

1 Metastable dendrite morphologies in Aluminum alloys

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8

9 Abstract

10 Cubic metallic alloys generally grow along $\langle 100 \rangle$ directions due to the anisotropy of the solid-
11 liquid interfacial energy. Under rapid solidification conditions, dendrites may deviate from $\langle 100 \rangle$
12 and develop unusual morphologies. Here, Al-alloy droplets (Al-4.5Cu, Al-10Si, Al-1.9Fe, Al-33Cu,
13 all in wt%) were rapidly solidified using Impulse Atomization to study the microstructures forming
14 at different cooling rates and undercoolings. Growth morphologies of Al-4.5Cu droplets were
15 characterized using X-ray micro-tomography and EBSD. Al-dendrites were found to grow along
16 either $\langle 100 \rangle$ or a more unusual $\langle 111 \rangle$ depending on the solidification conditions. Also, a transition
17 from $\langle 111 \rangle$ to $\langle 100 \rangle$ in the same droplet was observed. These uncommon growth directions were
18 also observed in other Al-alloys. In Al-1.9Fe droplets, a change in dendrite growth direction from
19 $\langle 100 \rangle$ to $\langle 111 \rangle$ was observed, while $\langle 110 \rangle$ growth directions were detected in Al-10Si samples.
20 These experimental observations will be related to their solidification conditions using
21 Solidification Continuous Cooling Transformation diagrams.

22

23 1. Introduction

24 Solidification is a complex phenomenon arising in many modern experimental techniques and
25 industrial technologies related to casting, joining and surfaces processing. Rapid solidification of
26 metallic alloys is an ongoing research interest in the metallurgical sphere. Such non-equilibrium
27 processing conditions can give rise to solubility extension or the formation of metastable phases
28 due to nucleation and/or growth kinetics, where the interface is still at local equilibrium. True
29 departure from equilibrium of the solid-liquid interface can also be achieved, such as solute

30 trapping where the phase diagram no longer applies for the interface and the chemical potentials
31 are no longer equal at the interface. The variation of different conditions of solidification (such
32 as undercooling or cooling rate) gives a possibility to control the morphology and size of crystal
33 structure, which substantially influence physical and chemical properties of alloys.

34 The microstructure evolution during rapid solidification processes depends on the velocity of the
35 solid-liquid interface, which in turn depends on the undercooling ΔT prior to solidification of
36 individual phases in the alloy. Undercooled melts see a large driving force for solidification
37 created from the difference in Gibbs free energy between the solid and the metastable liquid
38 states. A number of microstructural changes have been observed in several different systems
39 with increasing undercooling. Under certain conditions, dendrite growth deviates from $\langle 100 \rangle$ and
40 unusual and complex morphologies can develop. Such deviations were first observed during the
41 growth of ionic crystals in aqueous solution, transparent metal analogs. Kahlweit *et al.*
42 highlighted the growth orientation change of $\text{NH}_4\text{Cl-H}_2\text{O}$ crystals from $\langle 100 \rangle$ to $\langle 110 \rangle$ and then
43 $\langle 111 \rangle$ with increasing growth velocity [1]. The authors investigated the growth orientation during
44 isothermal solidification by varying the nucleation undercooling through changes in
45 supersaturation and solidification temperature. In a subsequent investigation of this system,
46 Chan *et al.* suggested that at low undercoolings, the growth orientation is imposed by the
47 anisotropy of interfacial free energy, with a transition from $\langle 100 \rangle$ to $\langle 110 \rangle$ observed as the
48 supersaturation increases [2]. At high undercoolings, the growth of $\langle 111 \rangle$ dendrite is attributed
49 to the “anisotropy of rate constant”, *i.e.* attachment kinetics. Further studies into the $\text{NH}_4\text{Cl-H}_2\text{O}$
50 system by directional solidification showed oscillations between these different growth modes

51 and the velocity range over which they occur correlates well with the isothermal experiments of
52 Kahlweit [3].

53 Unusual dendrite growth morphologies have also been observed in aluminum alloys. Herenguel
54 first reported twinned dendrites, or so-called feathery grains in semi-continuous castings [4].
55 Using Electron BackScattered Diffraction (EBSD), Henry *et al.* showed that feathery grains in Al-
56 Mg-Si, Al-Si-Ti and Al-Cu cast billets were made of $\langle 110 \rangle$ dendrites with their trunk split by a $\{111\}$
57 twin plane [5]. Using a Bridgman type directional solidification and a unidirectional solidification
58 setup, Henry *et al.* obtained other morphologies dependent on the alloy composition, growth
59 velocity and temperature gradient [6]. For instance, so-called degenerate feathery grains were
60 observed in Al-9wt%Si. These non-twinned dendrite morphologies grow along $\langle 110 \rangle$ with
61 secondary arms of $\langle 100 \rangle$ and $\langle 110 \rangle$ types.

62 Sémoroz *et al.* studied dendrite formation in Al-45wt%Zn coating on hot-dipped steel sheets [7].
63 They observed an eight-fold symmetry dendritic pattern in grains having a (001) plane parallel to
64 the surface. Their EBSD analysis unambiguously identified these as $\langle 320 \rangle$ dendrite growth
65 directions. These observations have been further analyzed in Al-Zn alloys by Gonzales *et al.* using
66 directional solidification and Bridgman solidification over a wide range of compositions [8]. $\langle 100 \rangle$
67 dendrites were found for compositions up to 25wt%Zn while $\langle 110 \rangle$ growth directions were seen
68 above 65wt%Zn. In between these two compositions, a continuous change of growth direction
69 from $\langle 100 \rangle$ to $\langle 110 \rangle$ was highlighted. This so-called Dendrite Orientation Transition (DOT) was
70 then successfully simulated by Haxhimali *et al.* using the phase field method [9]. The origin of the
71 DOT has been attributed to a change in the solid-liquid interfacial energy anisotropy [10].

72 More recently, the origin of twinned dendrites in Al alloys has been identified in the same Al-Zn
73 alloys as coming from icosahedral short-range order in the liquid [11], when very small amount
74 of Cr is added. This leads to quasicrystal formation from which the fcc phase grows with epitaxial
75 relationships leading to multiple twinning relationships of grains.

76 Becker *et al.* observed a DOT similar to that of Al-Zn in the Al-Ge system by conducting isothermal
77 solidification experiments in thin samples under slow cooling conditions [12]. Primary dendrite
78 arms are growing along $\langle 100 \rangle$ in Al-20wt%Ge alloys, along $\langle 110 \rangle$ in Al-46wt% Ge, while both
79 directions are growing simultaneously at an intermediate composition of Al-29wt% Ge. These
80 microstructures were successfully modelled using phase-field simulations with varying solid-
81 liquid interfacial energy anisotropies [12].

82 Lately, Wang *et al.* examined the change in dendrite growth direction in laser-melted Al-Sm alloys
83 as a function of samarium content [13]. A transition from $\langle 100 \rangle$ to seaweed was observed at
84 $\sim 2.2\text{at\% Sm}$, and to $\langle 110 \rangle$ at $\sim 5.6\text{at\% Sm}$, which is again attributed to changes in the interfacial
85 free energy anisotropy induced by the solute. This was confirmed using Molecular Dynamic
86 simulations to calculate the two parameters used to characterize the anisotropy, ϵ_1 (fourfold
87 symmetric contribution) and ϵ_2 (sixfold contribution). A clear variation in the $(\epsilon_1, -\epsilon_2)$ space is
88 predicted as the Sm content increases. The interfacial anisotropy changes such that the expected
89 dendritic growth directions shift from $\langle 100 \rangle$ to seaweed to $\langle 110 \rangle$.

90 All these experiments suggest that the transition to $\langle 110 \rangle$ is induced by the variation of interfacial
91 energy anisotropy, while attachment kinetics would explain the change to $\langle 111 \rangle$ orientation seen
92 with increasing solidification rate.

93 More than twenty-five techniques have been reported to induce rapid solidification [14], the
94 most widely used being atomization. Impulse Atomization (IA) is a single fluid atomization
95 technique that yield droplet solidification with large levels of undercooling over a wide range of
96 cooling rates [15] [16]. This paper compiles new and previously reported occurrences of
97 metastable dendrite morphologies observed in rapidly solidified Al-based alloys obtained by IA.
98 The aim is to explore common features in morphology in various hypoeutectic Al alloys, namely
99 Al-Cu, Al-Si, Al-Fe, Al-Zn, Al-Ge, as well as near-eutectic Al-Cu.

100 2. Experimental methods

101 Impulse atomization (IA) (a type of drop tube) is a containerless solidification technique (
102 Figure 1). It consists in the transformation of a bulk liquid into a spray of liquid droplets. A plunger
103 (or impulse applicator) applies a pressure (or impulse) to the melt in order to push it through a
104 nozzle plate with several orifices of known size and geometry. Liquid ligaments emanate from
105 each orifice, which in turn break up into droplets. Rapid solidification of the droplets then occurs
106 during free fall by heat loss to the surrounding gas (usually He, N₂ or Ar). The thermal history of
107 the liquid droplets is a function of both the droplet size and the gas in the atomization tower and
108 has been described mathematically by previous workers [17] [18]. The solidified samples can
109 finally be collected at the bottom of the tower and subsequently sieved into different size classes.
110 A detailed description of the process is available in [15].

111 3. Results and Discussion

112 An in-depth study of Impulse Atomized Al-4.5wt%Cu was carried out by Bedel *et al.* [19]. This
113 alloy is widely used as a model for solidification studies and its thermophysical properties are
114 well documented. Two sets of droplets were atomized using different gas atmospheres, helium
115 and argon. Post-mortem synchrotron X-ray micro-tomography was carried out on these droplets
116 at ESRF (European Synchrotron Radiation Facility, Grenoble, France). Careful analysis of more
117 than 200 particles revealed four main distinct dendritic morphologies (Figure 2). The highly
118 branched morphology (a) shows dendrites growing with a four-fold symmetry typical of $\langle 100 \rangle$
119 growth while microstructural features indicate that dendrite arms develop mostly along $\langle 111 \rangle$
120 directions in the other three morphologies (b-d).

121 The assumed growth directions from the tomography analysis have been confirmed using EBSD
122 on selected droplets. Figure 3 (left) shows the cross section of a 960 μm droplet atomized in
123 argon exhibiting the highly branched morphology along with its corresponding pole figures. This
124 cross section lies in a $\{001\}$ plane. The primary dendrite trunks directions, highlighted by the red
125 arrows, correspond to the $\langle 100 \rangle$ poles. Microtomography cross sections along the XZ and YZ
126 planes (*i.e.* after a 90° rotation around the X or Y axis of this particular cross section) exhibit a
127 similar morphology, suggesting that the dendritic microstructure developed along $\langle 100 \rangle$
128 throughout the whole droplet. Figure 3 (right) shows the cross section and pole figures of a 415
129 μm particle atomized in helium with the finger bundle structure. This particular cut lies this time
130 in a $\{011\}$ plane. Two of the four $\langle 111 \rangle$ poles are well aligned with the bundles, as shown by the
131 blue arrows. This is consistent with a $\langle 111 \rangle$ dendritic growth.

132 Only 7.8% of all the droplets analyzed solidified completely along $\langle 100 \rangle$ directions. The majority
133 of droplets showed at least some instances of $\langle 111 \rangle$ growth. The transition from $\langle 100 \rangle$ to $\langle 111 \rangle$
134 is attributed to the variation of the attachment kinetics anisotropy as the solidification growth
135 velocity increases. $\langle 100 \rangle$ arms develop at low solidification growth velocity (Figure 2.a). As the
136 cooling rate and/or undercooling increases, primary arms start growing along $\langle 111 \rangle$. However,
137 growth reverts to $\langle 100 \rangle$ for slower growing side arms formed after recalescence (Figure 2.b). At
138 even higher solidification rates, the droplet solidifies completely with a $\langle 111 \rangle$ growth direction,
139 as illustrated in Figure 2.c. Finally, at the highest speed, finger bundles stem from a growth
140 competition between different $\langle 111 \rangle$ dendrites originating from the same nucleation point
141 (Figure 2.d). The $\langle 111 \rangle$ dendritic structure is found again in regions of the droplet with slower
142 solidification velocity. The distribution of the four observed morphologies was established by
143 careful analysis of a large number of droplets solidified in argon from two size categories (73
144 droplets with a diameter smaller than 212 μm and 64 droplets between 250 μm and 300 μm ,
145 Figure 4). Due to the stochastic nature of nucleation, all morphologies are found within a single
146 size range. However, growth directions tend to shift towards the faster $\langle 111 \rangle$ morphologies when
147 the droplet size decreases. Indeed, when the droplets are smaller, they solidify faster due to the
148 larger surface to volume ratio. The same trend is observed when switching the gas atmosphere.
149 For a given size range, the morphology is more often of the finger-bundle type in droplets
150 solidified in helium, due to its higher thermal conductivity and hence higher cooling rates.

151 Subsequent work on the rapid solidification of Al-4.5wt%Cu and Al-4.5wt%Cu-0.4wt%Sc
152 estimated the undercooling temperatures of both the primary phase and the eutectic structures
153 in these alloys [20]. In order to represent the resultant microstructure and relate it to macro-

154 solidification conditions, Solidification Continuous Cooling Transformation (SCCT) diagrams were
155 developed [20]. To construct these maps, the liquid cooling rates of the samples are used and
156 the corresponding undercoolings are plotted on a CCT diagram (Figure 5). Also plotted on this
157 graph are the equilibrium liquidus and eutectic temperatures, TL and TE respectively. Finally, the
158 resulting proportions of droplets exhibiting any type of $\langle 111 \rangle$ growth ($\langle 111 \rangle$ to $\langle 100 \rangle$ transition,
159 $\langle 111 \rangle$ dendrites and finger bundles) are also indicated for each cooling rate. Over the range
160 studied, primary nucleation undercoolings are high, starting at 60 K at 500 K/s and steadily
161 increasing up to 86 K at 30,000 K/s. Dendrites growing along $\langle 111 \rangle$ can be found under all
162 solidification conditions analyzed and are the predominant morphology at all but the lowest
163 cooling rate. In addition, samples processed under lower cooling rates (~ 10 -50 K/s) and
164 undercoolings (~ 5 -45 K) using electromagnetic levitation only showed dendrites growing along
165 $\langle 100 \rangle$ [20]. This clearly indicates that the transition from $\langle 100 \rangle$ to $\langle 111 \rangle$ is related to the
166 solidification conditions and supports the hypothesis that the variation of the attachment kinetics
167 anisotropy is responsible for the observed transition.

168 Unusual growth directions were also observed in atomized Al-10wt%Si droplets. One such
169 example is shown in Figure 6 [21]. In this 230 μm particle atomized in He, the primary α -Al shows
170 a clear 6-fold symmetry. EBSD analysis shows that the whole particle is constituted of a single α -
171 Al grain, while the pole figure clearly indicates $\langle 110 \rangle$ growth directions.

172 Figure 7 shows the proportion of droplets exhibiting $\langle 110 \rangle$ growth. While it was not observed in
173 droplets solidified in argon, this $\langle 110 \rangle$ growth only occurred in the finer particle sizes that were
174 atomized in helium. This indicates that this growth mode is dependent on solidification

175 conditions, *i.e.* cooling rate and/or nucleation undercooling. However, it is different from the
176 $\langle 111 \rangle$ observed in Al-Cu and Al-Fe alloys. As mentioned above, $\langle 110 \rangle$ growth was reported in Al-
177 Si and Al-Ge alloys (Ge and Si having many similar properties). The Al-Ge results were obtained
178 at very low solidification rate, with this transition clearly attributed to changes in the surface
179 energy anisotropy with increasing solute content [12]. $\langle 110 \rangle$ growth in unidirectionally solidified
180 Al-9wt%Si was however obtained at a much higher solidification rate [6]. While kinetics effects
181 cannot totally be discounted for this particular transition, it is possible that rapid solidification
182 induces $\langle 110 \rangle$ growth by increasing the supersaturation, which could change the surface energy
183 anisotropy in favor of $\langle 110 \rangle$.

184 A seaweed-type of growth was also observed in many Al-10wt%Si droplets, as shown in Figure 8.
185 This seaweed morphology is characterized by a highly branched structure that grows outwards
186 from the primary arms at an angle below 90° . During seaweed growth, the growing dendrites
187 bend away from the primary arm and split, causing them to grow in a “zig-zag” fashion.

188 Seaweed structures were observed by Friedli *et al.* in Al-Zn alloys directionally solidified at low
189 growth velocities [10]. The zigzagging of the arms was considered a response to perturbations at
190 the solidification front, caused by the solute field of another arm. The rejected solute changes
191 the surface tension anisotropy at the solidification front, which forces the growing α -Al dendrite
192 to bend, switch sides and form a seaweed structure. Hence, solute effects play a role in the Al-Zn
193 alloys. It remains unclear how segregation and undercooling interplay between solute and
194 kinetic effects.

195 Using the phase-field method to model the solidification of pure metallic melts, Mullis *et al.*
196 highlighted a transition to seaweed growth with increasing undercoolings [22]. As undercooling
197 increases, the dendrite growth velocity would also increase and promote the onset of dendrite
198 tip splitting. Furthermore, an increase in growth velocity causes the initial side-branching to move
199 closer to the dendrite tip. Mullis *et al.* considered this shift in the perturbations at the
200 solidification front to be a competition between the surface tension anisotropy and the atom
201 attachment kinetics, where the influence of kinetics became more dominant at higher growth
202 velocities. Experimental work done by Assadi *et al.* [23] with Ni₅₁Al₄₉ alloys found that at higher
203 growth velocities seaweed growth was possible, where the formation of the seaweed structure
204 was also attributed to the increasing role of atom attachment kinetics.

205 To relate the solidification conditions to the seaweed growth observed in Al-10wt%Si, a visual
206 analysis was conducted to determine the number of powders, within a particle size range, that
207 displayed seaweed growth (Figure 9). Experimental conditions noticeably influence the growth
208 of the seaweed structure. Decreasing the particle size and using helium (instead of argon) will
209 make seaweed growth more prevalent. These results indicate that seaweed growth in atomized
210 Al-10wt%Si might be caused by the increasing role of atom attachment kinetics, due to an
211 increase in the α -Al dendrite growth velocity.

212 Similarly to Al-4.5wt%Cu, an SCCT diagram was developed. This combines the above results and
213 outlines the solidification pathway of Impulse Atomized Al-10wt%Si droplets (Figure 10). This plot
214 clearly shows that the metastable morphologies, be it seaweed structures or dendrites growing
215 along $\langle 110 \rangle$, are favored by high cooling rates and undercoolings. This again indicates that the

216 prevalence of metastable morphologies is related to solidification conditions. While attachment
217 kinetics effects are likely responsible for the formation of seaweed, the exact reason for the
218 transition to $\langle 110 \rangle$ remain unclear.

219 The microstructure evolution of impulse atomized powders of Al-0.61wt%Fe has been
220 investigated by Henein *et al.* [24] and Chen *et al.* [25]. The solidification cooling rate of a 355
221 μm droplet of Al-0.6wt%Fe atomized in helium was reported as about $4,000 \text{ K s}^{-1}$. Figure 11 shows
222 an optical image of a 355 μm droplet atomized in He. Primary trunks are extending in three
223 different directions from a single nucleation site. EBSD analysis around the nucleation site was
224 performed (Figure 11). On this reconstructed map, the orientation of the nucleation center was
225 chosen as a reference. The false coloring corresponds to the misorientation angle from the
226 nucleation center. The color scheme shows that misorientation for most of the region, especially
227 the three main trunks, is within 1 degree. While no clear side arms can be seen, it is apparent
228 that the growth direction of the primary trunks coincide with the three outer poles of the $\langle 111 \rangle$
229 pole figure. The misorientation of the side arms is attributed to fragmentation. Local remelting
230 and break-off of dendrite arms takes place in the mushy zone during recalescence [26].
231 Fragments that do not fully melt can grow further, forming new grains. As the free-falling droplets
232 are almost devoid of convection, their orientation would not deviate greatly from the parent
233 dendrite. Thus, no major misorientation is observed.

234 Another instance of deviation from $\langle 100 \rangle$ growth has been observed when atomizing an Al-Cu
235 melt of eutectic composition (Al-33wt%Cu). While the vast majority of droplets exhibited a fully
236 eutectic microstructure, some droplets showed nucleation of α -Al dendrites prior to eutectic

237 solidification. This indicates that the melt composition was somehow locally hypo-eutectic. An
238 example is shown in Figure 12. It is clear that the dendrite arms do not grow at 90° from each
239 other, indicating a growth direction different from $\langle 100 \rangle$. EBSD mapping of this sample shows
240 that two different nucleation events happened. The α -Al dendrites on the left being part of one
241 grain (blue) and the ones on the right constituting a second grain (green). Looking at the pole
242 figure, it appears that the dendrites are growing along $\langle 110 \rangle$. This is most intriguing as both
243 aluminum and copper are FCC crystals with interfacial energy anisotropies favoring $\langle 100 \rangle$ growth.

244 4. Conclusions

245 Al-alloy droplets of various compositions were rapidly solidified using Impulse Atomization to
246 study the microstructures forming at different cooling rates and undercoolings. Deviation from
247 the usual $\langle 100 \rangle$ growth direction was observed. In Al-4.5wt%Cu and Al-0.61wt%Fe, dendrites can
248 grow along $\langle 111 \rangle$ depending on the solidification conditions. Higher cooling rates and/or
249 undercooling favour a transition to $\langle 111 \rangle$. This suggests that anisotropy of attachment kinetics is
250 a contributing factor, as observed in transparent alloys. $\langle 110 \rangle$ growth directions were detected
251 in Al-10wt%Si, with occurrences increasing with the solidification rate. While attachment kinetics
252 cannot be discarded, it is usually linked to $\langle 111 \rangle$ growth. Germanium, which has similar properties
253 than silicon, was shown to induce a DOT in Al-Ge alloys. Thus it is suggested that rapid
254 solidification conditions might locally change the interfacial energy anisotropy with increasing
255 supersaturation. The presence of $\langle 110 \rangle$ α -Al in near eutectic Al-Cu is still puzzling, as both
256 elements have anisotropies that favour $\langle 100 \rangle$ growth.

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266 Figure 1: Schematic view of an impulse atomization apparatus.

267 Figure 2: Synchrotron X-ray tomography images depicting the four typical dendrite morphologies
268 observed in Al-4.5wt%Cu droplets solidified in IA: (a) $\langle 100 \rangle$ highly branched dendrites; (b) $\langle 111 \rangle$
269 to $\langle 100 \rangle$ dendrite transition; (c) $\langle 111 \rangle$ dendritic morphology; (d) $\langle 111 \rangle$ finger bundle morphology.
270 The black rings are reconstruction artifacts [20].

271 Figure 3: Left: 960 μm Al-4.5wt%Cu particle atomized in Ar and EBSD pole figures for α -Al showing
272 the $\langle 100 \rangle$ growth of the highly branched morphology; Right: 415 μm Al-4.5wt%Cu particle
273 atomized in He and EBSD pole figures for α -Al showing the $\langle 111 \rangle$ growth of the finger bundle
274 morphology [19].

275 Figure 4: Distribution of the four morphologies in Al-4.5wt%Cu droplets solidified in argon for
276 two diameter ranges (73 droplets for $0 < d < 212 \mu\text{m}$ and 64 droplets for $250 \mu\text{m} < d < 300 \mu\text{m}$).

277 Figure 5: Solidification Continuous Cooling Transformation curves of Impulse Atomized Al-
278 4.5wt%Cu.

279 Figure 6: $230 \mu\text{m}$ Al-10wt%Si particle atomized in He. (left): SEM micrograph; (right): $\alpha\text{-Al} \langle 110 \rangle$
280 pole figure.

281 Figure 7: Proportion of Al-10wt%Si powders exhibiting $\langle 110 \rangle$ growth directions as a function of
282 particle size and atomization gas.

283 Figure 8: Optical micrograph outlining the seaweed growth of $\alpha\text{-Al}$ phase in an $212\text{-}250\mu\text{m}$ Al-
284 10wt%Si droplet atomized in helium.

285 Figure 9: Proportion of Al-10wt%Si powders exhibiting seaweed-type growth as a function of
286 particle size and atomization gas.

287 Figure 10: Solidification Continuous Cooling Transformation curves of Impulse Atomized Al-
288 10wt%Si.

289 Figure 11: $355 \mu\text{m}$ Al-0.61wt%Fe particle atomized in He. (left): optical; (middle): EBSD false color
290 reconstruction; (right): $\alpha\text{-Al} \langle 111 \rangle$ pole figure.

291 Figure 12: $275 \mu\text{m}$ Al-33wt%Cu particle atomized in Ar. (left): SEM micrograph; (middle): EBSD
292 false color reconstruction; (right): $\langle 110 \rangle$ pole figure.

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