  STOP
    75
                  END .
End of
     1
           C
     2
           C
          .C•
                  ******* TUN3*******
           С
                STRESS DISTRIBUTION WITHIN THE PLASTIC ZONE
           С
                             - SOLVING CAUCHY PROBLEM
           t
                (INTERNAL PRESSURE WITH SHEAR IN % OF INSITU STRESS)
    8
           C
                             (MARCH 21, 1985)
    9
           C • •
    10
           С
   11
           C INPUT FILE:
                1. RAD, HT, PI(%), ANGLE, COHESION, DENSITY, POISSON'S RATIO, TOLEN.
   12
           C
   13
          С
                2. EOF
   14
   15
   16
          C OUTPUT INTERPRETATION
                +VE X-AXIS POINTING DOWNWARD, +VE Y-AXIS RIGHT
   17
   18
                R - RADIUS , THETA- ANGLE BETWEEN X AND R DIRECTIONS
   19
                DIVERGENCE OCCURED WHEN THE DIFFERENCE IN FED EXCEEDS
   20
                THE SPECIFED TOLERANCE, I.E. HYPERBOLIC CONDITION NOT
          C
   21
          C
                VALID AND DIFUTCATION INVOLVED
   22
          C
  23
                 COMMON TUNPRS, RAD, PHI, COH , W.
  24
  25
                        X(20.20), Y(20.20), TH(20.20), SP(20.20), SR(20.20),
  26
                        NP, NF , IXO, IYO, NBC, DIVAN, TOLEN
  27
  28
                WT OF SOIL
```

```
29
         C TUNPRS - INTERNAL SUPPORT PRESSURE %OF INSITU STRESS
         C TOLEN - TOLERANCE ALLOWED IN CALCULATION FOR ACCURACY
 30
 31
         C RAD - RADIUS OF TUNNEL
         C'HT - DEPTH OF OVERBURDEN
 32
         C PHI - INTERNAL FRICTIONAL ANGLE OF SOIL
 33
         C COH COHESION OF SOIL
 35
         C XKO - INITIAL INSITU STRESS COEFFICIENT
 36
         C P - NODAL POINT IN CHARACTERISTIC
 37
         C X - X COORDINATES
 38
         C Y - Y COORDINATES
         C SP - STRESS FUNCTION
 39
         C SR - STRESS FUNNCTION
 40
 41
         C TH -ANGLÉ BETWEEN X -AXIS AND PRINCIAPL STRESS
 42
 43
         C
            INPUT DATA
 44
                READ (5,501) RAD, HT, TUNPRS, PHI, COH, W, XKO, TOLEN
 45
         501
                FORMAT(8F10.0)
 46
                WRITE(6,601) RAD, HT, TUNPRS, PHI, COH, W, XKO, TOLEN
                FORMAT( ' RADIUS OF TUNNEL "'.F6.2./.
 47
         601
 48
               & ' DEPTH OF OVERBURDEN = ',F6.2./,
 49
                   TUNNEL SUPPORT PRESSURE =',F10.2./.
 50
                   FRICTIONAL ANGLE =',F6.2./.
                   51.
                / WEIGHT = ', F6,2\/,
' KO = ',F6.2/,
' TOL = ',F8.5,//)
 52
 53
 54
 55
         C ·
            INITIAL DATA
 56
 57
               NF = 10
 58
           NF - NO. OF FORWARD CALCULATION (PLASTIC ZONE EXTENT)
                PHI=PHI+3.1416/180.
 59
 60
                XNP * TAN(0, 25*3, 1416+0, 5*PHI) * TAN(0, 25*3, 1416+0, 5*PHI)
61
         С
                   WRITE 6,698)
FORMAT(10X,' X Y',' R',5X,'THETA',3X,'SX',3X,'SY','
S1 53 SR ST SSH ',/)
62
' 63
64
               8
6$
           BOUNDARY CONDITIONS
 еď
                NBC=9 /
67
         C NBC - NO. OR KNOWN BOUNDARY CONDITIONS
68
                IXO=1
69
                IY0#9
 7p
                DIVAN=3.14159/8 .
 7 h
                DO 11 I=1.NBC
72
                II=I-1
                MX=1X0+11
 73
74
                NY = I YO+ I'I
75
                X( MX,NY)=RAD*COS((II)*DIVAN)
76
                Y( MX .NY )=RAD*SIN((II)*DIVAN)
77
                TH( MX ,NY )=0.5 +3.4416+(11)+DIVAN
78
            INITIAL SXO .SYO STRESSES IN X . Y DIRECTIONS
                SYO = W*(HT+X(MX,NY))
79
           SXO = XKO*W*(HT+X(MX , NY))

INITIAL NORMAL AND SHEAR STRESSES ON THE SURFACE

SNO = 0.5*(SXO+SYO) + 0.5*(SXO-SYO)*CDS(2*TH(MX,NY))

STO = -0.5*(SXO-SYO)*SIN(2*TH(MX,NY))
80
81
82
83
84
            SUPPORT PRESSURE - % OF INSITU STRESS
85
                SN = SNO*TUNPRS
                ST = STO+TUNPRS
86
87
        C . CALCULATION OF P , R STRESS PARAMETERS
                          COH+(1./TAN(PHI))
88
```

```
VX1 = (SN*SN*SIN(PHI)*SIN(PHI)-ST*ST*COS(PHI)*COS(PHI))
  89
                 VX2 = 2*CX*SN*SIN(PHI)*SIN(PHI)+CX*CX*SIN(PHI)*SIN(PHI)
 ,90
                VX3 = SORT(VX1+VX2)
ROOT1=(SN+CX-VX3)/(COS(PHI)*COS(PHI))
  91
  92
          C
                 ROOT2=(SN+CX+VX3)/(COS(PHI)+COS(PHI))
  93
  94
                 SP( MX,NY) =ROOT2, - CX
  95
          С
                 P2=R00T1-CX
                 SR( MX,NY) =SP(MX,NY)*SĮN(PHI)+COH*COS(PHI)
  96
  97
                    R = SQRT(X(MX,NY)*X(MX,NY)+Y(MX,NY)*Y(MX,NY))
  98
                    AN = (ATAN(Y(MX, NY)/X(MX,NY)))
  99
 100
                    DAN = AN + 180./3.1416
 101
                    SX = SP(MX,NY) + SR(MX, NY) + COS(2+TH(MX,NY))
                    SY = SP(MX,NY) / SR(MX, NY) *COS(2*TH(MX,NY))

SXY = SR(MX,NY) / SIN(2*TH(MX,NY))
 102
 103
 104
                    S1 = (SP(MX,NY) + SR(MX,NY))
                    S3 = (SP(MX,NY) - SR(MX,NY))
 105
                    SGT = SX+(SIN(AN)+SIN(AN)) +SY+(COS(AN)+COS(AN))-
 106
 107
                     2*SXY*SIN(AN)*COS(AN)
108
                    SGR = SX*(COS(AN)*COS(AN)) +SY*(SIN(AN)*SIN(AN))+
109
                    2*SXY*SIN(AN)*EOS(AN)-
110
                    SGSH = 0.5*(SGT-SGR)*SIN(2*AN)+SXY*(COS(2*AN))
         C
.114
                 WRITE(6,612) MX, NY ,X(MX, NY),Y(MX, NY ),R,DAN,
112
                SX, SY,SXY,S1, S3,SGR,SGT,SGSH
FORMAT( ' P(',1'2,',',12,') =',3F6.2, F6.1,8F7.2)
113
114
         612
                 WRITE(6,613) CX, VX1, VX2, VX3, RODT1 .ROOT2
115
         C
116
         C613
                 FORMAT(6F10.3)
117
         C
                 WRITE(6,614) SXO. SYO, SNO, STO, SN, ST , P2
       , C614
118
                  FORMAT (7F 10.2) -
119
         C.
120
         11
                CONTINUE
121
         ·C
122
        - C
123
                   CALL NODE
         C
124
125
                   STOP
126
                   END
         C. . . .
127
128
         C
         C
129
             SUBROUTINE NODE
130
                 SUBROUTINE NODE
               COMMON TUNPRS,
131
                                RAD, PHI,COH, W.
                       X(20,2d), Y(20,20), TH(20,20), SP(20,20), SR(20,20),
132
              &
133
                       NP,NF , IXO,IYO, NBC,DIVAN,TOLEN
134
         С
135
                 XAN=(3.1416/4)-0.5*PHI
        C
136
             ICON - COUNTER TO LIMIT NO DE ITERATION
137
        C
138
        C
139
                 DO 10 I=1,NF
140
                   WRITE(6,699) I
FORMAT(' FORWARD CAL. NO.
141
        699
142
                   NNBC+NBC-1
143
                   DO 20 11=1 NNBC
144
                    ICONT 1=0
145
                    NC1=II-1
146
        C
               NC1 - COUNTER MOVING DIAGONAL DOWNARD
147
                    IM=IXO+NC1
148
                    JM=IYO+NC1
```

1

```
IN= I'XO+NC 1+1
149
                    JN=IYO+NC1+1
150
151
                    XA = X(IM,JM)
                    XB = X(IN,JN)
152
                    YA = Y(IM, JM)
153
                    YB - Y(IN, JN)
154
155
                    THA=TH(IM, JM)
                    THB=TH(IN, JN)
156
                    SPA=SP(IM.JM)
157
                    SPB=SP(IN.JN)
158
                    SRA=SR(IM,JM)
159
                    SRB=SR(IN, JN)
160
161
         C FIRST ESTIMATES OF XP. YP. SRP AT POINT P
162
163
                    XP1 = 0.5*(XA+XB)
164
                    YP1 = 0.5*(YA+YB)
165
                    THP 1=0.5* (THA+THB)
1666
                    SPP1=0.5*(SPA+SPB)
167
168
                    SRP1=0.5*(SRA+SRB)
                     WRITE(6,904) XP1, YP1, THP1, SPP1, SRP1
169
                     FORMAT(5F10.3,//)
         C904
170
171
         C SOLVE FOR SPP AND THP
         1000
                    ICONT 1 = ICONT 1+1
172
                    IF (ICONT1.GT.20) GO TO 999
173
174
                    E1=-SIN(2*XAN)
175
                    E2# SIN(2*XAN)
                    F1=-W*SIN(2*XAN)*(XP1-XA)
176
                    F2=-W*SIN(-2*XAN)*(XP1-XB)
177
                    G1=SRP1+SRA
178
                    G2=SRP1+SRB
179
                    H1=-W*COS(2*XAN)*(YP1-YA)
180
                    H2=-W*CDS(2*XAN)*(YP1-YB)
181
                    SPP2=(G2+(F1+H1)+G1+(F2+H2)+E1+G2+SPA-E2+G1+SPB+
182
                    G1*G2*(THA-THB))/(E1*G2-E2*G1)
THP2=(E2*(F1+H1)-E1*(F2+H2)+E1*E2*(SPA-SPB)+
              8
183
184
                      E2*G1*THA-E1*G2*THB)/(E2*G1-E1*G2) 4
185
                    SRP2= SPP2*SIN(PHI)+COH*COS(PHI)
186
                     WRITE(6,901) E1, E2,F1,F2,G1,G2,H1,H2
         С
187
         C901
                    FORMAT(8F10.3,//)
188
         C
                    WRITE(6,902) SPP2, THP2, SRP2
189
         C902
                    FORMAT(3F10.3,//)
190
         С
191
             SOLVE FOR XP, YP
192
         С
193
                     = (0.5*(THP2+THA)-XAN)
194
                   B = (0.5 + (THP2 + THB) + XAN)
195
                   XP2 = (YB-YA+XA*TAN(A)-XB*TAN(B))/(TAN(A)-TAN(B))
196
                   YP2=((XB-XA)*TAN(A)*TAN(B)+YA*TAN(B)-YB*TAN(A))/
197
                    (TAN(B)-TAN(A))
.198
                    WRITE (6,905) A. B. XP2,YP2
199
         Ç
                    FORMAT(4F1-.3',7/)
200
         C905
         C
201
            TEST FOR CONVERGENCE
202
         С
203
                    TEST = ABS(XP2-XP1)+ABS(YP2-YP1)+ABS(SPP2-SPP1)
204
205
                    +ABS(THP2-THP1)
                    IF (TEST .LT TOLEN) GO TO 2000
206
                    WRITE(6,906) TEST
207
         C905
                    FORMAT(F10.2)
208
```

```
209
 210
 211
 212
                    SPP1 =SPP2
 213
                   GD TD 1000
 214
 215
          2000
                    M = I M + 1
 216
                    X( M, JM )
 217
                    Y( M, JM
 218
                   TH( M, JM
                               =THP2
 219
                   SP( M, JM ) *SPP2
                   SR( M, JM") =SRP2
 220
 221
                   DTH = TH(
                              M, JM) * 180./3.1416
                   R = SQRT(X(M, JM)*X(M, JM)+Y(
 222
                                                      _M, JM\^Y(
                   AN = (ATAN(Y( M, JM)/X( M, JM)))
 223
 224
                   DAN = AN+180./3.1416
                                          M, JM )*COS(2*TH(
 225
                   SX = SP(M, JM) + SR(
                                                                    JM ))
 226
                                         M, JM )+COS(2+TH(
                   SY = SP(M,JM)-SR(
                                                               M, JM ))
                   SXY = SR(M, JM) + SIN(2 + TH(M, JM))
 227
 228
                   $1 =(SP(
                             M, JM)+SR( M, JM))
 229
                            M, JM)-SR( M, JM))
                   S3 =(SP(
 230
                   SGT = SX*(SIN(AN)*SIN(AN)) +SY*(CDS(AN)*COS(AN))-
 231
                   2*SXY*SIN(AN)*COS(AN)
 232
                   SGR = &X*(COS(AN)*COS(AN)) +SY*(SIN(AN)*SIN(AN))+
 233
                   2*SXY*SIN(AN)*COS(AN)
 234
                   SGSH = 0.5*(SGT-SGR)*SIN(2*AN)+SXY*(COS(2*AN))
 235
         С
 236
                WRITE(6,603)
                              M, JM, X( M, JM), Y( M,
                                                           JM ),R,DAN,
237
                 SX, SY, SXY, . $1, S3, SGR, SGT, SGSH
 238
         603
               FORMAT( ' P('.12,',',12,') =',3F6.2, F6.1,8F7.2)
239
         С
             INTERCHANGE IMAGES IF II=1 OR NBC-1
240
         C
241
        , C
                  IF (II.EQ. () GO 40 93
242
243
                 GD TO 94
244
                      IM, JM-1) = X(IM+1, JM)
         93
                  X(
245
                  Y (
                      IM, JM-1) =-Y(IM+1, JM)
246
                TH( IM, ...UM-1) =TH(IM+1, UM )-
                                                 NAVIO
247
                  SP( IM, JM-1) *SP(IM+1, JM )
248
                `~-SR( IM, JM-1) =SR(IM+1, JM )
249
                  MUM=UM-1
250
                  DTH = TH(IM, MJM) + 180./3.1416
251
                  R = SQRT(X(IM,MJM)*X(IM,MJM)*Y(IM,MJM)*Y(IM,MJM))
252
                  \dot{A}N = (ATAN(Y(IM,MUM)/X(IM,MUM)))
253
                  DAN =AN+180./3.1416
254
                  SX = SP(IM,MJM) + SR(IM,MJM) + CDS(2+TH(IM,MJM))
255
                  SY = SP(IM,MJM) - SR(IM, MJM) + COS(2+TH(IM,MJM))
                  SXY = SR(IM,MJM)*SIN(2*TH(IM,MJM))
256
257
                  S1 = (SP(IM,MJM)+SR(IM,MJM))
258
                  S3 = (SP(IM, MJM) - SR(IM, MJM))
259
                  SGT = SX*(SIN(AN)*SIN(AN)) +SY*(CDS(AN)*CDS(AN))-
260
                   2*SXY*SIN(AN)*COS(AN)
                  SGR = SX*(CDS(AN)*CDS(AN)) +SY*(SIN(AN))+
261
                   2*SXY*SIN(AN)*COS(AN)
262
263
                  SGSH = 0.5*(SGT-SGR)*SIN(2*AN)+SXY*(CDE(2*AN))
264
        C
               WRITE(6,604) IM, MJM ,X(IM, MJM),Y(IM, MJM),R,DAN;
265
266
                SX.SY.SXY.S1.S3,SGR.SGT.SGSH
267
        604
              FORMAT( ' P(',12,',',12,') =',\3F6.2, F6.1,8F7.2)
268
        С
```

```
TF (II.EQ.(NBC-1)) GO TO 95
269
       . 94
                  '60 TO 20
270
                       IM+2, JM+1) = X(IM+1, JM)
271
         95
                   Х(
                       IM+2, JM+1) =-Y(IM+1, JM
272
                   TH( IM+2, JM+1) =TH(IM+1, JM )+ DIVAN
273
                   SP( IM+2, JM+1) *SP(IM+1, JM )
SR( IM+2, JM+1) *SR(IM+1, JM )
274
275
276
                   NIM=IM+2
277
                   NUM=JM+1
                   DTH = TH(NIM, NJM) * 180./3.1416
278
                   R = SQRT(X(NIM,NJM)*X(NIM,NJM)*Y(NIM,NJM)*Y(NIM,NJM))
279
                   ((MUN,MIN)X/(MUN,MIN)Y)NATA) = NA
280
                   DAN = AN+180./3.1416
281
                   SX = SP(RIM, NJM) + SR(NIM, NJM) + COS(2*TH(NIM, NJM))

SY = SP(NIM, NJM) - SR(NIM, NJM) + COS(2*TH(NIM, NJM))
282
283
                   SXY = SR(NIM,NJM)+SIN(2+TH(NIM,NJM))
284
                   SPP1 = SPP2
285
                   S-1 = (SP(NIM,NJM) + SR(NIM,NJM))
286
                   S3 = (SP(NIM, NJM) - SR(NIM, NJM))
287
                   SGT = SX*(SIN(AN)*SIN(AN)) +SY*(COS(AN)*COS(AN))-
288
                    2+SXY+SIN(AN)+COS(AN)
289
                   SGR = SX*(COS(AN),*COS(AN)) +SY*(SIN(AN)*SIN(AN))+
290
291
                    2*SXY*SIN(AN)*COS(AN)
                   SGSH = 0.5*(SGT-SGR)*SIN(2*AN)+SXY*(COS(2*AN)) 7
292
         Ċ
293
                 WRITE(6,605) NIM, NUM ,X(NIM, NUM),Y(NIM, NUM),R,DAN.
294
               8 SX.SY.SXY.S1.S3.SGR.SGT.SGSH
295
                FORMAT(. ' P(',12,',',12,') =',3F6.2, F6.1,8F7.2)
         605
296
         20
297
                 CONTINUE
         С
298
         C COUNTER ( MOVING CENTRIOD IXO, IYO SIDEWAY RIGHT BY 1)
299
300
                  IYO= IYO- 1
301
         С
302
303
                 DO 30 LL=1, NBC
                   ICONT2=0
304
305
                   NC2=LL-1
                NC2 - OUNTER MOVING DIAGONAL DOWNARD
306
                    IM=IXO+NC2
307
                    JM=IYO+NC2
308
309
                    IN=IXO+NC2+1
                    JN= IYO+NC2+1
310
                    XA = X(IM, JM)
311
312
                    XB = X(IN,JN)
                    YA = Y(IM.JM)
313
                    YB = Y(IN,JN)
314
                    THA=TH(IM.JM)
315
                    THB=TH(IN, JN)
316
                    SPA=SP(IM, JM)
317
                     SPB=SP(IN.JN)
318
                    SRA=SR(IM.JM)
319
320
                     SRB=SR(IN, JN)
321
         C FIRST ESTIMATES OF XP. YP. SRP AT POINT P
322
323
                    XP1 = 0.5*(XA+XB)
324
                     YP1 = 0.5*(YA+YB)
325
326
                     THP1=Q.5+(THA+THB)
                     SPP1=Q.5+(SPA+SPB)
327
                     SRP1=0.5+(SRA+SRB)
328
```

```
329
                               SOLVE FOR SPP AND THP
   330
                                            ICONT2 * ICONT2+1
                      3000
   331
                                            IF (ICONT2.GT.20) CO TO 999
   332
   333
                                            E1=-SIN(2*XAN)
                                            E2= SIN(2*XAN)
   334
                                            F1=-W*SIN(2*XAN)*(XP1-XA)
   335
                                            F2=-W*SIN(-2*XAN)*(XP1~XB)
   336
                                            G1=SRP1+SRA
   337
                                            G2=SRP1+SRB
   338
                                            H1=-W+COS(2+XAN)+(YP1-YA)
   339
                                            H2 = - W + COS (2 + XAN) + (YP1 - YB)
   340
                                            SPP2=(G2*(F1+H1)-G1*(F2+H2)+E1*G2*SPA-E2*G1*SPB+
   341
                                              G1*G2*(THA-THB))/(E1*G2-E2*G1)
  342
                                8
                                            THP2=(E2*(F1+H1)-E1*(F2+H2)+E1*E2*(SPA-SPB)+
   343
                                                E2*G1*THA-E1*G2*THB)/(E2*G1-E1*G2)
  344
                                            SRP2= SPP2*SIN(PHI)+COH*COS(PHI)
  345
  346
                     C
                     C
                              SOLVE FOR XP, YP
  347
                     C
  348
                                         A = (0.5*(THP2+THA)-XAN)
  349
  350
                                         B = (0.5 \cdot (THP2 + THB) + XAN)
  351
                                         XP2 = (YB-YA+XA+TAN(A)-XB+TAN(B))/(TAN(A)-TAN(B))
                                         YP2* ((XB-XA)*TAN(A)*TAN(B)+YA*TAN(B)-YB*TAN(A))/
  352
  353
                                           (TAN(B)-TAN(A))
4 354
                    C
                           TEST FOR CONVERGENCE
  355
                    C
  356
                                           TEST = ABS(XP2-XP1)+ABS(YP2-YP1)+ABS(SPP2-SPP1)
  357
                                8
                                              +ABS(THP2-THP1)
  358
  359
                                           IF (TEST .LT. TOLEN) GO TO 4000
                                           XP1 =XP2
  360
                                           YP1=YP2
 361
  362
                                           SRP1 =SRP2
                                           SPP1 #SPP2
  363
                                         GD TD 3000
  364
  365
                    C
                      4000
                                           M= IM+ 1
  366
                                           X(M,JM) = XP2
  367
  368
                                           Y( M, JM
                                                               ) =YP2
                                         TH( M, JM ) =THP2
  369
                                         SP( M, JM ) =SPP2
  370
  371
                                         SR( M, JM ) = SRP2
                                                                .M. JM) * 180./3.1416
                                         DTH = TH(
  372
                                         R = SQRT(X(-M, JM)*X(-M, JM)+Y(-M, JM)+Y(-M,
                                                                                                                        M, JM)*Y( M, JM))
  373
  374
                                         AN = (ATAN(Y(M, JM)/X(M, JM)))
                                         DAN = AN+180./3.1416
 375
                                                               M, JM )+SR( M, JM )*COS(2*TH( M, JM )*COS(2*TH(
                                         SX = SP(
                                                                                                                                          M , UM ))
 376
 377
                                         SY = SP(
                                                                                                                                            M, JM ))
                                         SXY = SR(M, JM) + SIN(2+TH(
                                                                                                          M, JM))
 378
                                        S1 =(SP(
 379
                                                               M, JM)+SR( M, UM))
                                         S3 =(SP(
                                                               M, JM)-SR( M, JM))
 380
                                         SGT = SX*(SIN(AN)*SIN(AN)) +SY*(COS(AN)*COS(AN))-
 381
 382
                                          2*SXY*SIN(AN)*COS(AN)
                                         SGR = SX^*(COS(AN)^*COS(AN)) + SY^*(SIN(AN)^*SIN(AN)) +
 383
                                          2*SXY*SIN(AN)*COS(AN)
 384
 385
                                         SGSH = 0.5*(SGT-SGR)*SIN(2*AN)+SXY*(COS(2*AN))
 386
                   C
                                   WRITE(6,606) M, JM ,X( M, JM),Y( M, JM ),R,DAN,
 387
 385
                               & &SX.SY.SXY.S1.S3.SGR. SGT.SGSH
```

```
389
           606
                 FORMAT( ' P(',12,',',12,'), =',3F6.2, F6.1,8F7.2)
   390
           С
   391
           30
                     CONTINUE
           C . COUNTER ( MOVING CENTRIOD IXO, IYO UPWARD BY 1)
   392
   393
   394
   395
                    IXO=IXO+1
                   CONTINUE
   396
           10
   397
                   GO TO 888
   398
           999
                   WRITE(6,610)
                   FORMAT(' SOLUTION DIVERGING',/)
   399
           610
                   RETURN
   400
           888
   401
                   END
End of
           C
           C * *
     4
           C
     5
           С
           С
                VERTICAL STRESS DISTRIBUTION ALONG TUNNEL DEPTH
           С
     8
                (USING BIN ARCHING THEORY, REF: HANDY, ASCE, MAR. 1985, P302)
     9
           С
    10
           C*********************
    11
           C
           С
                INPUT FILE:
    12
    13
           С
                 1. DEPTH(TUNNEL), RADIUS, INCLINED ANGLE(YIELDING EAR).
                 2. DENSITY, SOIL-SOIL COEFF., SOIL ANGLE
           С
    14
    15
           С
                 3. DITTO (REPEAT LINES 1 and 2)
    16
           C
                 4.0.0.
                 5 END OF FILE
    17
           C
           С
    18
                 DUTPUT FILE INTERPRETATION:
    19
           С
           С
                 1. VERTICAL STRESS IN TERM OF PI/YH
   20
   21
           €
   22
           C * *
   23
           C
                 READ(5,1) H, RAD, ALPHA
   24
   25
                 IF(RAD LE .O .O) GO TO 999
                 FDRMAT (3F10.0)
   26
   27
                 READ(5,2) DEN, XK, PHI
   28
29
           Ź
                 FORMAT (3F10.0)
                 WRITE(6,101) H, RAD, ALPHA
   30
           101
                 FORMAT(3X, ' HEIGHT= ', F5.1, '
                                                 RADIUS=', F7.2,' INCLINED=', F6.1)
   31
                 WRITE(6, 102), DEN, XK, PHI
   32
                 FORMAT(3X, 'DENSITY =',F5.2,'
           102
                                                XK≃′,F4.2,
                & ' SOIL FRICTION=',F5.1,//)
   33
   34
             XK - COEFF. OF HORIZ. STRESS / VERTICAL STRESS AT FAILURE BOUND
   35
           С
   36
           С
              ALPHA - INCLINED ANGLE OF THE CONE
   37.
                 PHI=PHI+3.1416/180.
   38
   39
                 ALPHA=ALPHA*3.1416/180.
   40%
           С
   41
                 R=H*(1.0/TAN(ALPHA))
   42
                 DZ=H/20.
   43
                 SIGV=0.0
   44
           С
   45
                 WRITE(6,601)
                 FORMAT(' DEPTH ',40X,' P/DEN(Z) ',/)
   46
           601
```

```
С
     47
                   DO 22 I=1,20
     48
     49
                   Z=DZ•I
                   AREA=2. *RAD+2.*(H-Z)/TAN(ALPHA)
     50
                   IF (AREA LE .O.O) GO TO 44
    51
                   FSOIL *XC*SIGV*TAN(PHI)*DZ
DSIGV*(1) O/AREA)*(DEN*DZ*AREA-FSOIL)
    52
    53
                   SIGV=SIGV+DSIGV
    54
                   PV =SIGV/(DEN+Z)
    55
            C
    56
                   WRITE(6, 103) AREA, FSOIL, DSIGV, STOV, PV
    57
                   FORMAT(6F8.2)
          , 103
    58
                   CONTINUE
            22
            44
                   GO TO 11
    60
                   STOP
            999
    61
                   END
    62
End of
               IDENTIFY MODES OF FAILURE (AUG. 1985)
CALCULATE FORCES DUE TO GLOBAL GRAVITY EFFECT (YIELDING ZONES REACHING SURFACE FORMING COLLAPSE MECHANISM)
     5
     7
                SPECIFY KO
                INCLUDING VARIATION OF STRESS ALONG TUNNEL PERIPHERY
     9
              REDUCING SUPPORT PRESSURE BY PROPORTIONS ALONG POSITION
     1.0
            С
    11
              INPUT FILE
    12
            C
                     1. ANGLE OF FRICTION OF SOIL, KO, DEPTHRATIO,
            С
    13
                        ( ALL ARE INITIAL VALUES)
    14
            С
                    2. END OF FILE
    15
            С
            С
    16
    17.
            C
            C OUTPUT INTERPRETATION
   -18
                    1. GRAVITY EFFECT RESULTS ONLY IF PI/HY LESS THAN
    19
            C·
                        P2(I) (MODE: WEDGE)
            С
    20
                     2. MAX. P2(I) GIVES THR CRITICAL YIELDING ANGLE
    21
            С
    22
            С
            С
    23
                    DIMENSION P2(10), FETA(10), XFETA(10), XT(10)
    24
                   READ (5,501) PHIO, XKOO, DEPO
    25
                    FORMAT (3F10.0)
            501
    26
    27
            C
               DEP -DEPTH/RADIUS RATIO
    28
    29
                    DDEP *1.0
                    DPHI #10.
    30
                    DFETA # 10.
    31
                    DXK0=0.1
    32
            C
    33
                    DO 11 I1=1,3
    34
                    PHI = PHIO + DPHI + (I1-1)
    35
                    XPHI = PHI+3.1416/180.
    36
                    WRITE(6,602) PHI
    37
            602
                    FORMAT(//, ' ANGLE PHI =',F5.1)
    38
    39
            С
                    DO 22 I2=1.8 -
    40
                    DO 22 12-1.0
XKO=XKOO+DXKO*(12-1)
    41
                    WRITE(6,666) XKO
    42
                    FORMAT( ' XKO =', F5 2)
            666
    43
```

```
DO 33 T3=1,9
                   DEP=DEPO+DDEP*(13-1)
    45
    46
                   WRITE(6,601) DEP
    47
           601
                   FORMAT(' DEPTH RATIO = ',F5.1)
    48
           ٠C
    49
           С
    50
                   GF = 0.5*(DEP-1.0)*(DEP-1.0)
    51
           С
    52
                   DO 55 L=1,10
    53
                   FETAD = 45.+0.5*PHI ~ 5.0
    54
                   FETA(L) =FETAO +(L-3)+DFETA
                   XFETA(L) = FETA(L)+3.1416/180.
    55
                   XT(L) = TAN(XFETA(L)-XPHI)
    56
    57
           55
                   CONTINUE
                   WRITE(6,612) FETA(1), FETA(2), FETA(3), FETA(4), FETA(5),
    58
    59
                 & FETA(6), FETA(7), FETA(8), FETA(9), FETA(10).
    60
                   FORMAT('ANGLE', 10F8.1,/)
           612
    61
           С
           С
    62
    63
           С
              CH - RATIO OF LATERAL SUPPORT STRESS TO PO
              CV - RATIO OF VERTICAL SUPPORT STRESS TO PO
    64
           С
           С
    65
    66
                   CH= (1,0-0.5/DEP)*XKD
                   CV= (1.0-0.5/DEP)
    67
           С
    68
    69
           С
               MODE 1- SLIDING WEDGES ON SIDE
    70
           15
    71
           С
    72
                   DO 66 M=1,10
                   PART1 = 0.5*DEP*DEP*(1./TAN(XFETA(M)))+ (DEP -3.1416/4.)
    73
                   PART2 = (XKO*GF+XKO*DEP)
    74
                   P2(M) =((PART1*XT(M) - PART2)/(CH+CV*XT(M)-XKO))/DEP
    75
    76
                   CONTINUE
           66
    77
           С
    78
           , C
                   WRITE(6,611) P2(1),P2(2),P2(3),P2(4),P2(5).
    79
                 & P2(6),P2(7),P2(8),P2(9),P2(10)
    80
    81
           611
                   FORMAT('PI/YH', 10F8.3)
    82
           44
                   CONTINUE
    83
           33
                   CONTINUE
    84
                   DEP =DEPO
    85
           22
                   CONTINUE
                   XKO=XKOO
    86
   87
                   CONTINUE
   -88
                   STOP
                   END
    89
End of
           С
           C******** SFT1 *********
     4
           C SHAFT - GROUND REACTION CURVE FORMULATION*
     5
     6
     7
           С
              SOIL MODEL - LADANYI (1974)
                        (MAY, 1985)
     8
           C
     9
           C+
             "ASUMMING:
    10
           C
              HOLE IN PLATE 2-D PLANE STRAIN
    11
           С
    12
           С
              CONSTANT VOLUME CHANGE
              ASSOCIATED FLOW RULE
    13
```

```
INPUT FILE:
 15
             1. KO, MODULUS, DENSITY, RADIUS(=1.0, NORMALISED)
 16
         С
             2. PEAK ANGLE(SOIL), RESIDUAL ANGLE, PEAK COHESION, RESI. COH
 17
         С
                  POISSON RATIO
 18
         Ct
             3 DITTO
         С
 19
 20
         С
             4.40.
             5. END OF FILE
 21
         С
 22
 23
         C APH1= SUPPORT PRESS(PI)/INT. STRESS(PO)
 24
         C APH2= COHESION(CE)/INT. STRESS(PO) ELAST
 25
         C APH3= COHESION(CP)/INT. STRESS(PO) PLAST
 26
 27
         C CE/PO *KA**.5/2*KO
 28
         C APH4=PO/E
         C FANGLE=FRICTION ANGLE ELASTIC
 29
 30
         C PANGLE*FRICTION ANGLE PLASTIC
 31
        C XMIL*POISSON RATIO
        C XKO=COEFF. EARTH PRESSURE AT REST
32
 33
        C E = MODULUS
        C YDEN#DENSITY
 34
       C DEP = DEPTH/RADIUS OPENING
 35
 36
 37
               DIMENSION APH(50), PSIGR(10), PSIGT(10), ESIGR(10)
38
               DIMENSION ESIGT(10), PXR(10), EXR(10)
39
               REAL NOE, NOP, NUM
 40
               WRITE(6,7)
41
               FORMAT( 'GROUND CONVERGENCE CURVE CALCULATION'./)
               READ(5,77) XKD,E,YDEN,RD
42
43
        77
               FORMAT (4F10,0)
44
               WRITE(6,777) XKO,E, YDEN, RD
             FORMAT(3X, 'KO = '.F6.2, 'E = '.F8.1, 'DENSITY = '.F6.2.
8 'RADIUS = '.F6.2)
        777
45
46
47
               IF(XKO.LE.O.O) GO TO 1000
        С
48
49
        333
               READ(5,1) EANGLE, PANGLE, CE, CP, XMIL
               FORMAT (5F 10.0)
50
              IF (EANGLE.LE.O.O) GO TO 1000
51
52
               WRITE(6, 101) EANGLE, PANGLE, CE, CP, XMIL
53
        101
              FORMAT (4X.
             8' PEAK ANGLE =',F5.2,' RESID. ANGLE =',F5.2,
8' CE =',F6.2,' CP =',F6.2,
54
55
             &' POISSON RATIO =',F4.2,/)
56
        C
57
58
              EANGLE=EANGLE+3.1416/180.
59
              PANGLE = PANGLE + 3. 1416/180.
              NOE=(1.0+SIN(EANGLE))/(1.-SIN(EANGLE))
60
61
              NOP=(1.0+SIN(PANGLE))/(1.-SIN(PANGLE))
              DEP=0.00
62
       C
63
            APH1 INPUT DATA
64
            NAPH1 = NUMBER OF APH(I) ENTRY
65
              NAPH1 = 27
              APH(1)=0.0002
66
67
              APH(2)=0.0004
68
              APH(3)=0.0006
69
              APH(4)=0.0008
7.0
              APH(5)=0.0010
71
              APH(6)=0.0015
72
              APH(7)=0.002
```

APH(8)=0.0025

```
APH(9)=0.003
 75
              APH(10)=0.0035
 76
              APH(11)=0.004
 77
              APH(12)=0.005
 78
              APH(13)=0.006
              APH(14)=0.007
 79
 80
              APH(15)=0.008
 81
              APH(16) *0.009
              APH(17)=0.01
 82
              APH(18)=0.02
 83
 84
              APH(19)=0.03
              APH(20)=0.04
 85
              APH(21)=0.05
 86
 87
              APH(22)*0.06
 88
              APH(23)=0.07
              APH(24)=0.08
 89
              APH(25)=0.09
 90
 91
              APH(26)=0.10
 92
              APH(27)=0.15
 93
        C PRINT TITLE COLUMN
 94
              WRITE(6,601)
 95
              FORMAT(',PI/PO ',3X,' HT/RADIUS ',4X,'RP',8X,'UI',4X,
             & 'LAT PRESS'./)
 96
 97
        С
 98
          DEPTH(SHAFT)/RADIUS = 15
        С
 99
        C
100
              DO 44 JJ=1,15
101
              DEP=DEP+1.
102
              PO=YDEN*DEP*RD
              APH2=CE/PO
103
104
              APH3*CP/PO
105
              APH4=PO/E
        C SCP=A (PO)
106
107
              A=APH3/TAN(PANGLE)
108
        C SIGMAC=B (PO)
              B=APH2*2*COS(EANGLE)/(1.O-SIN(EANGLE))
109
110
        C MC*SIGMAC=C (PD)
111
              C=(B+(NDE-1.0)*XKO)/(NDE+1.0)
        C UENOS= ELASTIC DISPLACEMENT AT WALL OF NO SUPPORT
112
113
              UENOS=2.0*APH4*(1.0+XMIL)*XKO
114
        C PLIM=. PI/PO AT PLASTIC YIELD INITIATION (*APH1).
              PLIM=XKO-C
115
116
        С
117
              DD 999 I=1, NAPH1
              APH1=APH(I)
118
119
              IF(APH1.GE.PLIM) GO TO 44
        C RP(EXTENT OF PLASTIC ZONE)=RAD(PLAST)/RAD(OPENING) (RE/RI)
120
              EX=1.0/(NDP-1.0)
121
              RP=((XKO+A-C)/(APH1+A))**EX
122
123
              IF(RP.LE.1.0) GO TO 44
        C RSTS(RAD. STRESS AT ELAST/PLAST INT) = D (PO)
124
125
              RSTS=(XKO-C)
126
          CALCULATE WALL DISPLACEMENT UI
127
        С
128
129
        C APH4= PO/E
130
131
        C UE=RAD. DISPLACEMENT AT ELAST/PLAST INT.
              UE=(1.0+XMIL)+C+APH4+RP
132
133
              DC= -SIN(EANGLE)
```

```
RR=2.0*DC*ALOG(RP)
    135
    136
                  €0 TO 33
                  XA-1.1*DC
DEN=2.0*((1.+XMIL)*C*APH4)*RP*RP
    137
            99
    138
            33
                  NUM=((RP*RP)-1.0)*(1.0+1.0/XR)
    139
    140
                  EAV=DEN/NUM -
    141
            C UI PAD. DISPLACEMENT AT WALL/RAD. OPENING
    142
                A=(2.0*UE/RP-EAV)*RP*RP
    143
                   I=(1.0~SQRT(ABS((1.0-EAV)/(1.0+AA))))
    144
    145
            С
    146
                  HORP - APH1 PO
    147
            C
   148
                  WRITE(6,602) APH1, DEP, RP, UI, HORP
    149
            602
                  FORMAT(F10.6,F10.1,F10.2,F10.5,F10.3)
                  WRITE(6, 102) APH1, APH2, APH3, APH4, DEP, RP, UENOS, UI, PLIM
    150
            С
                 & DEN, NUM, AA, RSTS, PLIM
    151
            С
         _ C102 FORMAT(9F8.5)
    152
    153
           C CALCULATE STRESS DISTRIBUTION
    154
            C XR-WITHIN PLASTIC ZONE XRR-WITHIN ELASTIC ZONE
            C PSIGMAR-PLASTIC STRESS RADIAL
    155
   156
            С
                  XINR=(RP-1.0)/4
    157
            C
                  DO 55 IJ=1,5
                  PXR(IJ)=1.0+XINR*(IJ-1)
   158
            C
                  EXR(IJ) = (RP+IJ)
   159
           С
   160
                  PSIGR(IJ)=(APH1+A)*((PXR(IJ)**(NDP-1.0)))~A
                  ESIGR(IJ)=XKQ-C/((EXR(IJ)/RP)+(EXR(IJ)/RP))
   161
           C
                  PSIGT(IJ)=NOP+(APH1+A)+((PXR(IJ))++(NOP-1.0))-A
   162
   163
           С
                  ESIGT(IJ)=XKO+C/((EXR(IJ)/RP)+(EXR(IJ)/RP))
           C 55
   164
                    CONTINUE
   165
           C
                   WRITE(6,104) PXR(1),PXR(2),PXR(3),PXR(4),PXR(5),EXR(1),
   166
           С
                  &EXR(2),EXR(3),EXR(4),EXR(5)
   167
           C
             104
                    FORMAT (10F8.2)
                   WRITE(6,105) PSIGR(1), PSIGR(2), PSIGR(3), PSIGR(4), PSIGR(5),
   168
           C
   169
           С
                  &ESIGR(1).ESIGR(2).ESIGR(3).ESIGR(4).ESIGR(5)
   170
           С
             105
                   FORMAT(10F8.4)
   171
                   WRITE(6, 106), PSIGT(1), PSIGT(2), PSIGT(3), PSIGT(4), PSIGT(5),
           C
   172
                 &ESIGT(1),ESIGT(2),ESIGT(3),ESIGT(4),ESIGT(5)
   173
           C 106
                   FORMAT(10F8.4)
   174
           999
                 CONTINUE
   175
           44
                 CONTINUE
   176
                 GO TO 333
           1000
                 STOP
   177
   178
                 END
End of
           C*******
           C******* SFT1#********
     5
                VERTICAL AND HORIZONTAL STRESS DISTRIBUTION ALONG CONE DEPTH
           C
                (USING BIN ARCHING THEORY, REF: HANDY, 'ASCE, MAR. 1985, P302)
     7
           C
                (PLASTIC RADIUS CAN BE VARIED IN SHAPE)
    9
           C
           C*********************************
    10
           С
    11
                INPUT FILE:
           C
    12
                1. DENSITY(SOIL), SOIL-SOIL COEFF., SOIL-WALL COEFF., SOIL ANGLE
    13
           C
           С
                    WALL FRICTION ANGLE.
```

IF(RP.GE. 1.7321) GO TO 99

134

```
2. DEPTH(SHAFT), RADIUS.
    16
                  3. NO. OF ENTRY FOR PLASTIC RADIUS(RP)
    17
                  4. RP.
    18
                  5. RP.
    19
                  6. DITTO.
           С
    20
           ¢
                  7. 0.,0.,
    21
           С
                  8 END OF FILE
    22
    23
           С
    24
                  DIMENSION RP(50)
    25
                  READ(5,1)DEN, XK, XKW, PHI, DELTA
                  FORMAT (5F 10.0)
    26
                  READ(5,2) H. RAD
    27
    28
           2
                  FORMAT (2F10.0)
                  WRITE(6, 101) H, RAD
    29
    30
           101
                  FORMAT(3X, ' HEIGHT= ',F5.2,' RADIUS=',F7.3,/)
                 WRITE(6,102) DEN, XK, XKW, PHI, DELTA
FORMAT(3X, 'DENSITY=', F5.1, 'XK=', F4.2, 'XKW=', F4.2.
    31
    32
                 8 ' SOIL FRICTION=',F5.1,' WALL FRICTION=',F5.1,//)
    33
    34
    35
              XK - COEFF. OF HORIZ. STRESS / VERTICAL STRESS AT FAILURE BOUND
    36
           С
              XKW -
    37
           С
    38
           С
    39
           С
                 READ PLASTIC RADIUS (RP)
    40
                  READ(5,3) NRP
    41
           3
                  FORMAT(13)
    42
                  DO 55 I = 1 NRP
    43
                  READ(5,4) RP(I) -
   44
                  FORMAT(F10.0)
    45
           55
                  CONTINUE
   46
           С
   47
           C
    48
                 DZ=H/(NRP+1)
           С
   49
               ∴PHI*PHI*3.1416/180.
   50
   51
                 DELTA=DELTA+3.1416/180.
           С
   52
   53
                 SIGV=0.0
   54
           С
                 WRITE(6,601)
   55
   56
           601
                 FORMAT(' DEPTH '.40X.'VSTRESS',' P/DEN(Z) './)
   57
   58
                 DO 22 I=1,NRP
   59
                 Z=DZ+I
   60
                 EXR=RP(I)
                 ARE4=3.1416*(EXR*EXR-RAD*RAD)
   61
   62
                 IF (AREA.LE.O.O) GO TO 999
   63
                 FSDIL=2.0+3.1416+EXR+XK+SIGV+TAN(PHI)+DZ
   64
                 FWALL=2.0=3.1416=RAD=XKW=SIGV=TAN(DELTA)=DZ
   65
                 DSIGV=(1.0/AREA)*(DEN*DZ*AREA-FWALL-FSOIL)
                 SIGV=SIGV+DSIGV
   66
   67
                 PD=SIGV/(DEN*Z)
   €8
                 HORP=XKW*SIGV
           С
   69
   70
                 WRITE(6, 103) Z.RP(I), EXR, AREA, FWALL, FSOIL, DSIGV, SIGV, PD, HORP
   71
           103
                 FORMAT(11F8.2)
                 CONTINUE
   72
           22
   73
           999
                 STOP
   74
                 END
End of file
```